Particle acceleration in Galaxy Clusters & connection with cluster mergers



Institute of Radioastronomy –INAF, Bologna, ITALY

Outline

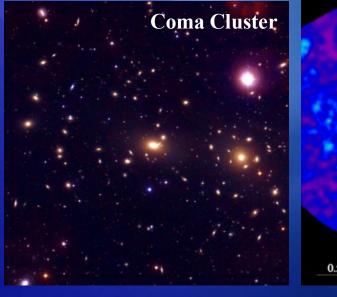
- Galaxy Clusters & non thermal phenomena
- Physics & dynamics of CR in galaxy clusters
- CR protons (limits)
- CR electrons : Radio Halos (turbulence?)
- General Picture of non-thermal clusters (calculations of broad band non thermal emission)
- Expectations : ... steep spectrum diffuse radio emission & LOFAR+LWA

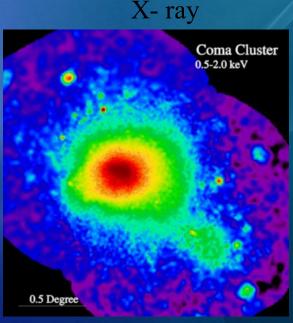
Clusters of galaxies:

• Galaxy Clusters are the largest concetrations of matter in our Universe.

- They form by accretion & mergers between sub-units
- They extend over 2-4 Mpc and have a total mass of ~ 10¹⁴-10¹⁵ M_{sol}
- They contain thousands of galaxies, hot diffuse gas and expecially dark matter

Optical





Galaxy cluster mass: Barions 10% of stars in galaxies 15-20% of hot

diffuse gas

Dark Matter 70%

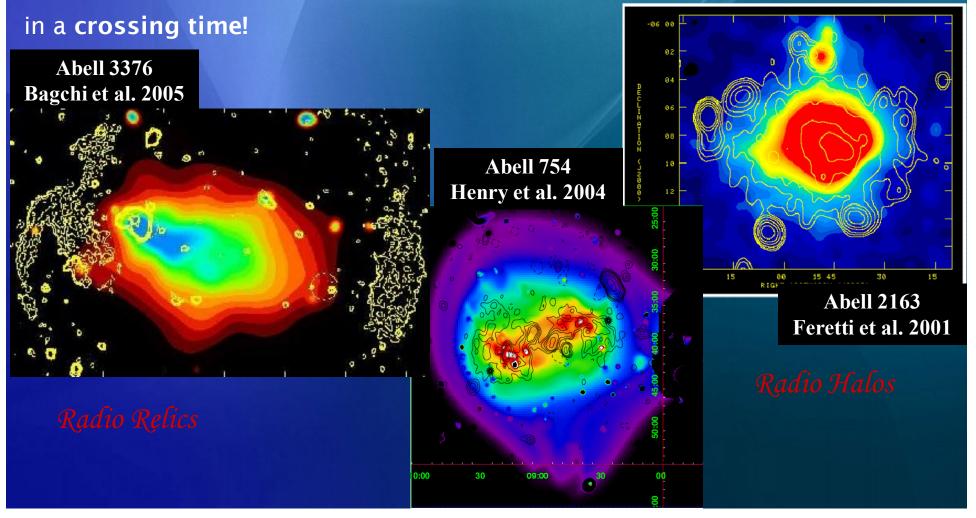
stars + dark matter

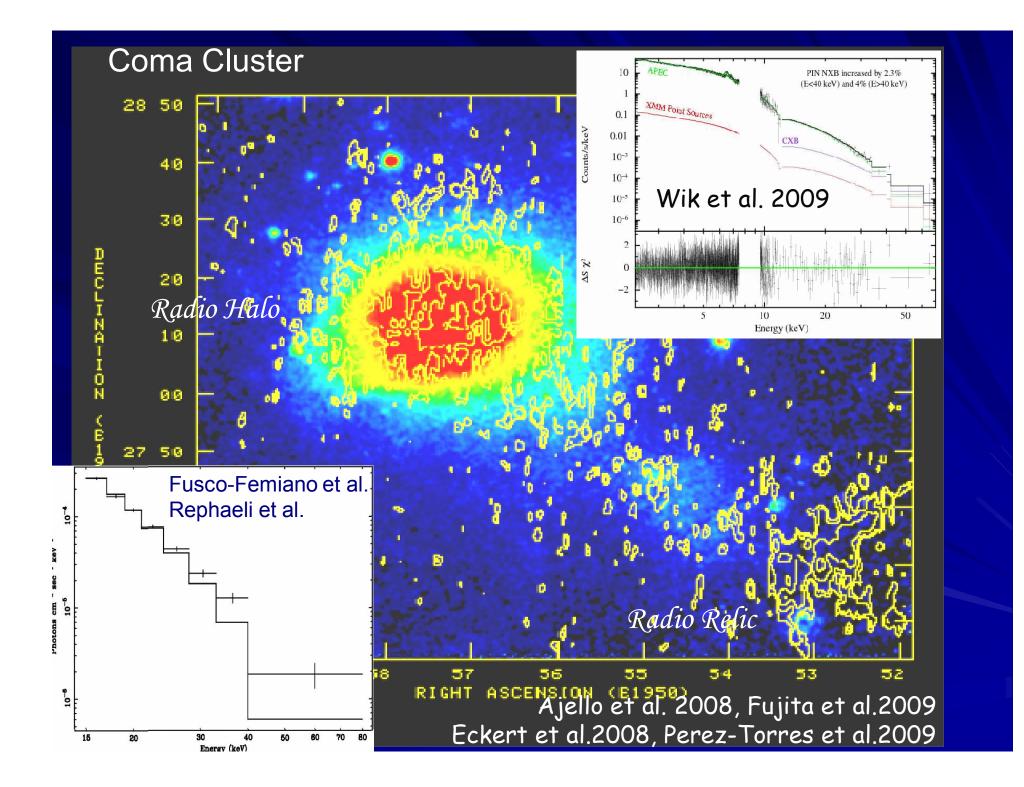
hot diffuse gas

Non-thermal components & cluster mergers

Radio Halos and Radio Relics are only found in non-relaxed clusters with recent /ongoing cluster mergers (e.g. Buote 2001)

Mergings are the most energetic events in the Universe \approx few 10⁶³ergs





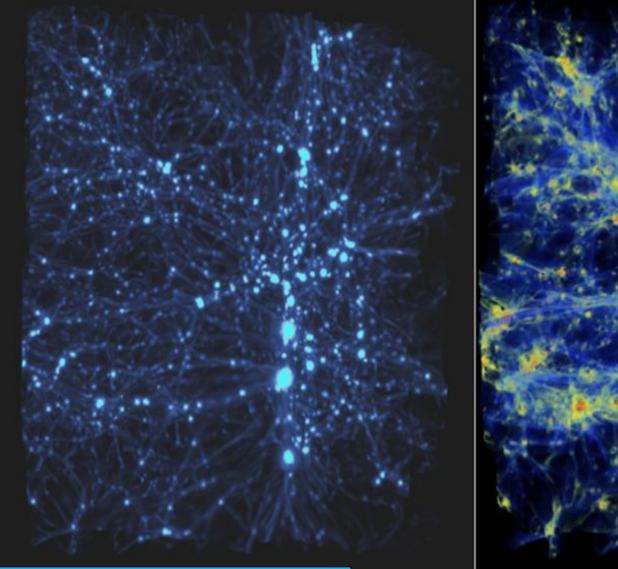
Injection & Dynamics of Cosmic Rays in GC

Cosmological Shocks

(e.g. Sarazin 1999, Miniati et al. 2001, Blasi 2001, Gabici & Blasi 2003, Ryu et al. 2003, Pfrommer et al. 2006, 2008, Vazza, Brunetti, Gheller 2008)

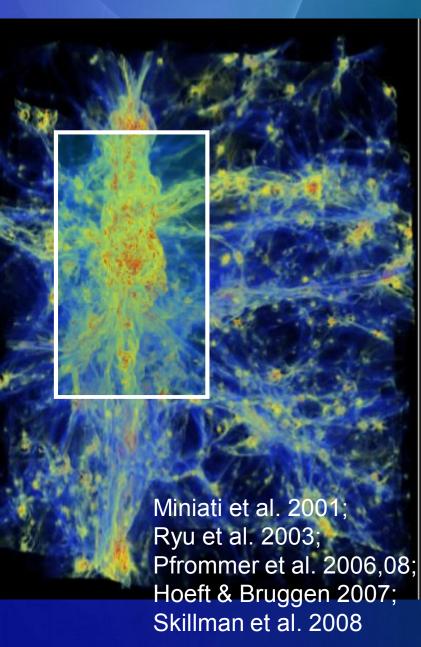
AGN, Galactic Winds (e.g. Ensslin et al. 1998; Voelk & Atoyan 1999)

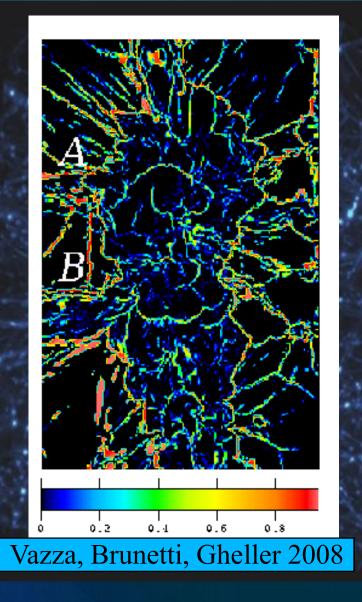
Simulations : Shocks in Galaxy Clusters



Vazza, Brunetti, Gheller 2008

Shocks in Galaxy Clusters

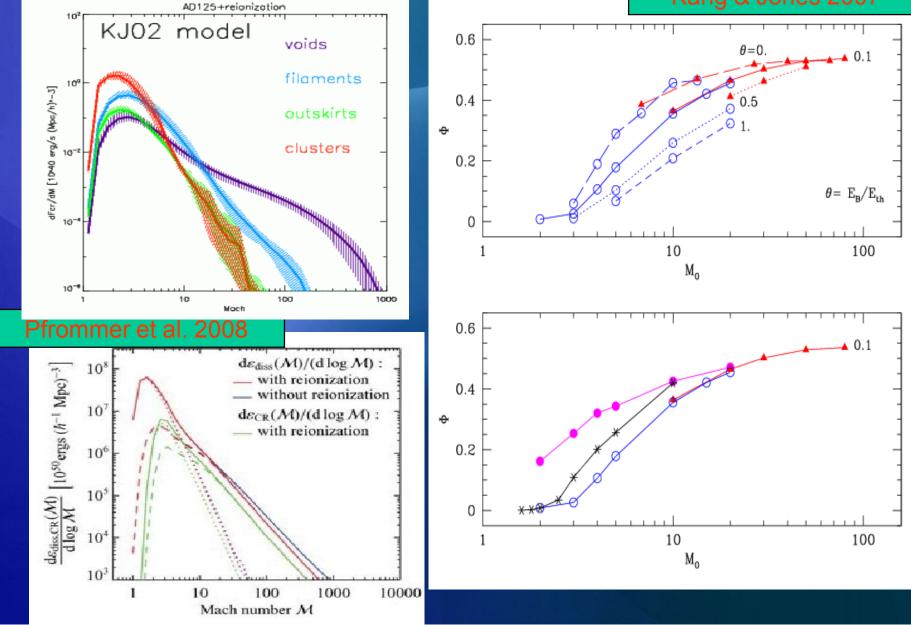


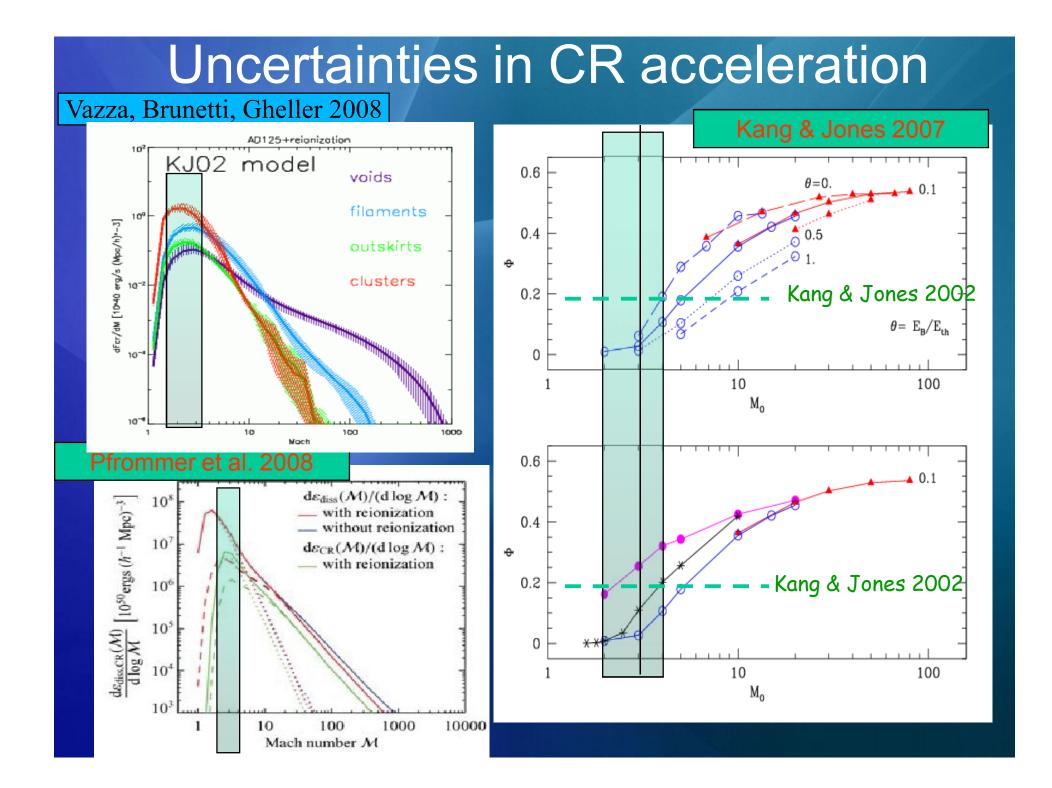


Uncertainties in CR acceleration

Vazza, Brunetti, Gheller 2008

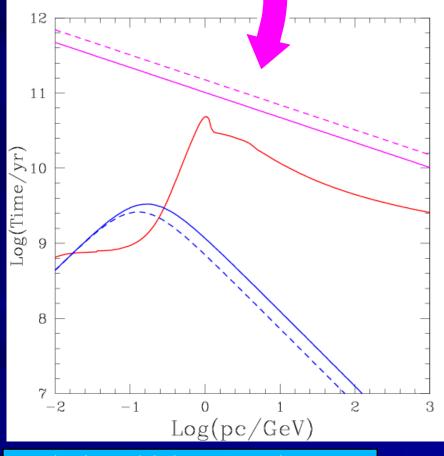






Cosmic Ray Physics & Confinment

Diffusion time



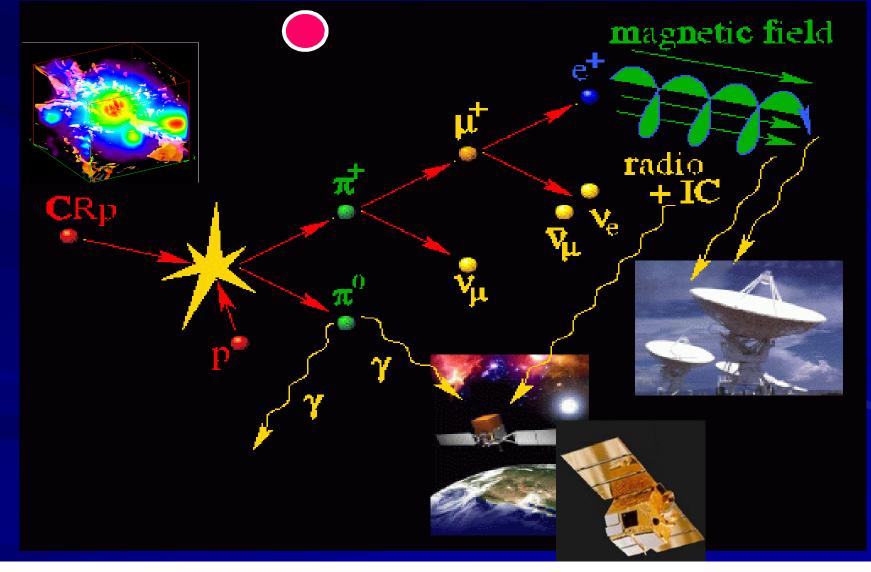
CR protons are <u>long living</u> particles and are <u>confined</u> (Voelk et al 1996; Berezinsky, Blasi, Ptuskin 1997)

CR electrons are <u>short living</u> particles and <u>accumulated</u> at $\gamma \approx 100-300$ (e.g., Sarazin 1999)

Blasi, Gabici, Brunetti 2007

Secondary Particles in the IGM

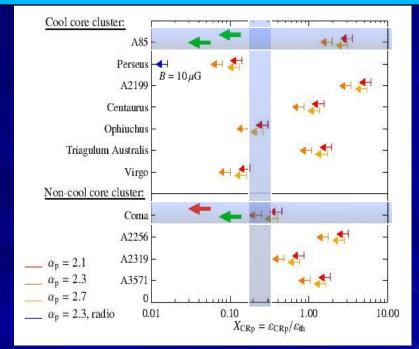
Dennison 1980; Blasi & Colafrancesco 1999; Pfrommer & Ensslin 2004; Brunetti & Blasi 2005; Wolfe & Melia 2007; Brunetti 2008



CRp: Limits from gamma rays

Aharonian + al. 2008 Declination (J2000) Abell 85 4 3 -9°00' 2 _ 10 -1 -9°30' -2 -3 -4 -10°00' -5 00h44m 00h42m 00h40m Right Ascension (J2000)

Reimer +al. 2003; Pfrommer & Ensslin 2004



H.E.S.S.

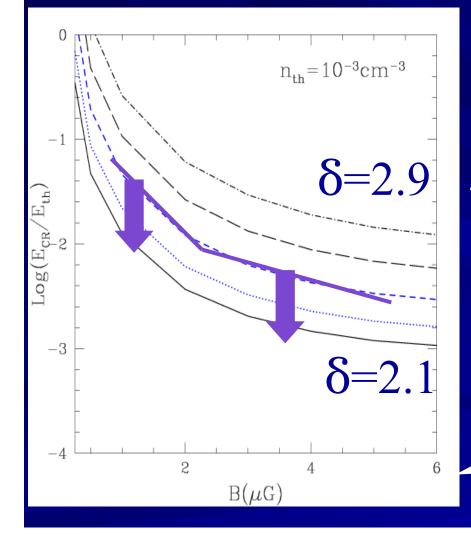
A 85 : Ecr/Eth < 6-15% (hard spectra Coma : Ecr/Eth < 12%

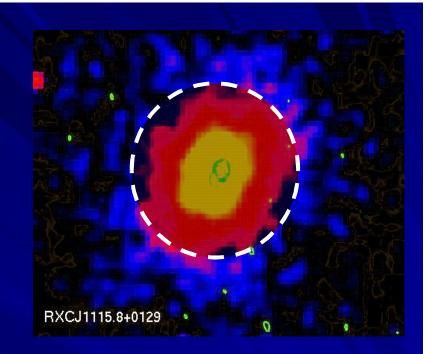
VERITAS (Perkins +al. 2008) Coma : Ecr/Eth < 5% (hard spectra

CRp: limits from Radio

Brunetti +al. 2007

Making use of clusters in the GMRT survey



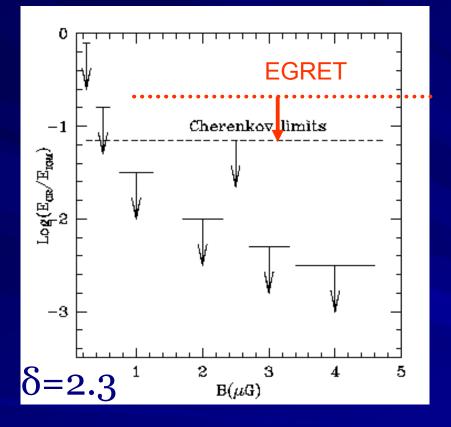


Assuming that secondary particles are injected in the IGM, their synchrotron emission should be smaller than upper limits to the diffuse radio emission.

> limit on : B, E_{CRp}, δ $N(p)=K p^{-\delta}$

Limits for CR protons

Brunetti, Blasi, Cassano, Gabici 2008



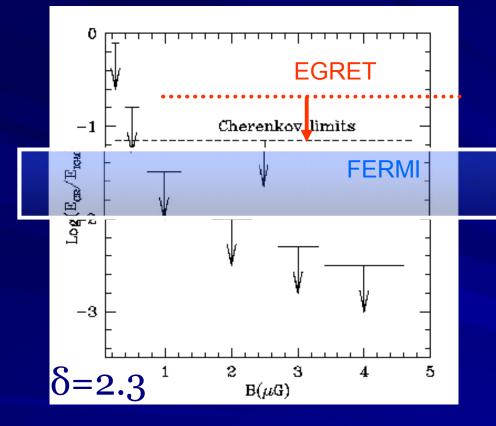
Reimer et al. 2003 Reimer et al. 2004 Pfrommer & Ensslin 2004 Perkins et al. 2006 Brunetti et al. 2007 Brunetti et al. 2008 Perkins et al. 2008 Aharonian et al. 2008

In case of flat proton spectrum present limits constrain E_{CR} < few % Ethermal

Additional limits from cluster dynamics (e.g. Churazov et al. 2008) constrain $E_{CR}+E_B+E_{turb}$ below 10% Ethermal.

Limits for CR protons

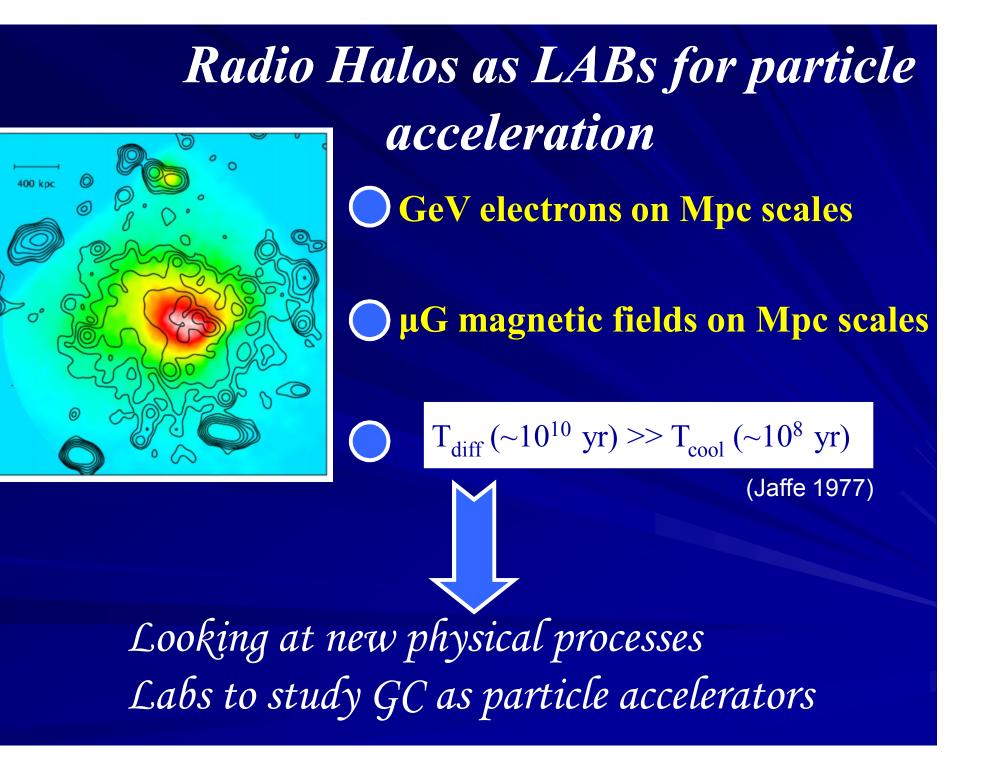
Brunetti, Blasi, Cassano, Gabici 2008



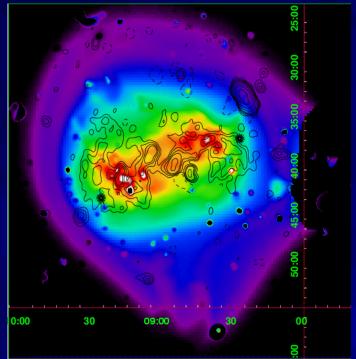
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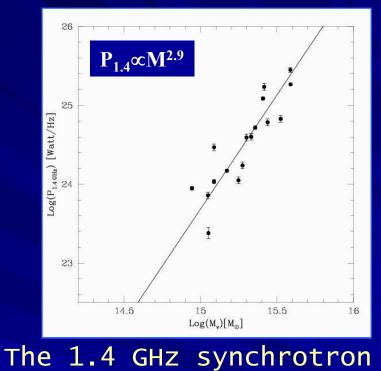
In case of flat proton spectrum present limits constrain E_{CR} < few % Ethermal. In case of flat spectra FERMI is expected to step into the region of the Radio limits, providing better constraints in case of steeper spectra.

Additional limits from cluster dynamics (e.g. Churazov et al. 2008) constrain $E_{CR}+E_B+E_{turb}$ below 10% Ethermal.



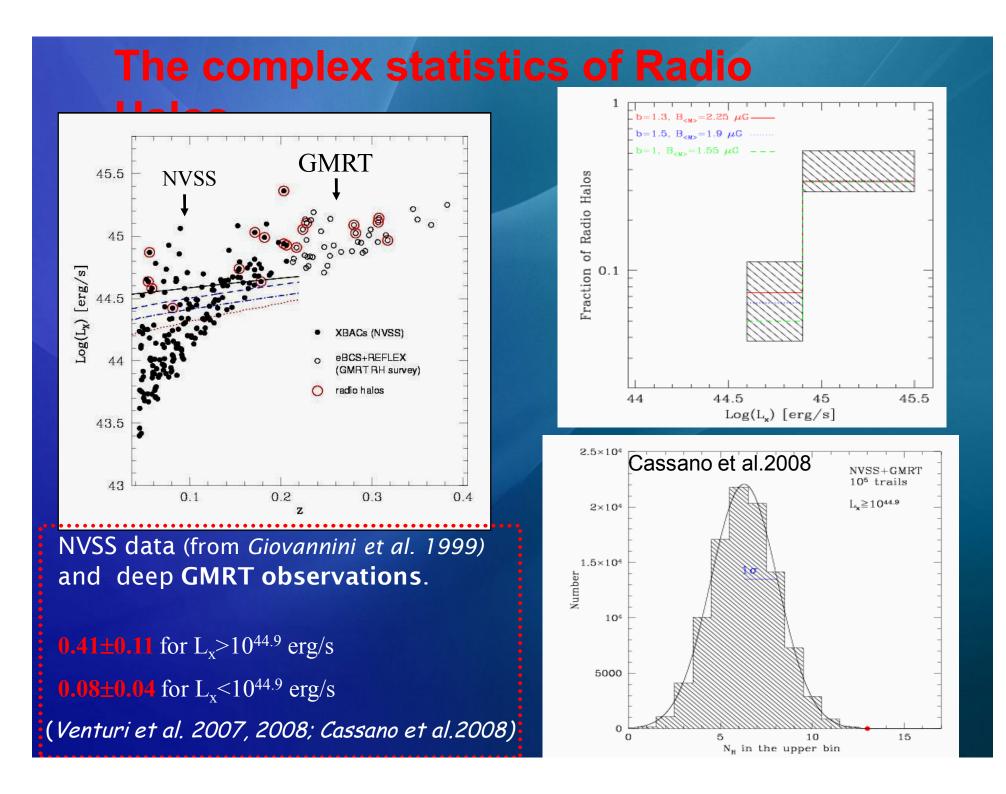
Radio Halos and cluster-cluster mergers





Connection with cluster
mergers (e.g. Buote 2001, Schuecher et al.The 1.4 GHz synchrotron radio
power of GRH increases with
the cluster mass (L_x,T) (e.g. Liang
1999, Bacchi et al. 2003, Cassano et al. 2006)2003 Govoni et al. 2004)

Gravitational - driven processes ?

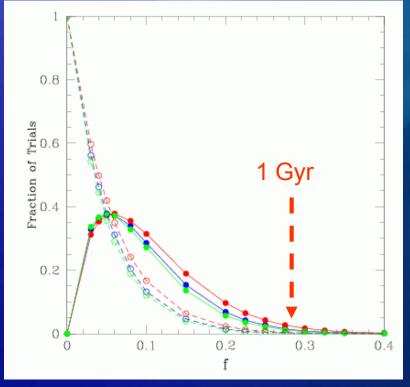


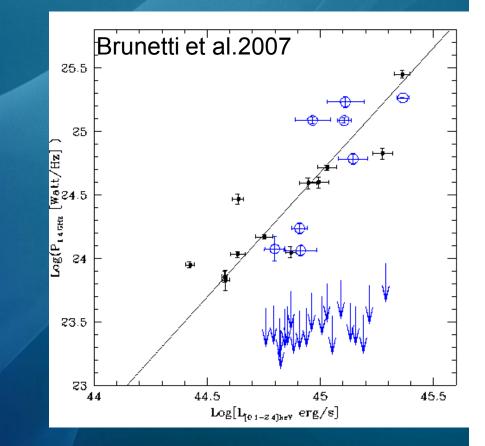
The complex statistics of Radio

Connection with cluster mergers (e.g. Schuecher et al. 2001, Markevitch et al. 2002, Boschin et al. 2003 Govoni et al. 2004) Which is the difference between RH-clusters and ULimits ?

Does non thermal emission evolve ?

Which time-scale ?



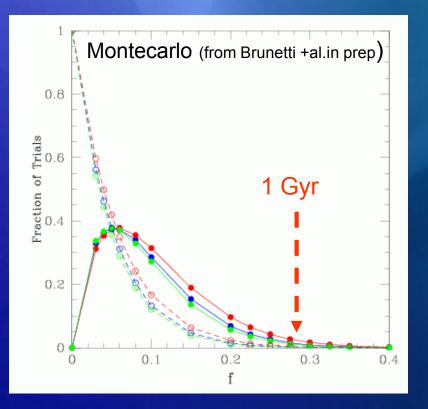


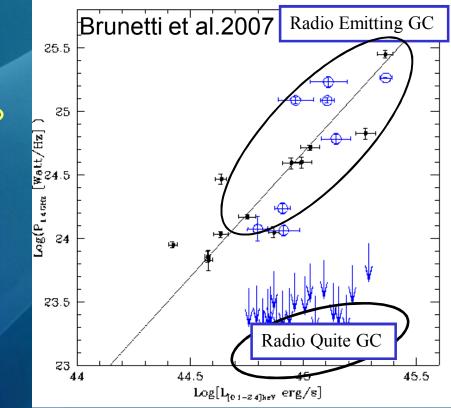
Magnetic field dissipation ?

The complex statistics of Radio

Connection with cluster mergers (e.g. Schuecher et al. 2001, Markevitch et al. 2002, Boschin et al. 2003 Govoni et al. 2004)

Does non thermal emission evolve ? Which time-scale ?



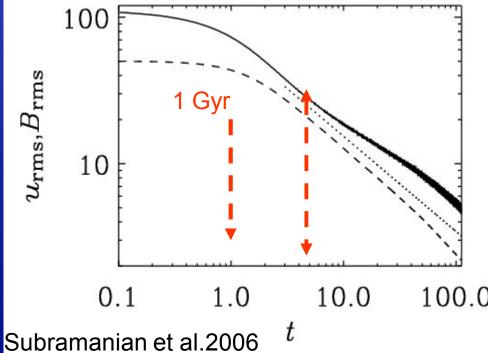


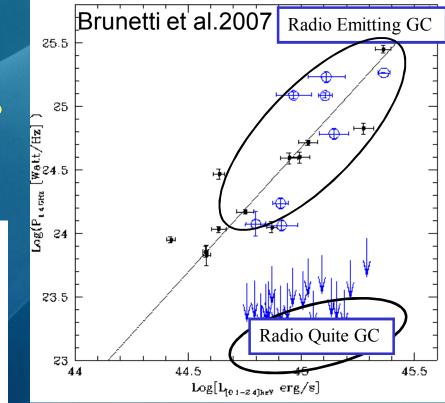
Magnetic field dissipation in the IGM must happen in << 1 Gyr



Connection with cluster mergers (e.g. Schuecher et al. 2001, Markevitch et al. 2002, Boschin et al. 2003 Govoni et al. 2004)

> Does non thermal emission evolve ? Which time-scale ?





Thus the "bi-modality" is due to the evolution (acceleration and cooling) 100.0 of the **emitting electrons**

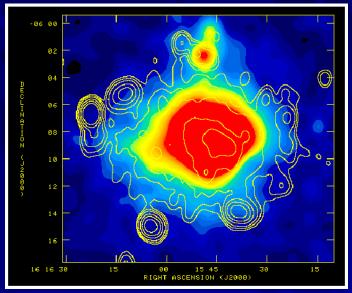
Origin of the emitting electrons

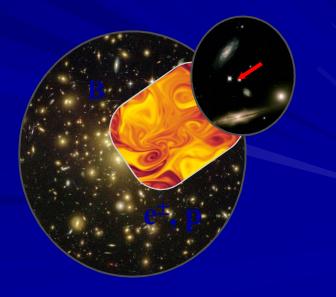
In situ re-acceleration of primary & secondary particles by cluster MHD turbulence (e.g., Brunetti et al. 2001, 2004; Petrosian 2001; Ohno et al. 2002; Fujita et al. 2003; Brunetti & Blasi 2005; Cassano & Brunetti 2005; Brunetti & Lazarian 2007; Petrosian & Bykov 2008)



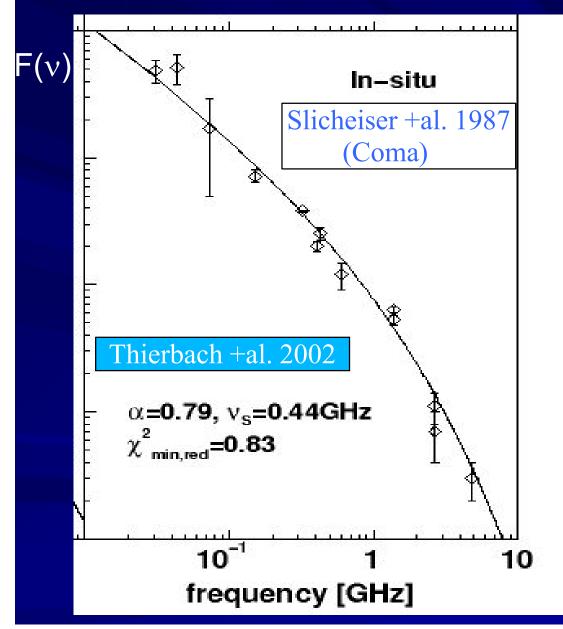
Radio halo – merger connection

Fast dissipation (evolution) of synchrotron emission (CRe cooling)

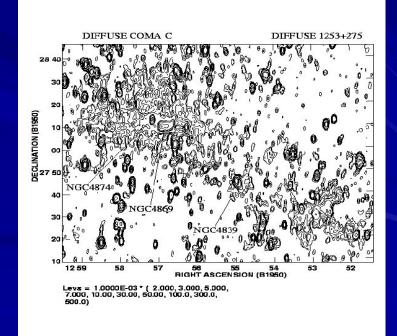




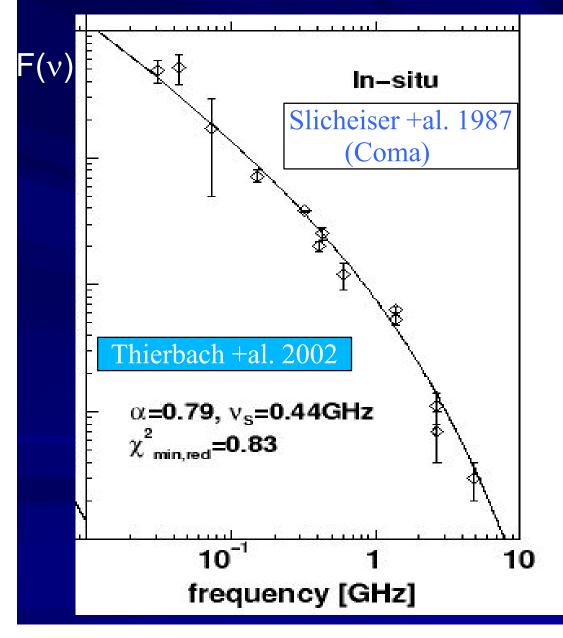
Radio Halos : are they generated by "inefficient" mechanism of CRe acceleration ?



Evidence of break in the synchrotron spectrum

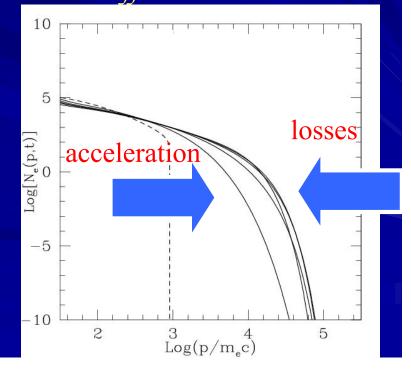


Radio Halos : are they generated by "inefficient" mechanism of CRe acceleration ?

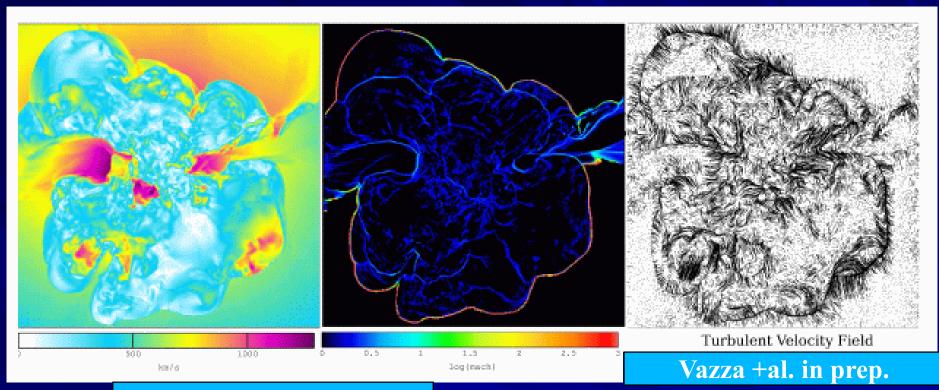


Evidence of break in the spectrum of the emitting electrons at energies of few GeV

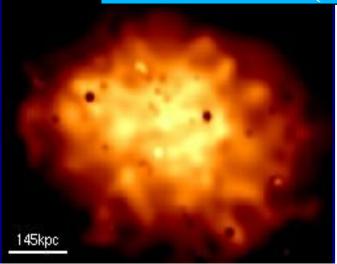
Acceleration mechanism not efficient !

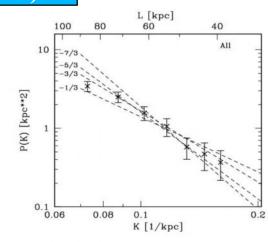


Turbulence in galaxy clusters & cluster-mergers

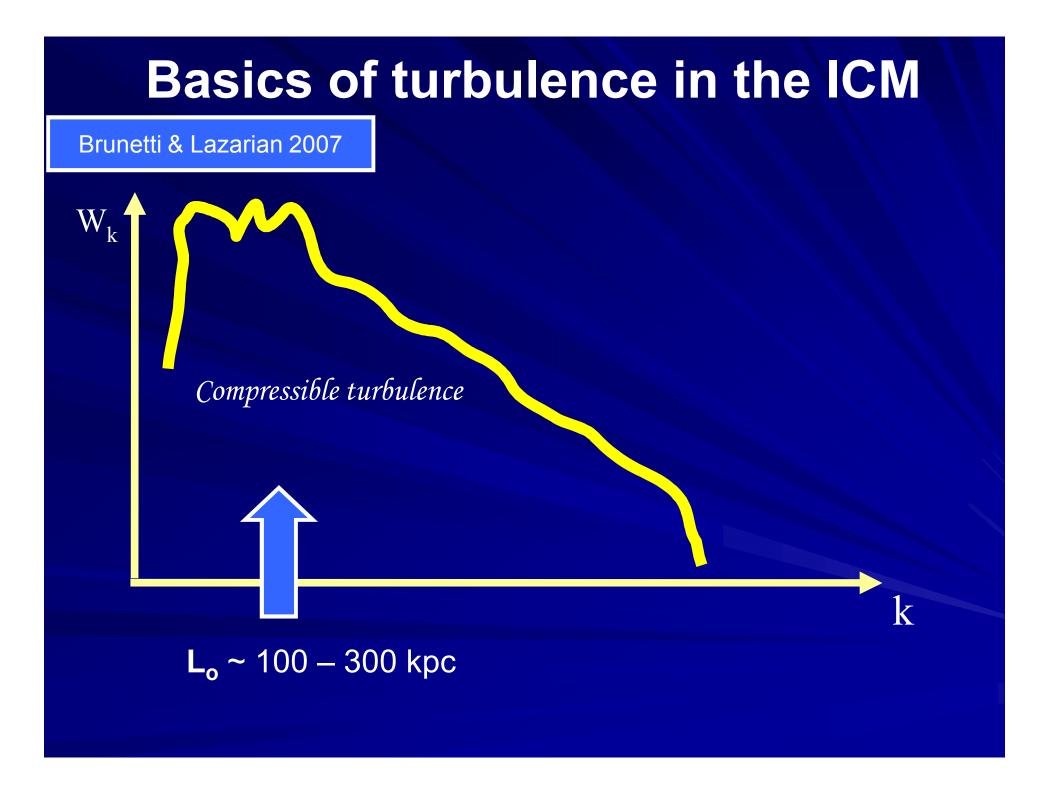


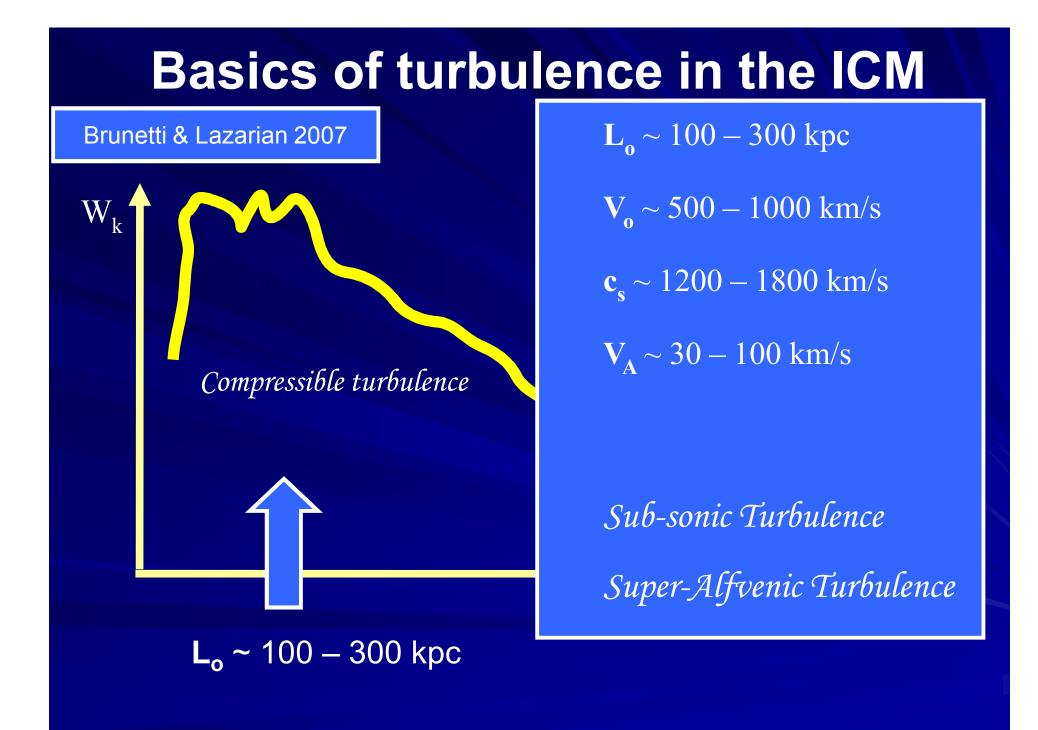
Schuecker +al. (2004)

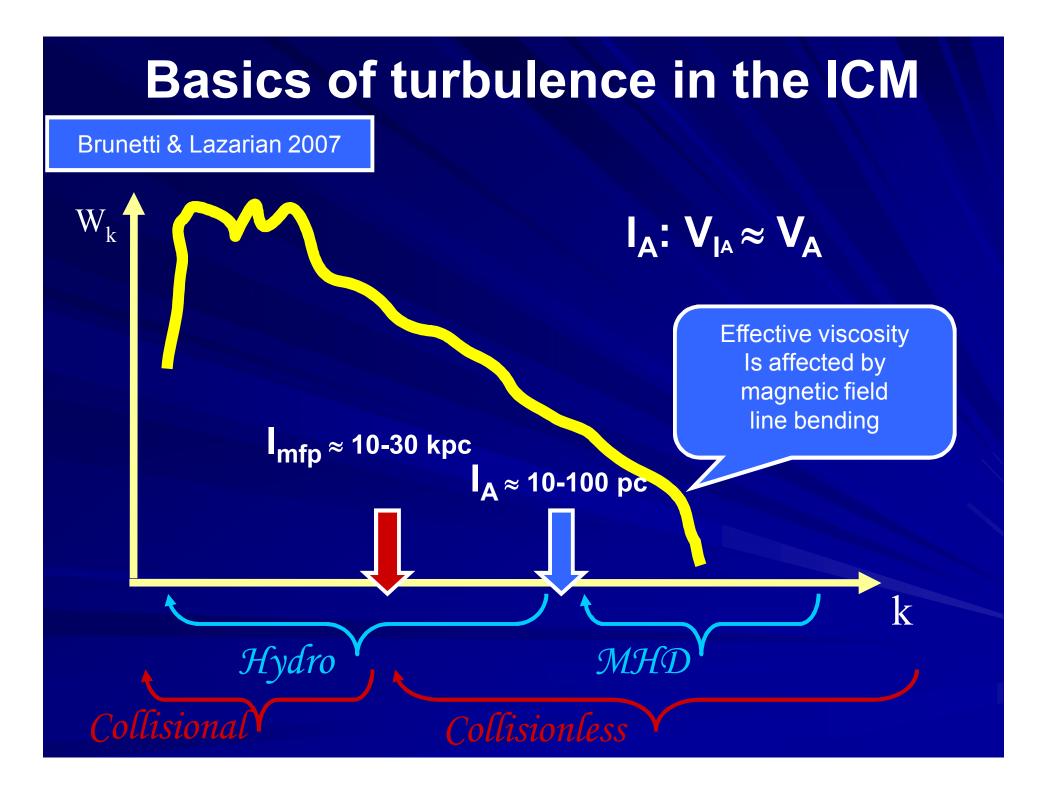


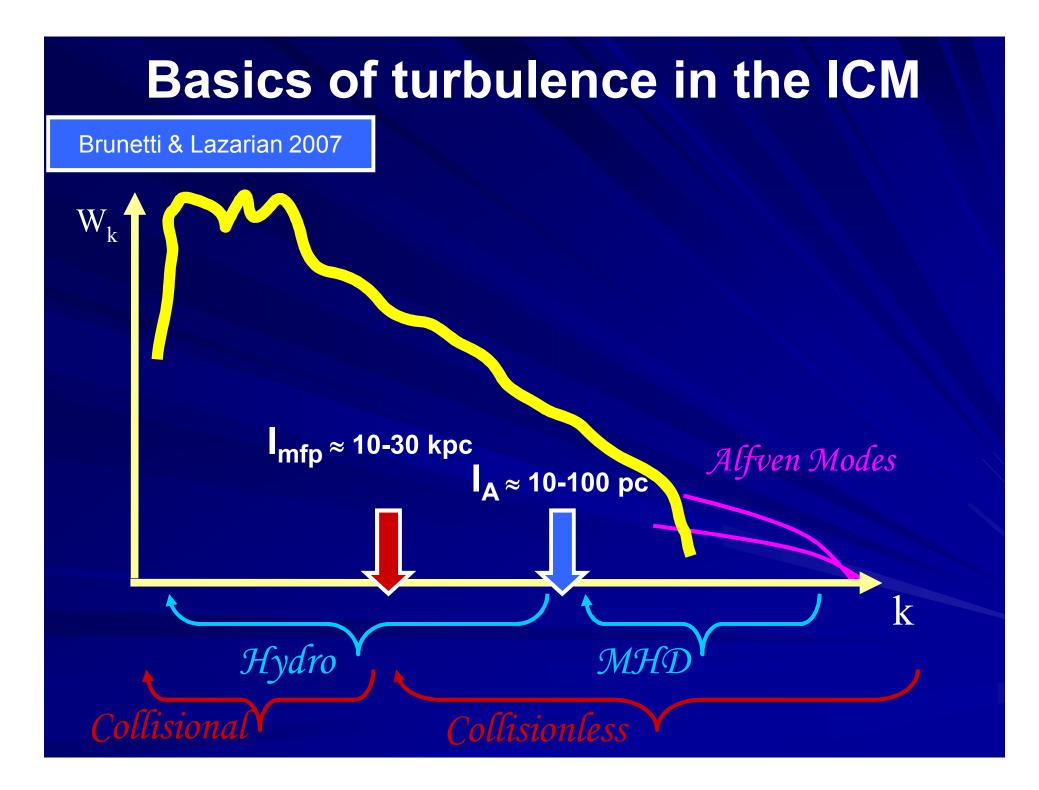


Bryan & Norman 1998; Ricker & Sarazin 2001; Sunyaev et al.2003; Cassano & Brunetti 2005 Dolag et al. 2005; Vazza et al. 2006; Nagai et al. 2007; Brunetti & Lazarian 2007; Iapichino & Niemeyer 2008



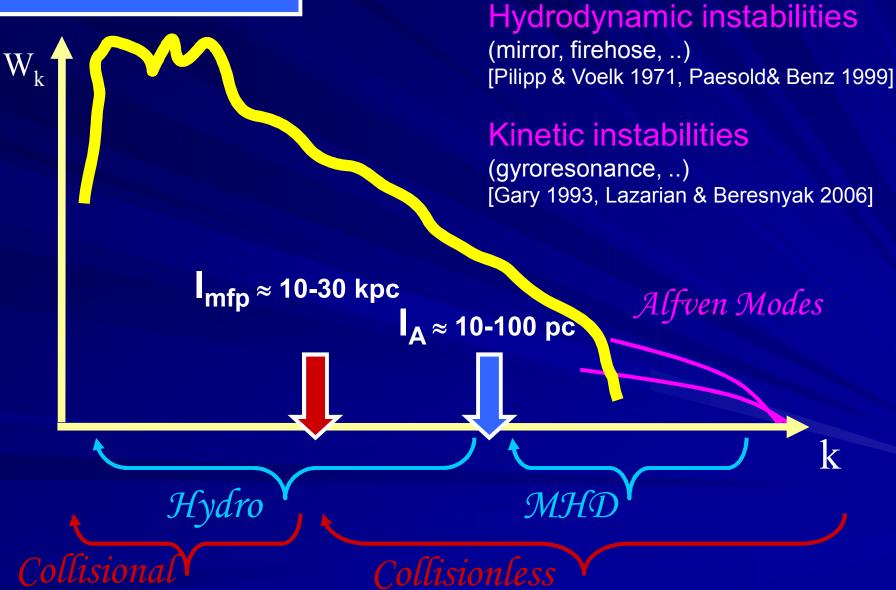


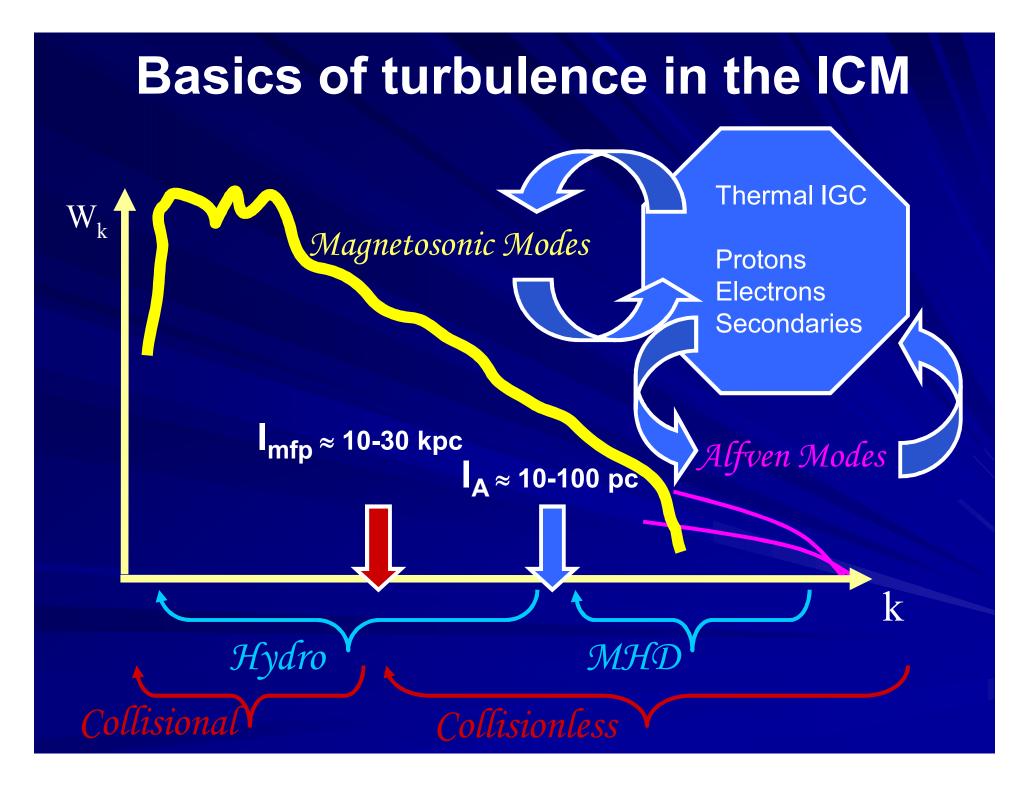




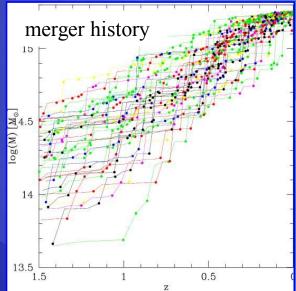


Brunetti & Lazarian 2007

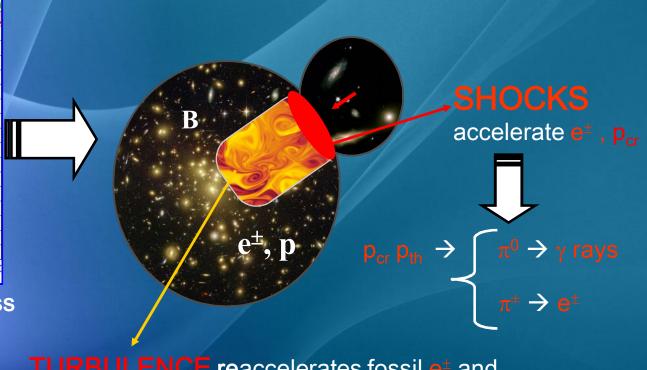




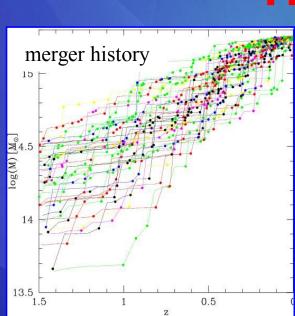
The general picture



clusters increase their mass via merger with smaller subclusters

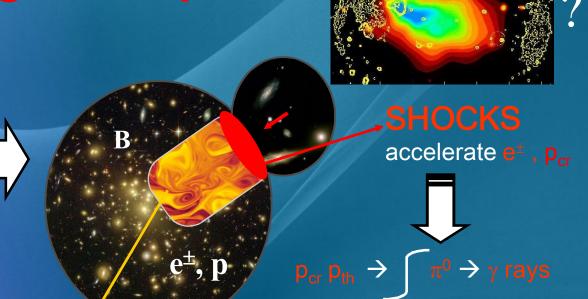


TURBULENCE reaccelerates fossil e[±] and secondaries e[±] on Mpc scales

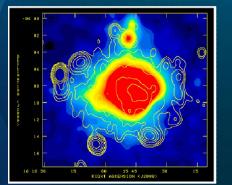


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TURBULENCE reaccelerates fossil e[±] and secondaries e[±] on Mpc scales



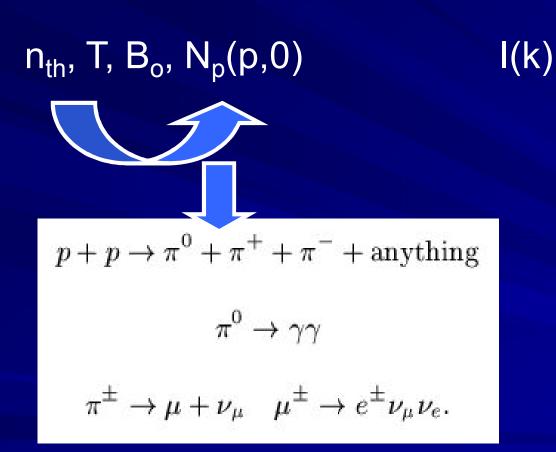
?

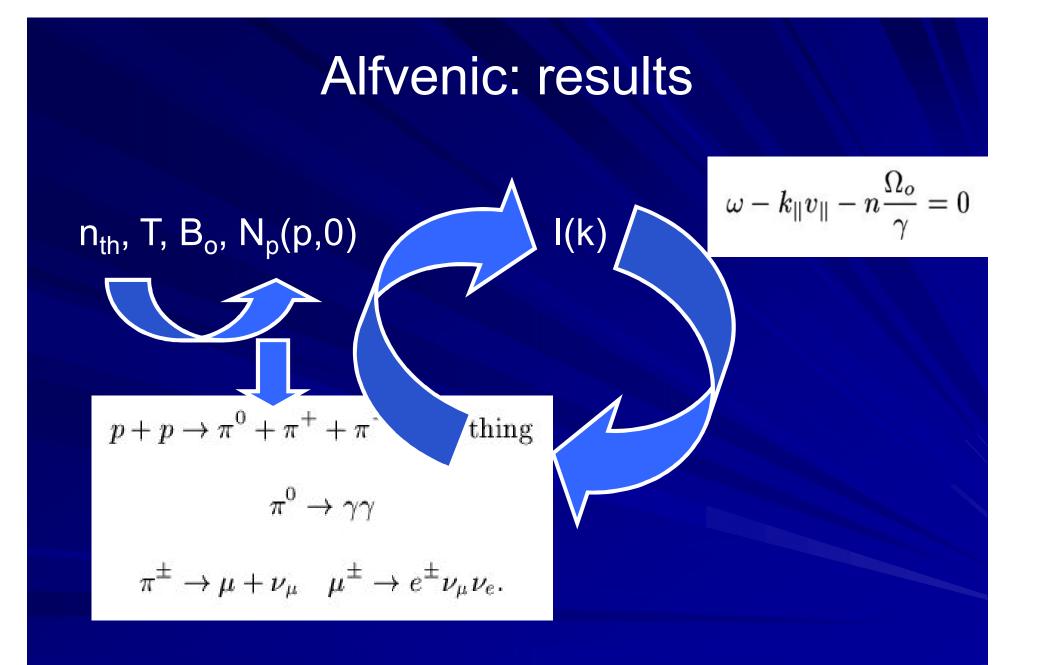
Alfvenic: results

n_{th}, T, B_o, N_p(p,0)

l(k)

Alfvenic: results





Alfvenic: results

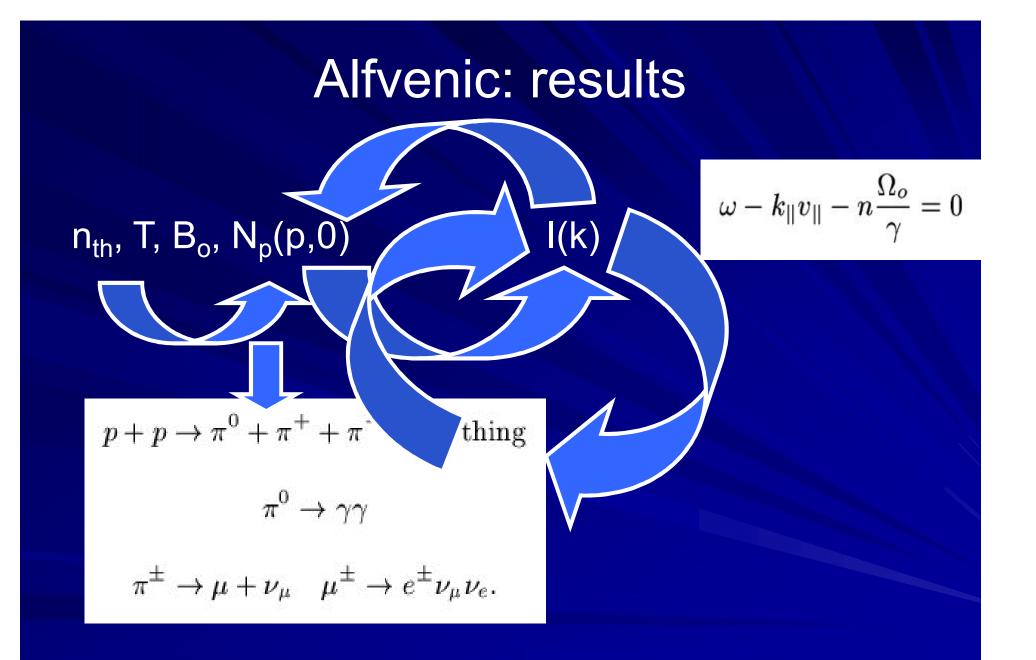
$$n_{th}, T, B_{o}, N_{p}(p, 0) \qquad I(k)$$

$$p + p \rightarrow \pi^{0} + \pi^{+} + \pi^{-} + anything$$

$$\pi^{0} \rightarrow \gamma \gamma$$

$$\pi^{\pm} \rightarrow \mu + \nu_{\mu} \quad \mu^{\pm} \rightarrow e^{\pm} \nu_{\mu} \nu_{e}.$$

$$\omega - k_{\parallel} v_{\parallel} - n rac{\Omega_o}{\gamma} = 0$$



Stochastic Particle Acceleration (formalism)

Brunetti +al. 2004, Brunetti & Blasi 2005, Brunetti +al. 2009

Electrons/Positrons

$$\frac{\partial N_e(p,t)}{\partial t} = \frac{\partial}{\partial p} \left(N_e(p,t) \left[\left(\frac{dp}{dt} \right)_{rad} + \left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_{rad} + \left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_{rad} + \left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right] + \frac{\partial}{\partial p} \left(\frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(\frac{\partial N_e(p,t$$

Protons

 Q_e : secondaries from (CR)p-p collisions

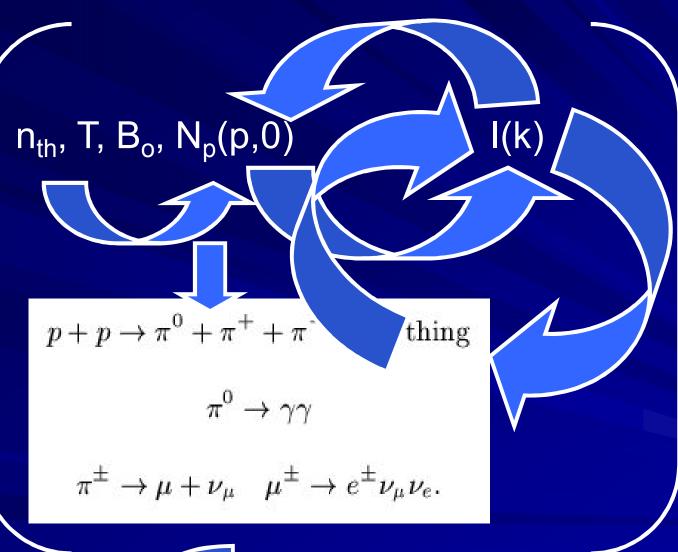
$$\frac{\partial N_p(p,t)}{\partial t} = \frac{\partial}{\partial p} \left(N_p(p,t) \left[\left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_p(p,t)}{\partial p} \right) + Q_p(p,t)$$

Waves

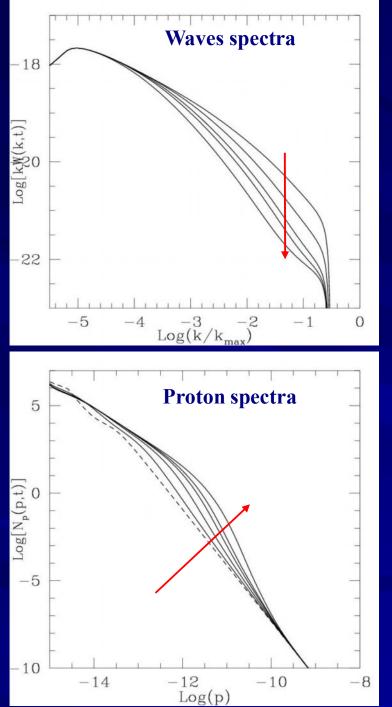
$$\frac{\partial \mathcal{W}(k,t)}{\partial t} = \frac{\partial}{\partial k} \left(k^2 D_{kk} \frac{\partial}{\partial k} \left(\frac{\mathcal{W}(k,t)}{k^2} \right) \right) - \sum_i \Gamma_i(k,t) \mathcal{W}(k,t) + I(k,t)$$

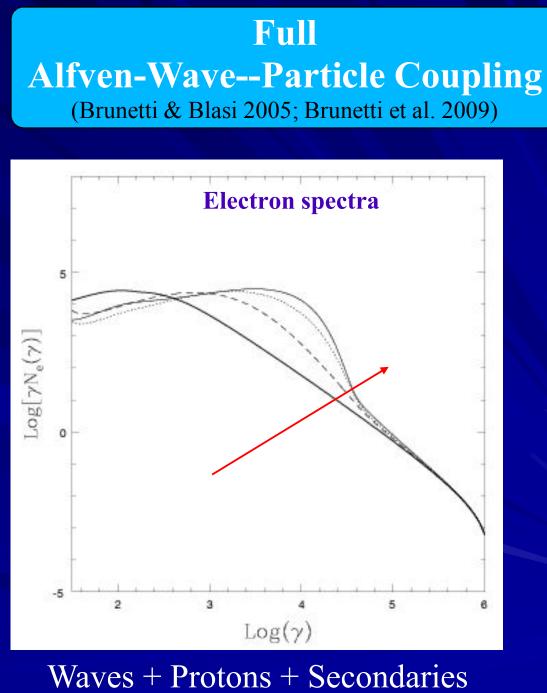
most could be with CR protons

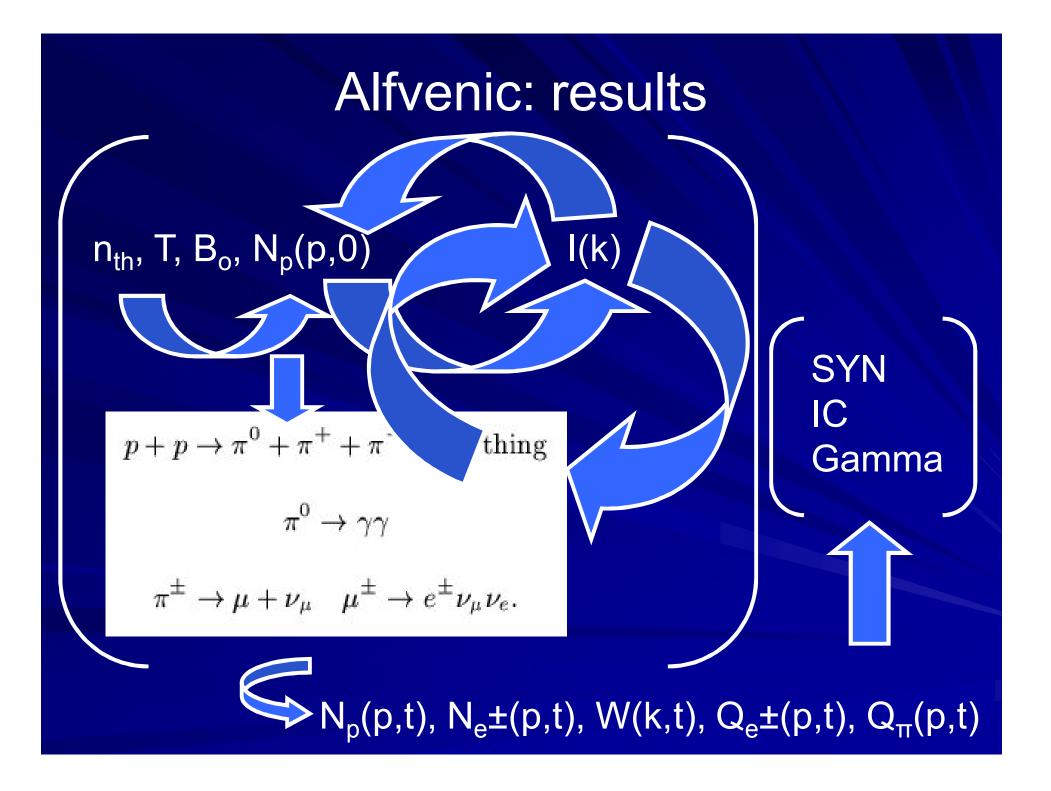




 $\gg N_p(p,t), N_e \pm (p,t), W(k,t), Q_e \pm (p,t), Q_{\pi}(p,t)$

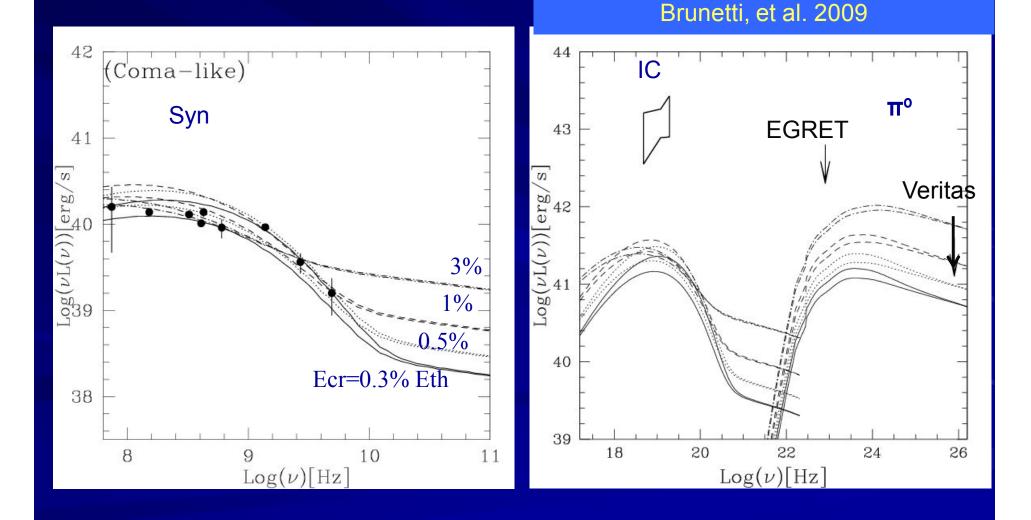






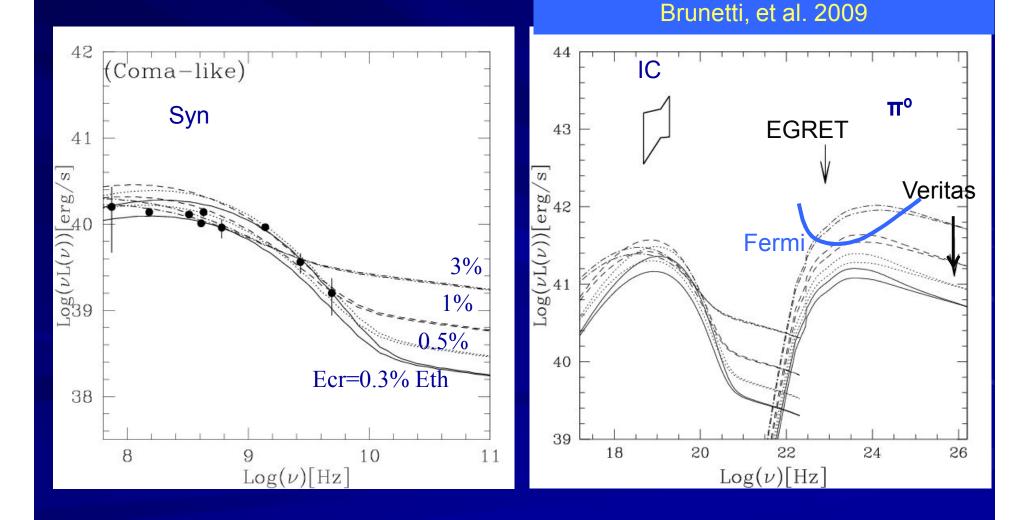
Alfvenic: results

Toy Model: β -profile, $B_o \approx A n_{th}$, $B_o(0)=2\mu G$, $W_{CR} \approx f W_{th}$, $P_A \approx Q n_{th}^{5/6}$



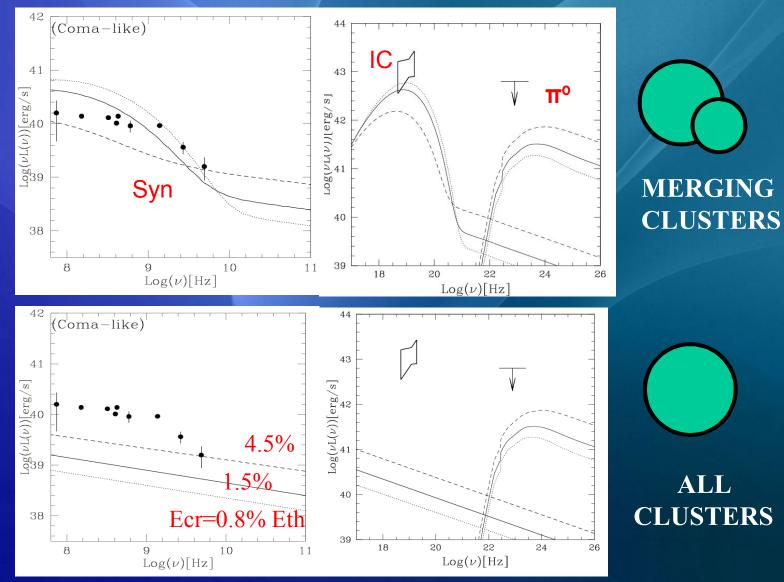
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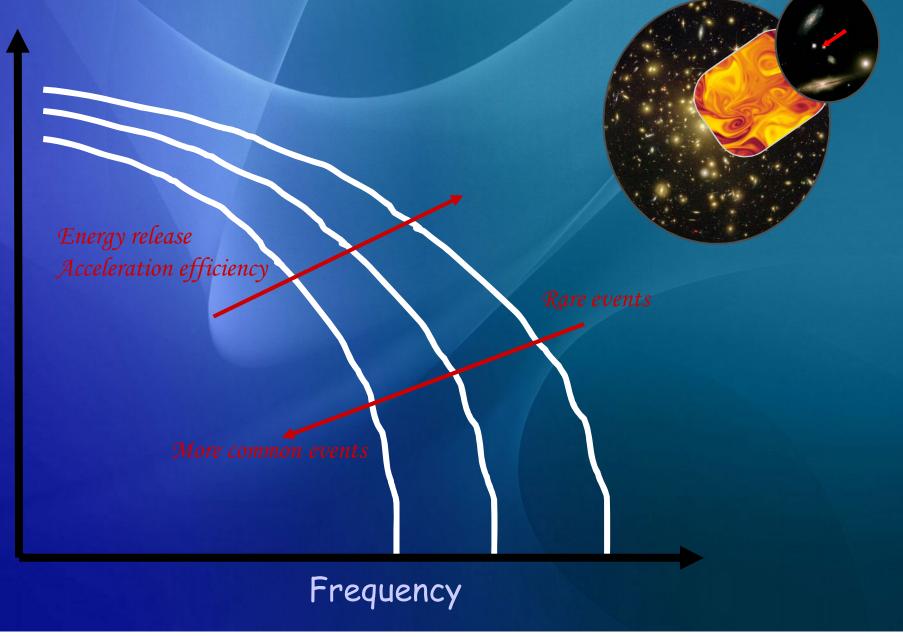


General multiwavelength expectations

Brunetti 2008; Brunetti +al. 2009



The case of the "ultra steep" spectrm radio halos



The case of the "ultra steep" spectrm radio halos







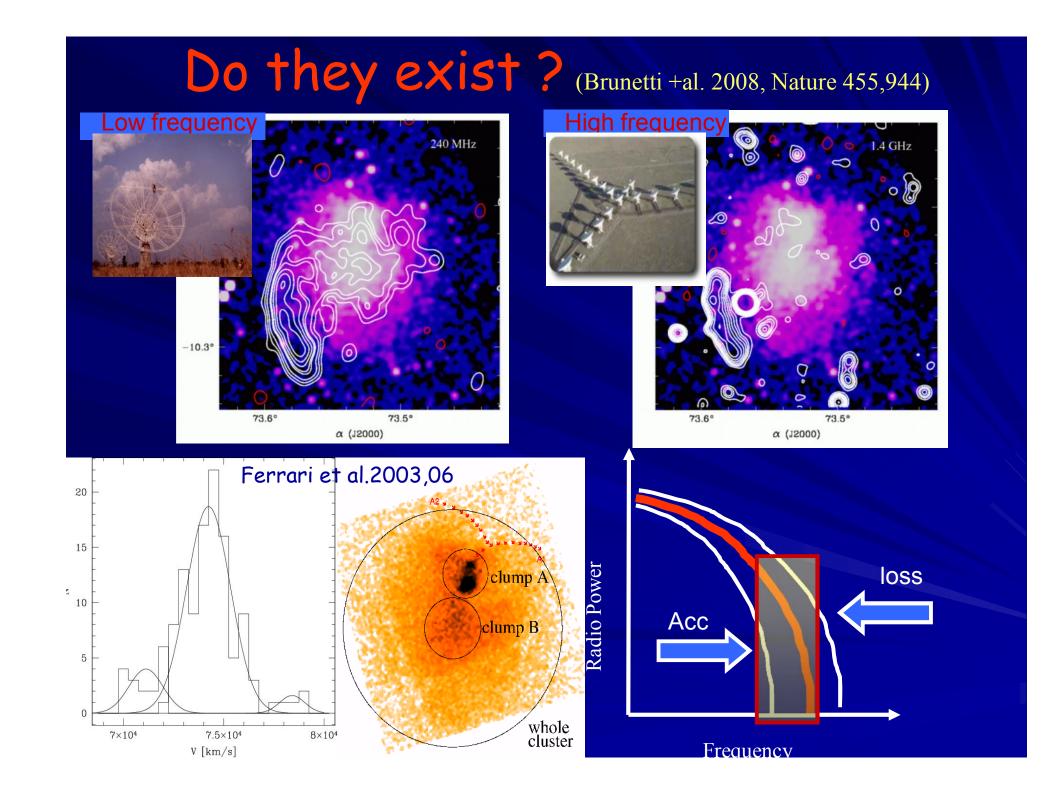
Frequency

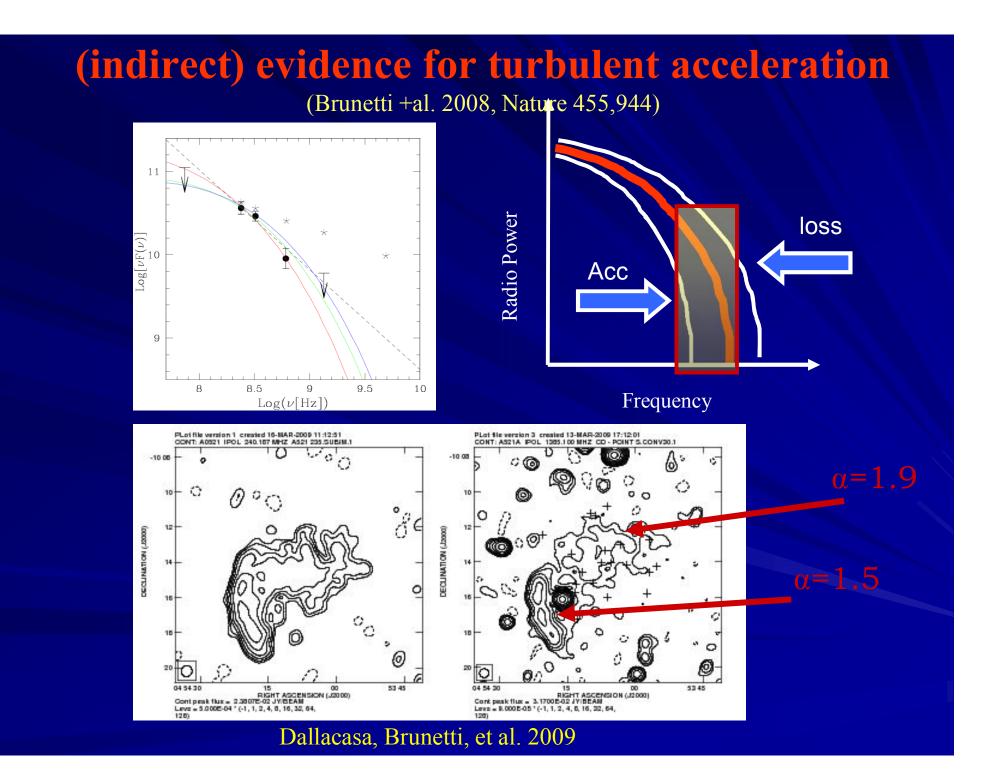
The case of the "ultra steep" spectrm radio halos

LOFAR is expected to detect 10-100 times more radio halos that present VLA surveys .

We expect a populations of radio halos with steeper spectrum that is (better) visible at low frequencies (Cassano, Brunetti, Setti 2006)

Frequency





Conclusions

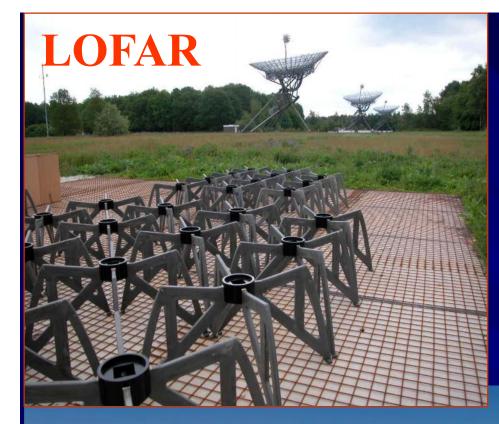
Theoretically <u>protons should be the dominant</u> non thermal particle-component in the IGM (yet only limits)

Theoretically <u>shocks may be an important</u> source of non thermal particles in the IGM, in addition present data suggest that <u>turbulence plays a role in the acceleration</u> of particles in galaxy clusters

The non-thermal emission from galaxy clusters is a mixture of a transient (turbulnce) and long-living component

LOFAR & FERMI will step into an essentially unexplored territory where theories can be tested (eg "ultra steep halos")





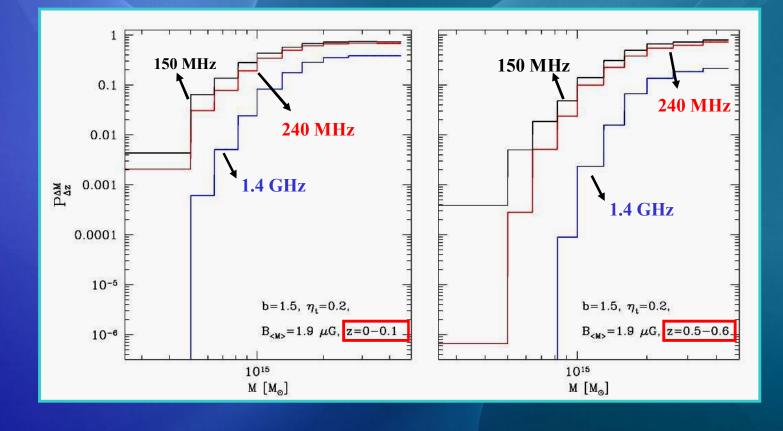
Cherenkov

The FUTURE

Fermi

Fraction of galaxy clusters with radio halos at low v

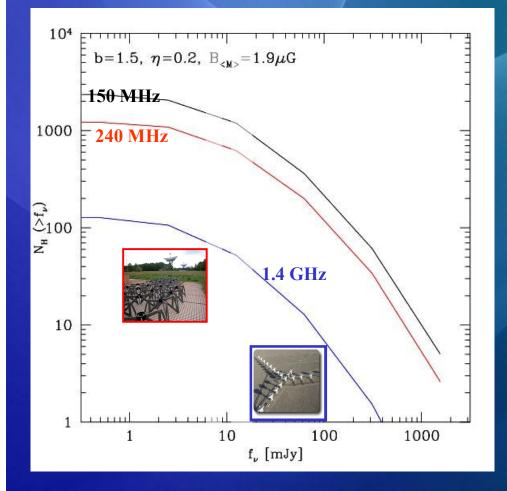
Cassano et al. 2008

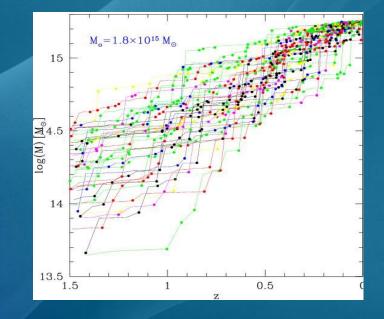


 \clubsuit The expected fraction of clusters with radio halos increases at low **v**

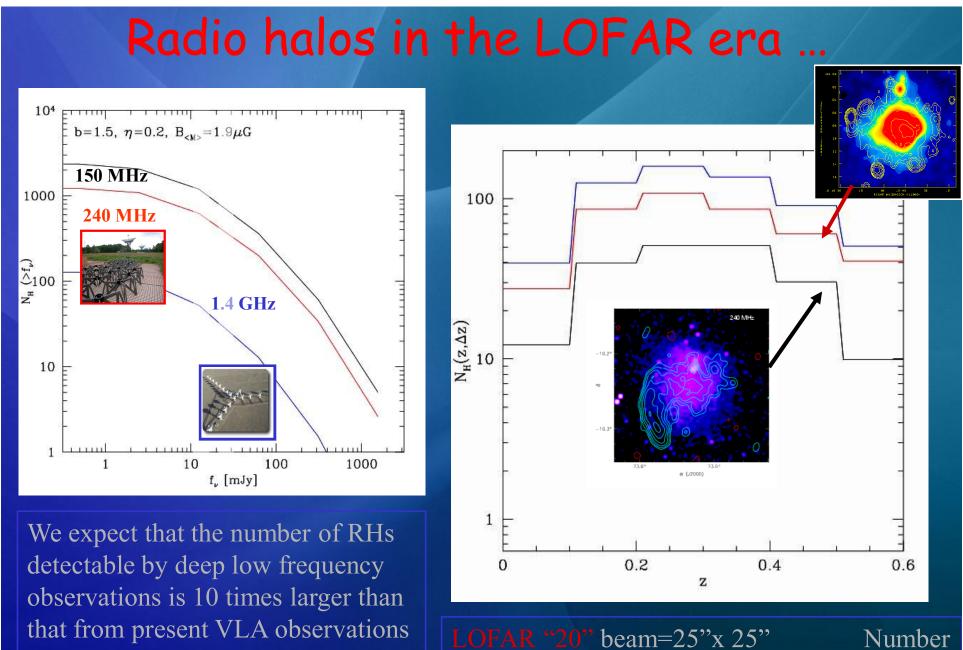
♦ This increase is even stronger for smaller clusters (M<10¹⁵ M_{\odot})

Radio halos in the LOFAR era ...



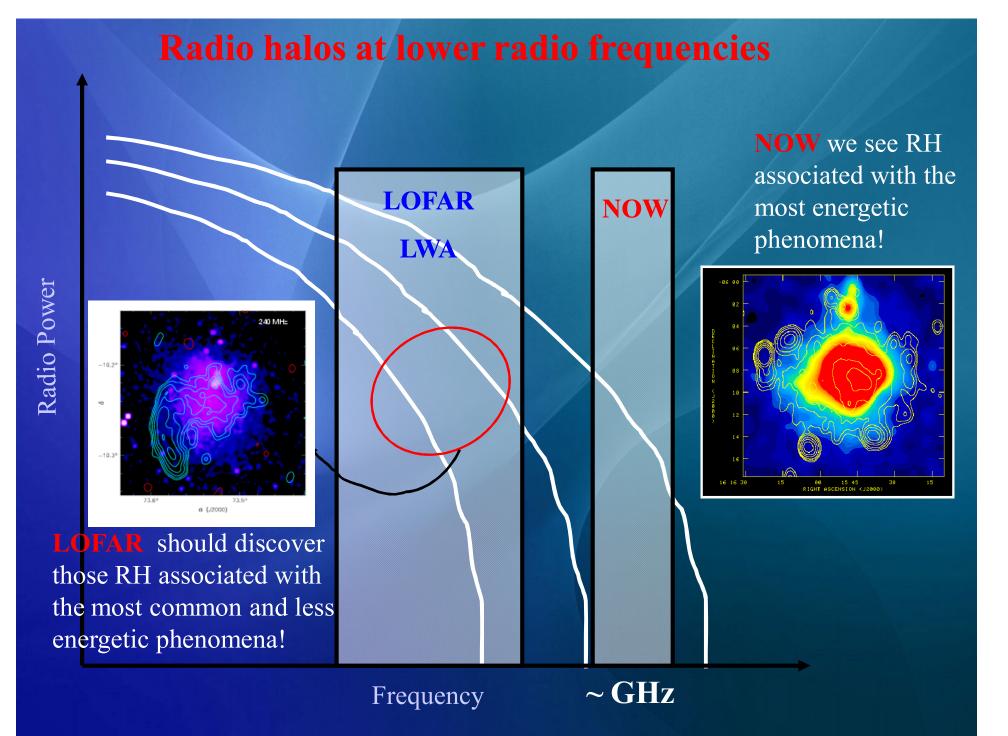


The number of radio halos detectable by deep low frequency observations is expected 10 times larger than that from present VLA observations at 1.4 GHz : LOFAR and LWA will catch these sources!

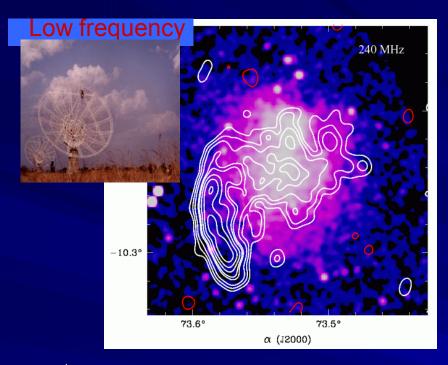


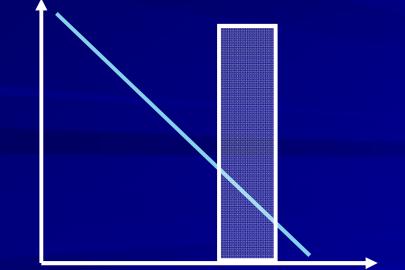
LOFAR "20" beam=25"x 25" counts at 120 MHz (Cassano et al., in prep.).

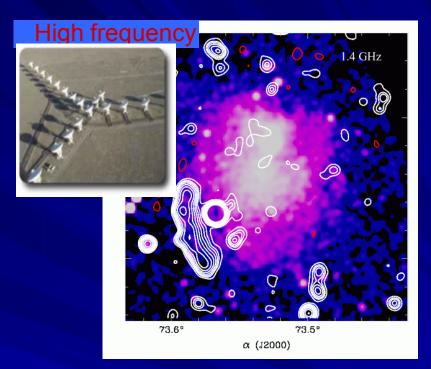
that from present VLA observations at 1.4 GHz : LOFAR and LWA would catch these RHs !



Do they exist? (Brunetti +al. 2008, Nature 455,944)



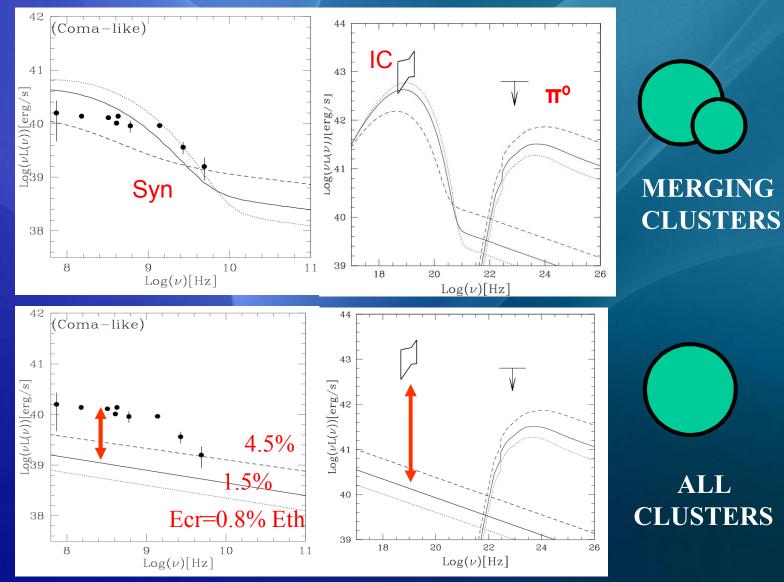




Supposing that the observed synchrotron radiation is from secondary electrons generated during p-pcolisions in theIGM .. we arrive at the untenable scenario of a relativistic proton population whose energy density exceeds that of the thermal IGM

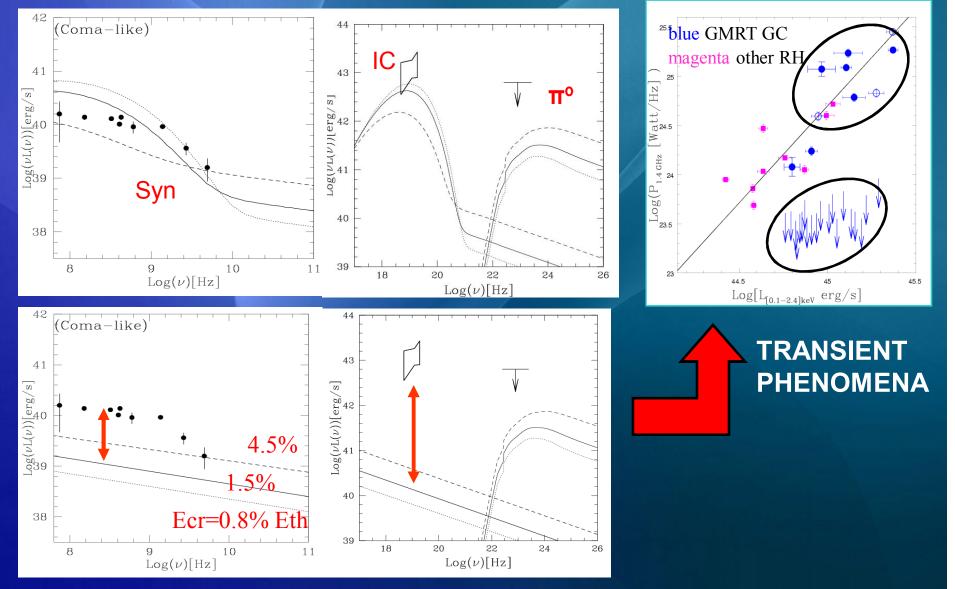
General multiwavelength expectations

Brunetti 2008; Brunetti +al. 2009



General multiwavelength expectations

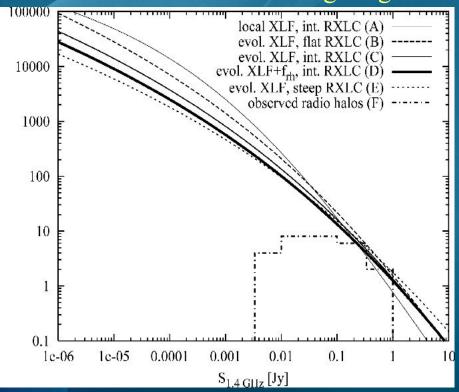
Brunetti 2008; Brunetti +al. 2009



Why low frequency ?

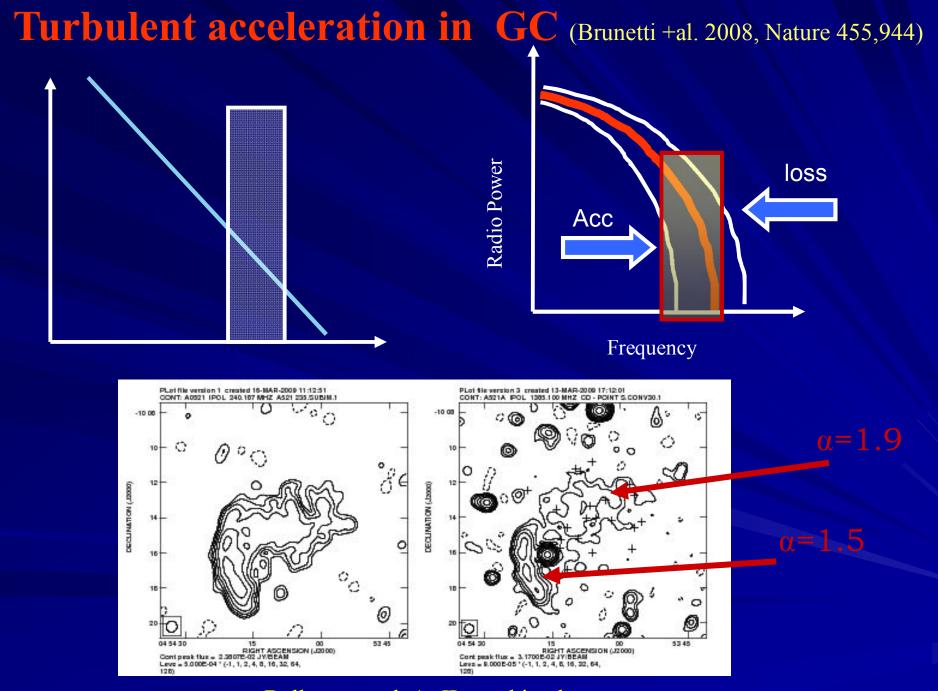
Ensslin & Roettgering 2002

Regardless of the origin of Radio Halos, extrapolations of their number counts at 1.4 GHz based on the Radio-X ray correlation observed for Radio Halos suggest that a large fraction of these Halos is at faint fluxes.



Due to their steep synchrotron spectrum, faint Radio Halos should appear more luminous at low frequencies and thus LOFAR and LWA are expected to discover a large number of these objects.

dlog10S1.4 GHz



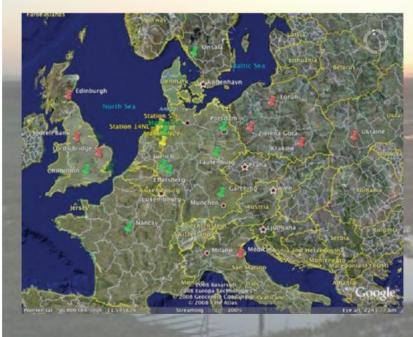
Dallacasa + al. ApJL sumbitted



AST(RON



The LOFAR Array



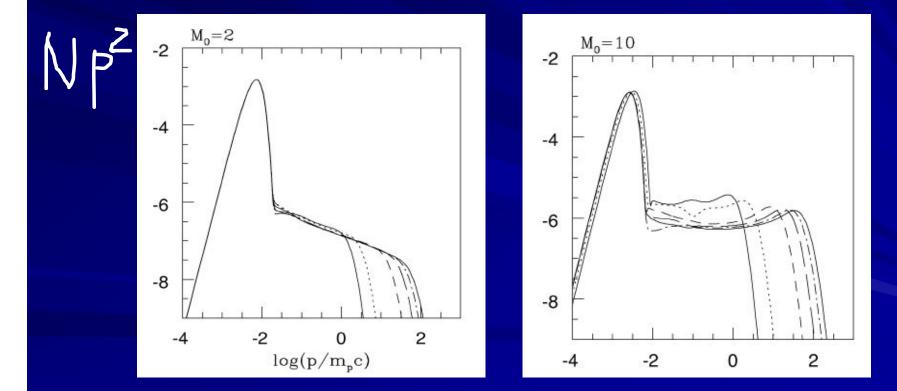
- Core area 2 km around Exloo
 - Min. 18 stations
- Remote stations -100 km
 - Min. 18 stations
- International stations
 Min. 8 stations

Acceleration of CRp at shocks

 $N(p) \sim p^{-d}$

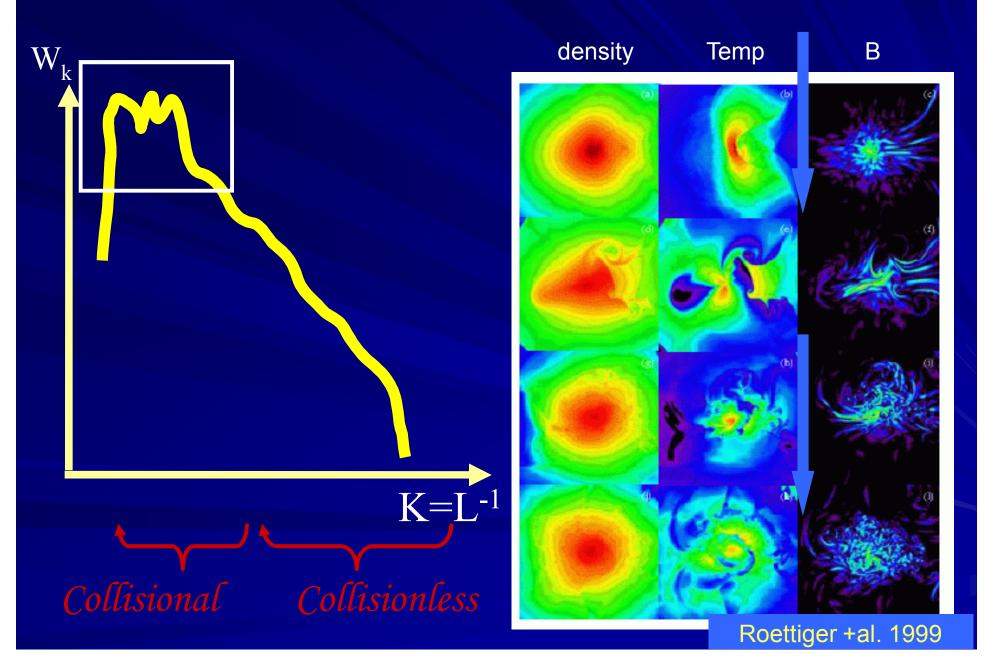
$$\delta = 2\frac{\mathcal{M}^2 + 1}{\mathcal{M}^2 - 1}.$$

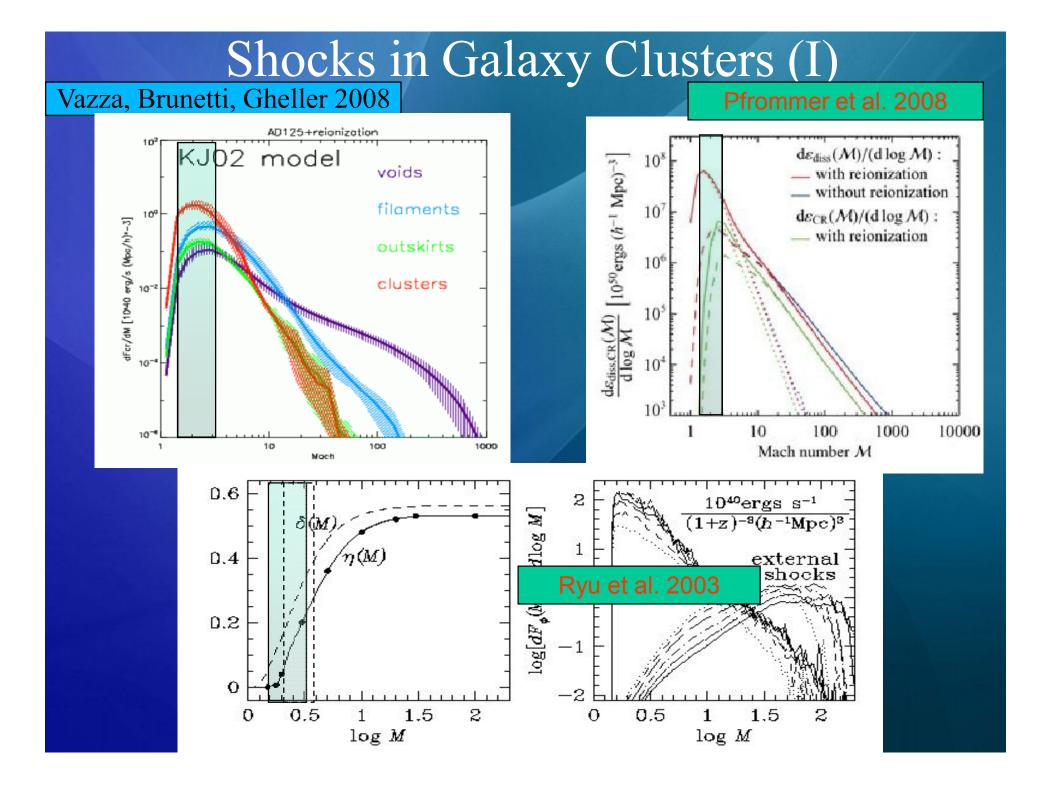
Linear Theory (e.g.,Blandford & Eichler 1987)



(Kang & Jones 2002; Malkov 1997, Kang & Jones 2005, 07, Amato & Blasi 2006)

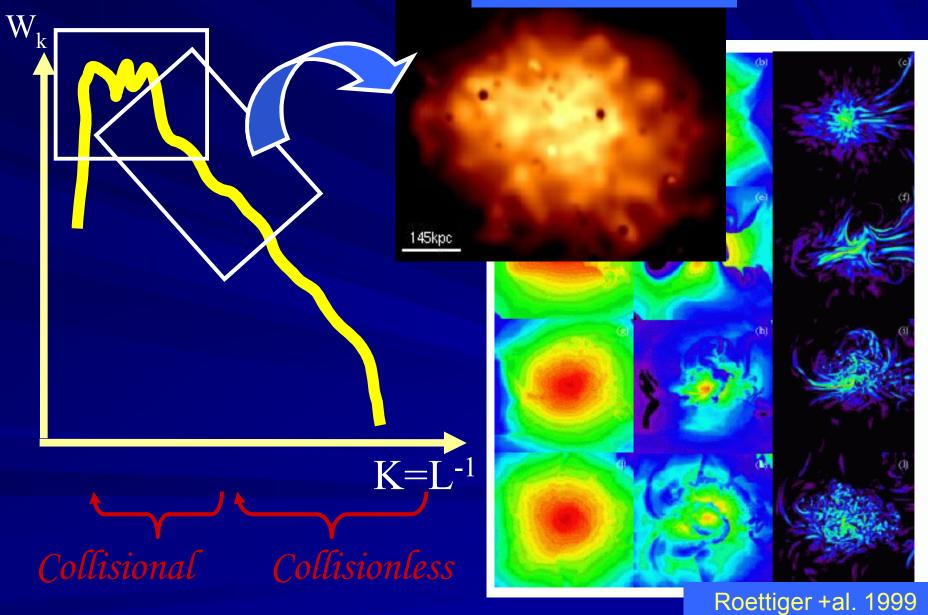
Turbulence in the ICM



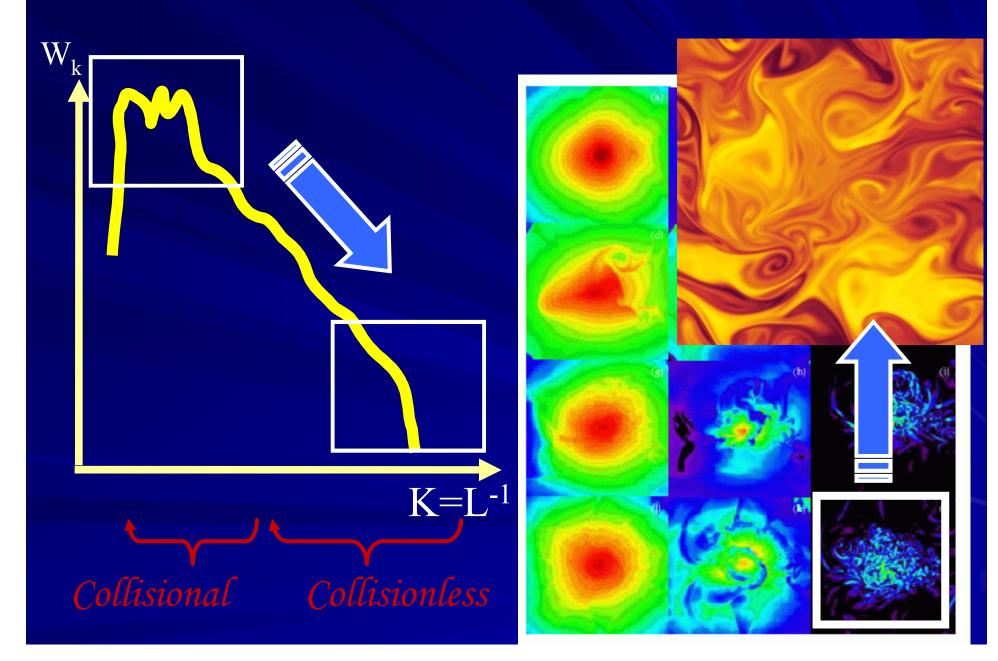


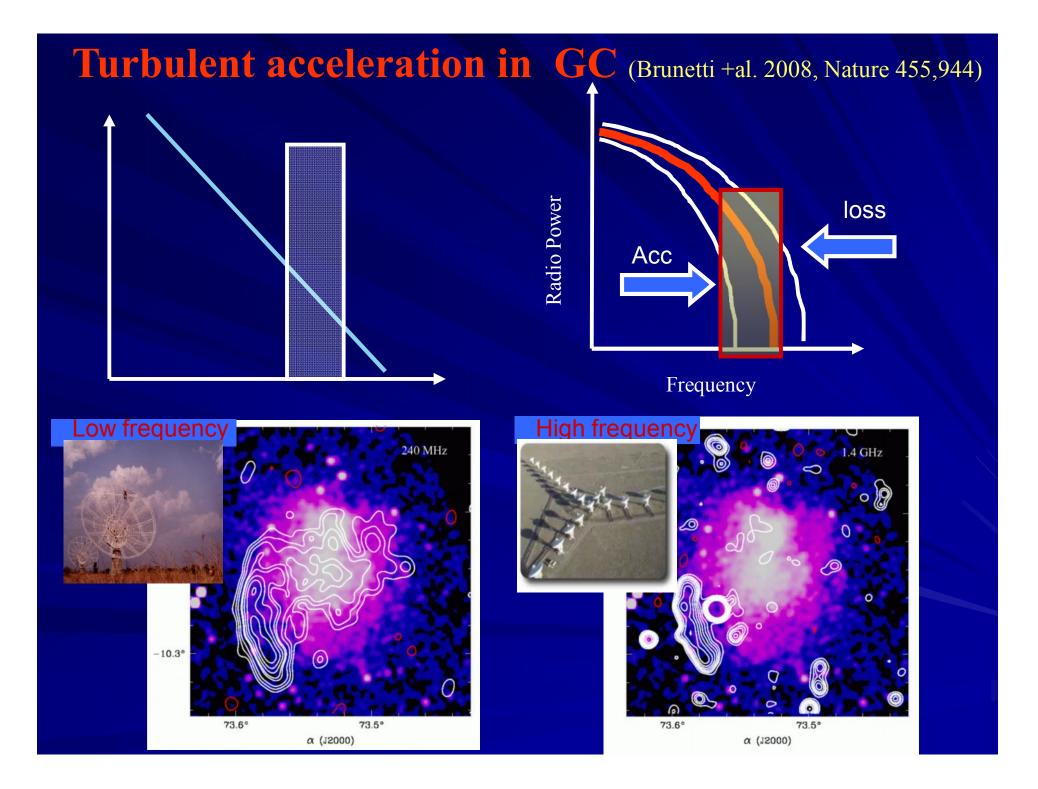
Turbulence in the ICM

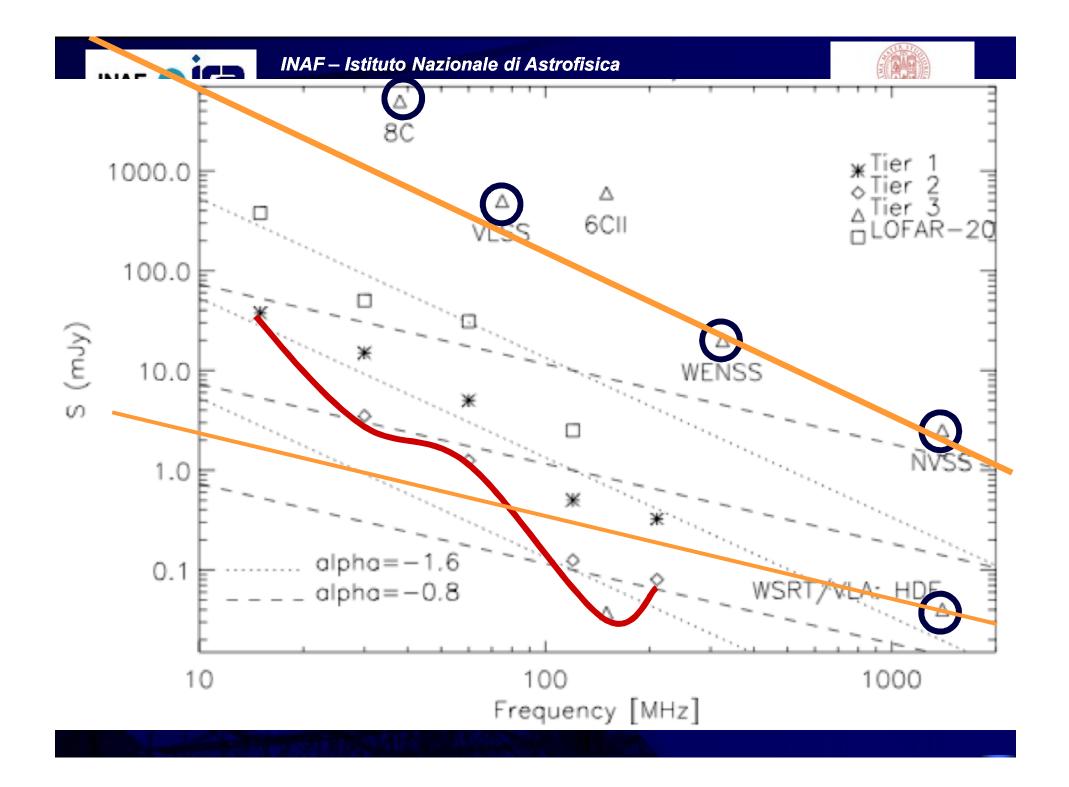
Schueker +al. 2004



Turbulence in the ICM

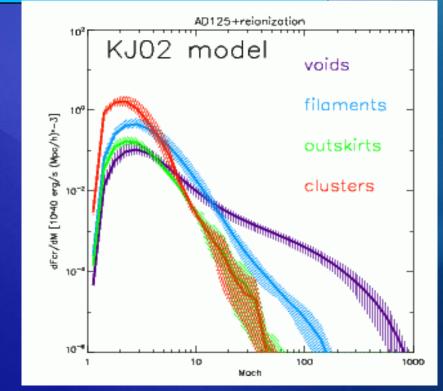


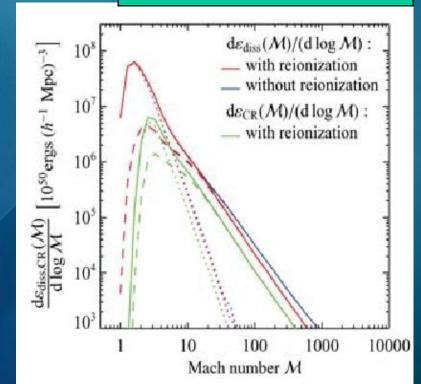




Shocks in Galaxy Clusters

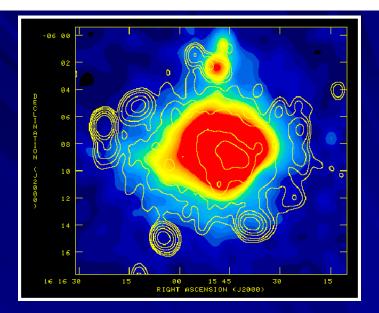
Vazza, Brunetti, Gheller 2008





Pfrommer et al. 2008

Origin of the emitting electrons

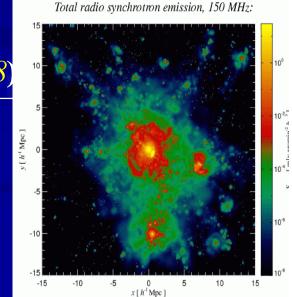


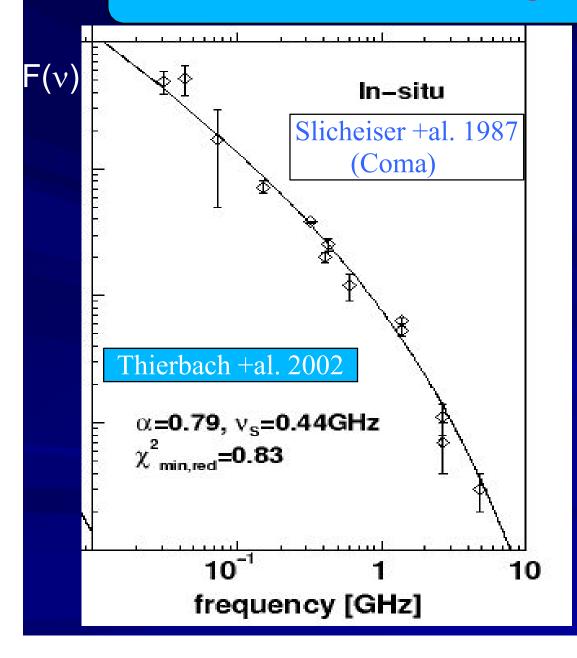
First possibility: *secondary models*, relativistic electrons continuously injected in the ICM by inelastic proton-proton collisions

through productions and decay of charged pions (e.g., Dennison 1980, Blasi & Colafrancesco 1999, Dolag & Ensslin 2000; Pfrommer & Ensslin 2004,08)

Since protons accumulate in galaxy clusters and clusters are magnetised at muG level, we expect that the synchrotron emission from secondary electrons in the IGM should be common.

- Which is the level of this emission ?
- Are Radio Halos due to secondary emission?





We expect a break in the spectrum of the emitting electrons at energies of few GeV Acceleration mechanism not efficient ! 10 $\tau_{Acc}^{II} \approx \frac{L}{c} \left(\frac{c}{\delta v}\right)^2$ 5 losses $Log[N_{e}(p,t)]$ acceleration

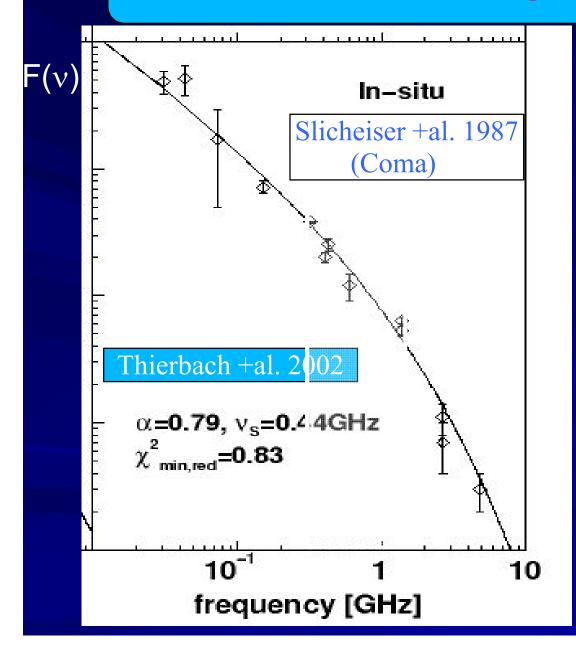
3 4Log(p/m_c)

5

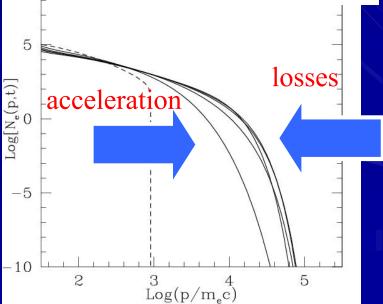
-5

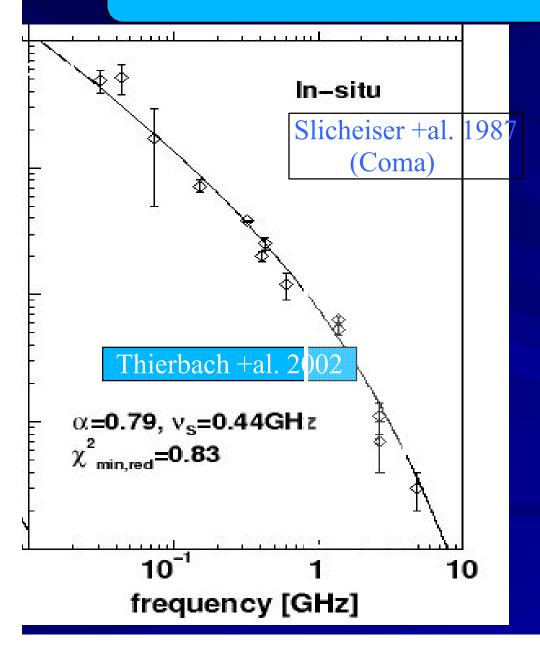
-10

2



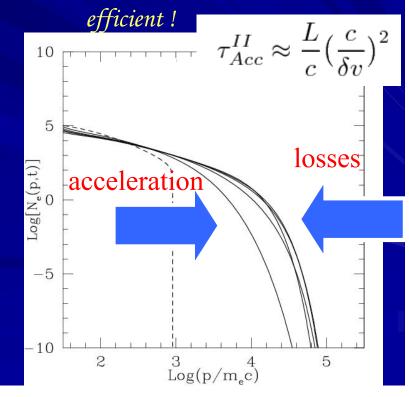
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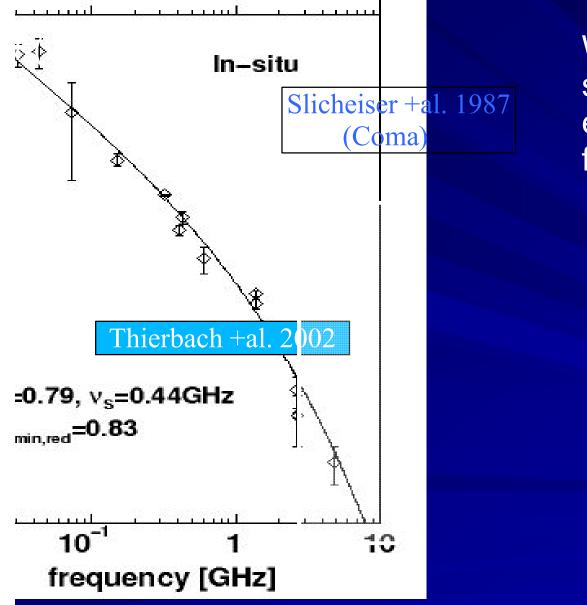




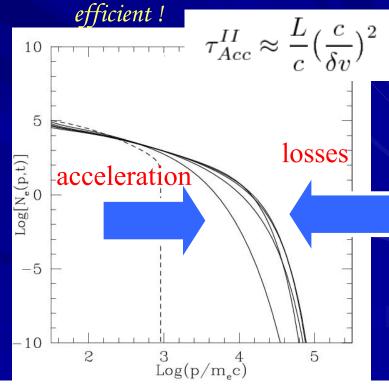
We expect a break in the spectrum of the emitting electrons at energies of few GeV

Acceleration mechanism not





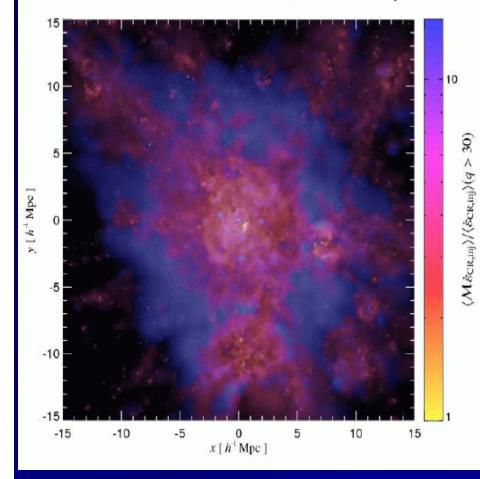
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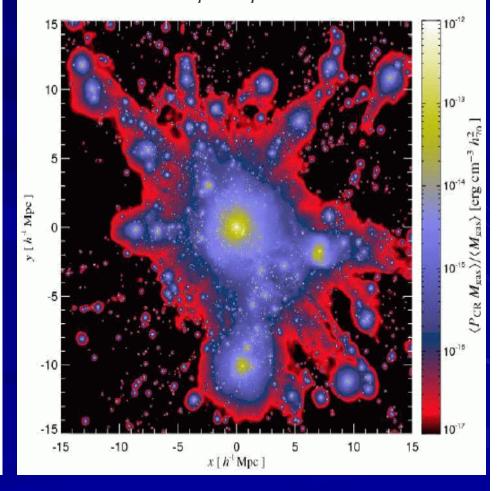
Acceleration of CR at shocks

Pfrommer et al. 2007, 08

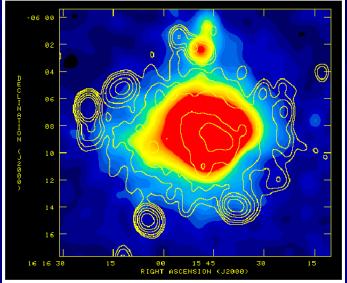
Shock Mach numbers weighted by $\dot{\epsilon}_{CR,inj}$:



CR proton pressure:



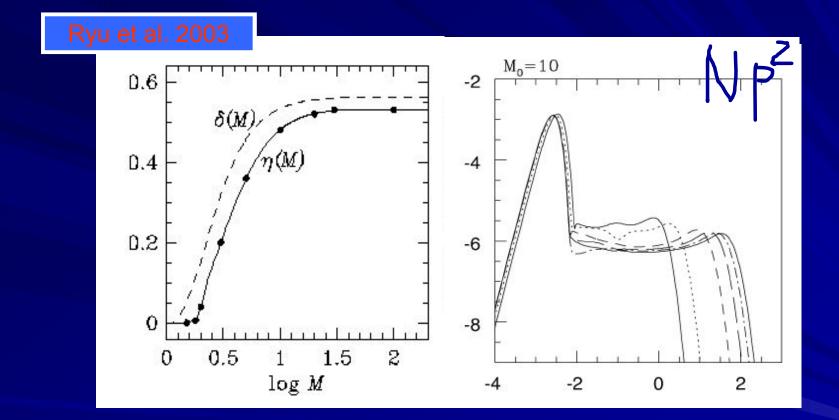
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First possibility: *secondary models*, relativistic electrons continuously injected in the ICM by inelastic proton-proton collisions through productions and decay of charged pions (e.g., *Dennison 1980, Blasi & Colafrancesco 1999, Dolag & Ensslin 2000; Pfrommer & Ensslin 2004*)

Second possibility : *in situ* re-acceleration by MHD turbulence developed in the cluster volume during the merger events (e.g., Brunetti et al. 2001, 2004; Petrosian 2001; Ohno et al. 2002; Fujita et al. 2003; Brunetti & Blasi 2005; Cassano & Brunetti 2005; Brunetti & Lazarian 2007; Petrosian & Bykov 2008)

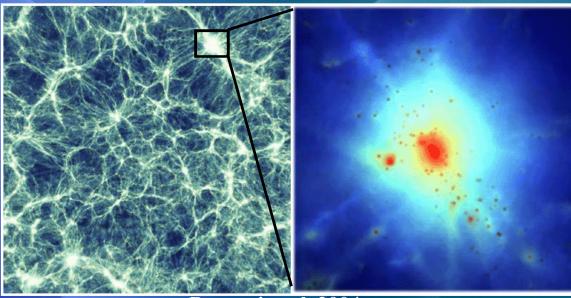
Acceleration of CRp at shocks



 $E_{CR} \approx \eta(M) \rho V_{sh}^3$

From SN and our Galaxy electrons take 1-10% of the proton energy.

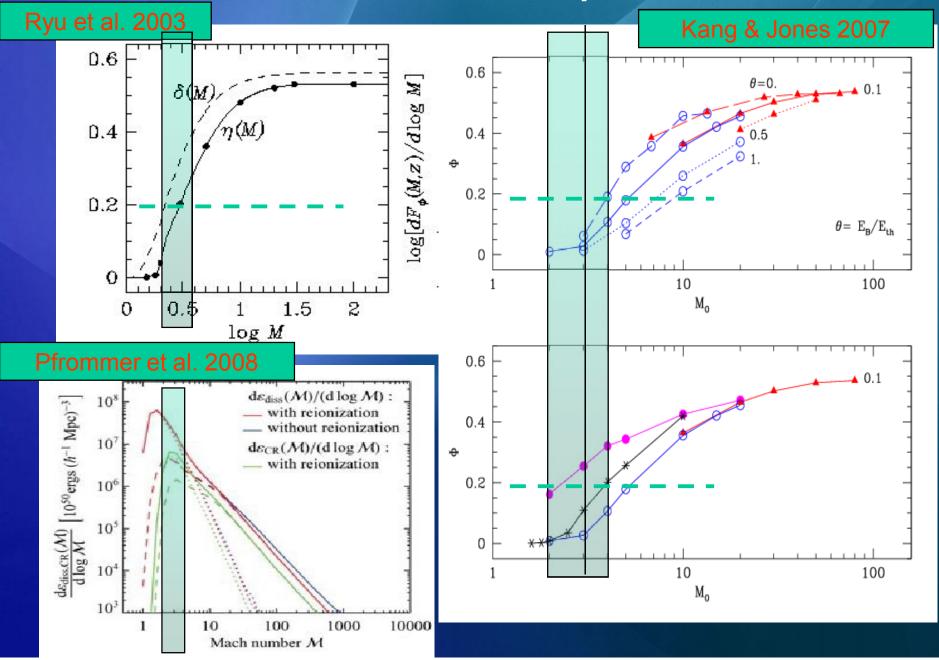
Clusters of galaxies How galaxy clusters form?



Borgani et al. 2004

- Cosmic structures forms as a result of the gravitational amplification of primordial density fluctuations.
- Clusters form by accretion of matter and mergers between sub-units at the intersection of filaments which make up the 'cosmic web'.
- Gas falling into deep potential wells is heated to ~10⁷-10⁸ [°]K by shocks (and adiabatic compression)

Acceleration of CRp at shocks



II - Conclusion

 (a) Protons are expected the dominant CR population in the IGM : they result from the accumulation of CR from the epoch of formation of galaxy clusters

(b) Despite (a) we have only upper limits to the energy content in the form of CR protons

Not only particle acceleration

Thermal conduction and kin. viscosity in the ICM (e.g., Lazarian 2006)

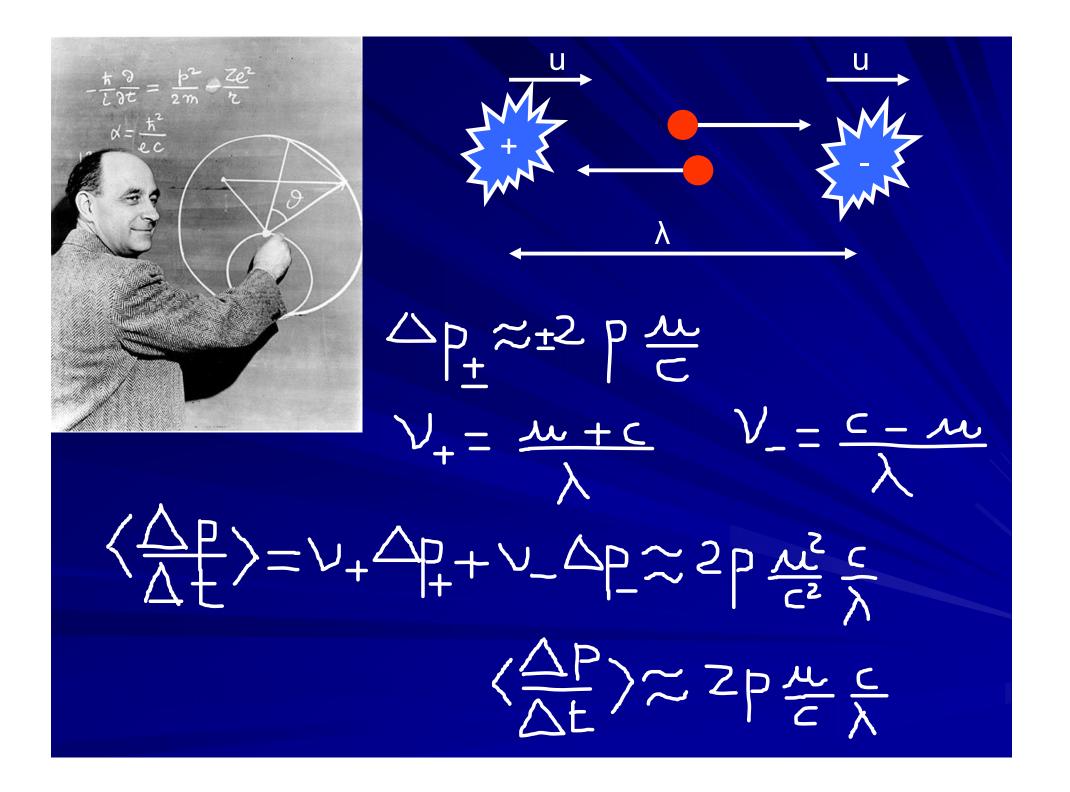
Heating of the ICM and "cooling flow" problem (e.g., Fujita , Matsumoto, Weda 2004)

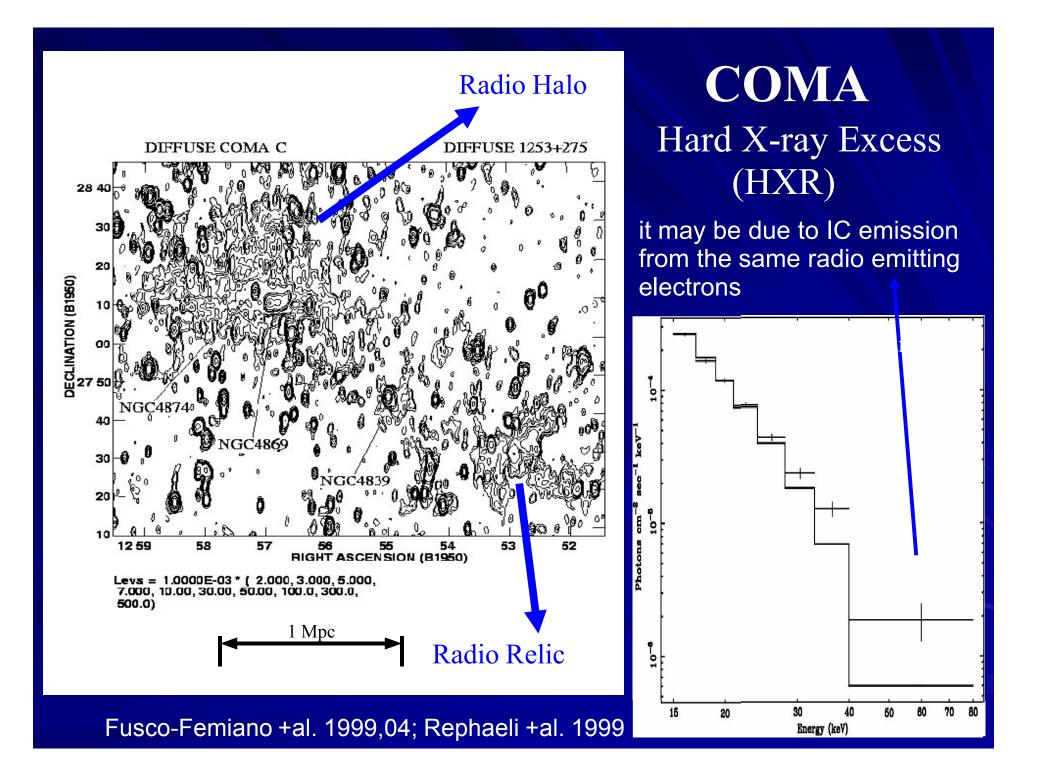
Diffusion and transport of metals in the ICM

(e.g., Voigt & Fabian 2004; Rebusco +al. 2005)

B-Amplification from Cosmological seed fields (e.g., Dolag +al. 1999,02; Subramanian +al. 2006)

Diffusion and scattering of HE & UHECR in the Universe (e.g., Sigl +al. 2005; Dolag +al. 2005)





CR populations in GC

Long living population of CR protons

 $p + p \rightarrow \pi^0 + \pi^+ + \pi^- + anything$

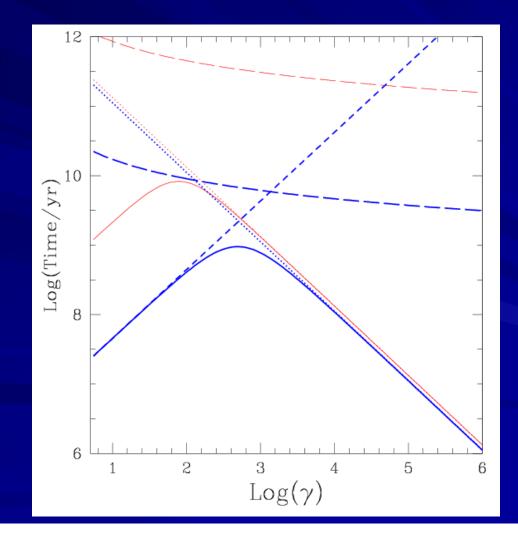
$$\pi^0 \to \gamma \gamma$$
$$\pi^{\pm} \to \mu + \nu_{\mu} \quad \mu^{\pm} \to e^{\pm} \nu_{\mu} \nu_e.$$

Stationary (continuous injection) population of CR (secondary) electrons

Transient population of CR electrons (< 1 Gyr)

Physics of CR Leptons

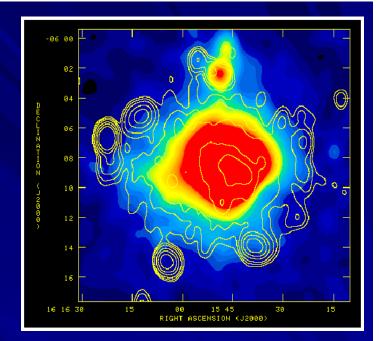
$(dE/dt) \sim E / \underline{Time} \sim m_e c^2 b$



Diffuse Radio Emission & origin of emitting electrons

(G.Brunetti, T.Venturi, S.Giacintucci, R.Cassano, D.Dallacasa, G.Setti, R.Athreya, N.Kassim, W.Lane, K.Dolag, S.Bardelli, B.Cotton, P.Mazzotta, M.Markevitch)

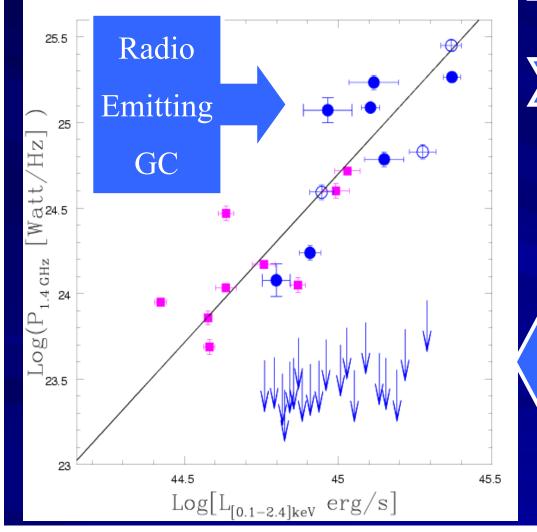




Sample of <u>50 massive GC z=0.2-0.4</u> (REFLEX + eBCS)

Similar z Similar X-luminosities

Brunetti +al. 2007



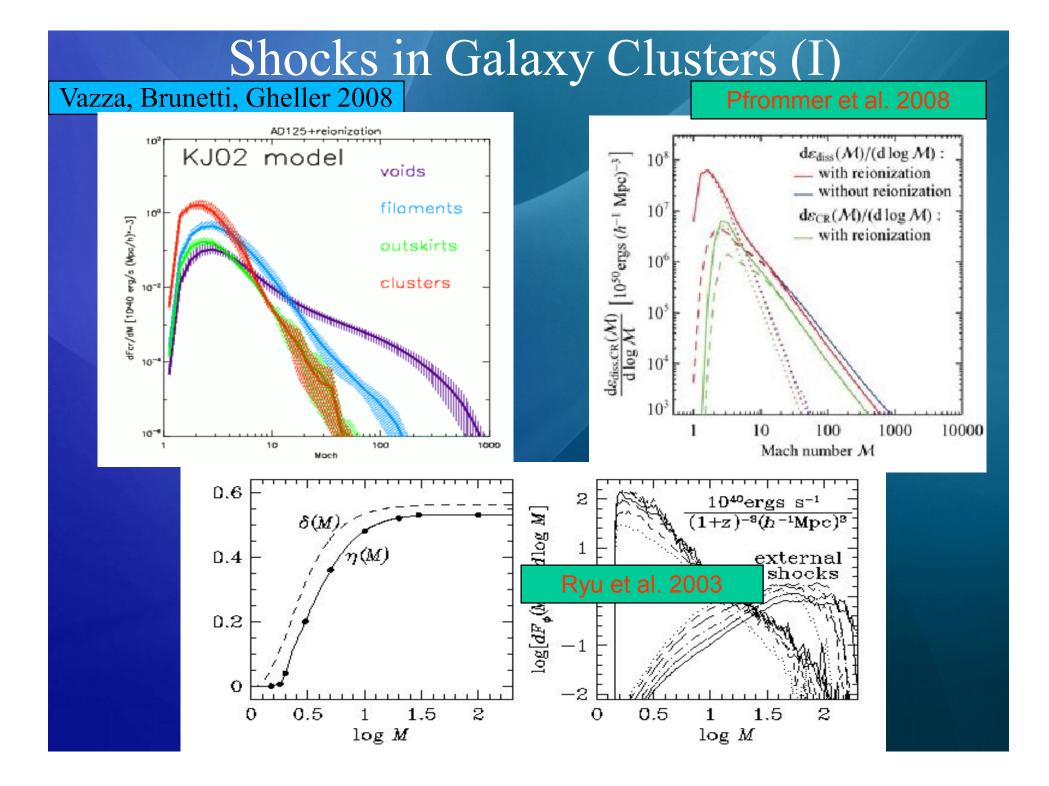
RH are not common

Bi-modality of GC

Radio

Quiet

GC



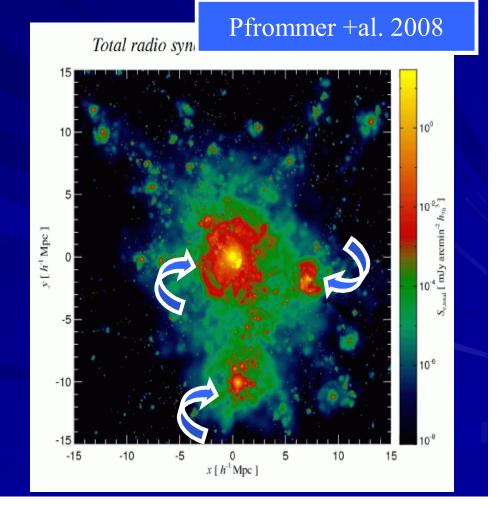
Disagreement with shock+secondary models

Brunetti +al. 2007 Pfrommer +al. 2008 Total radio syn 25.5 10⁰ 25 Hz \oplus (11°24.5 Mgt/ $\int_{neat}^{\infty} [mJy \operatorname{arcmin}^2 h_{\gamma_0}^{\circ}]$ y [h⁻¹ Mpc] $Log(P_{1.4 \text{ GHz}})$ 24 Ś 10-6 23.5 -10 23 44.5 45.5 45 15 -15 -10 -5 0 5 10 $x [h^{-1} Mpc]$ $Log[L_{[0.1-2.4]keV} erg/s]$

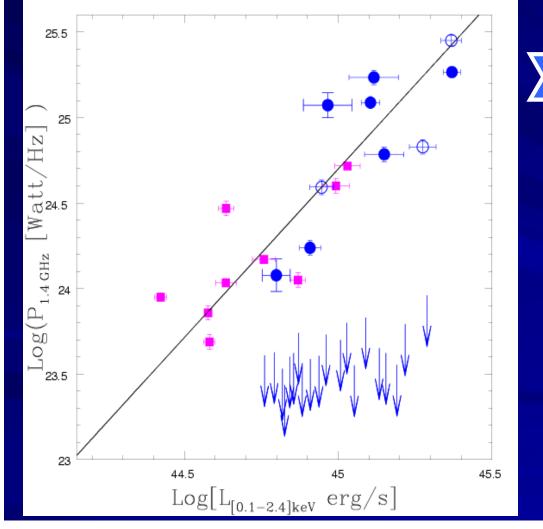
Disagreement with shock+secondary models

25.5 25 Hz \oplus (11°24.5 Mgt/ $Log(P_{1.4 \text{ GHz}})$ 24 23.5 23 44.5 45.5 45 $Log[L_{[0.1-2.4]keV} erg/s]$

Brunetti +al. 2007



Brunetti +al. 2007

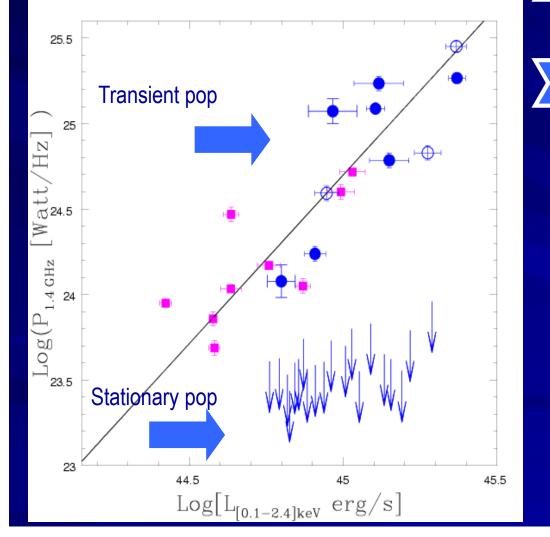


RH are not common

Bi-modality of GC

γ≈10⁴ e[±] are not common in GC (on-off mechanism)

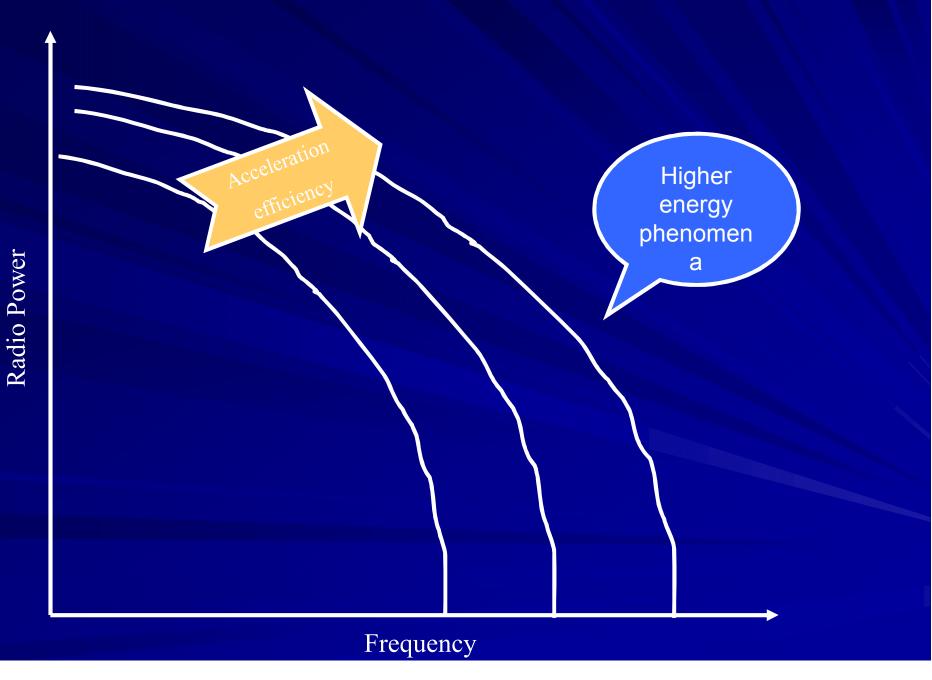
Brunetti +al. 2007

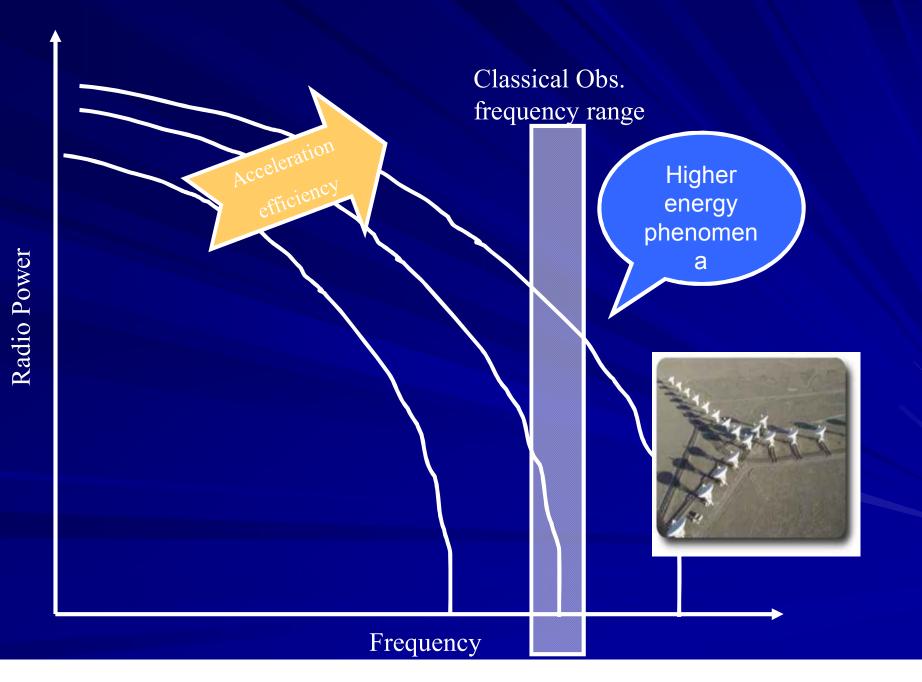


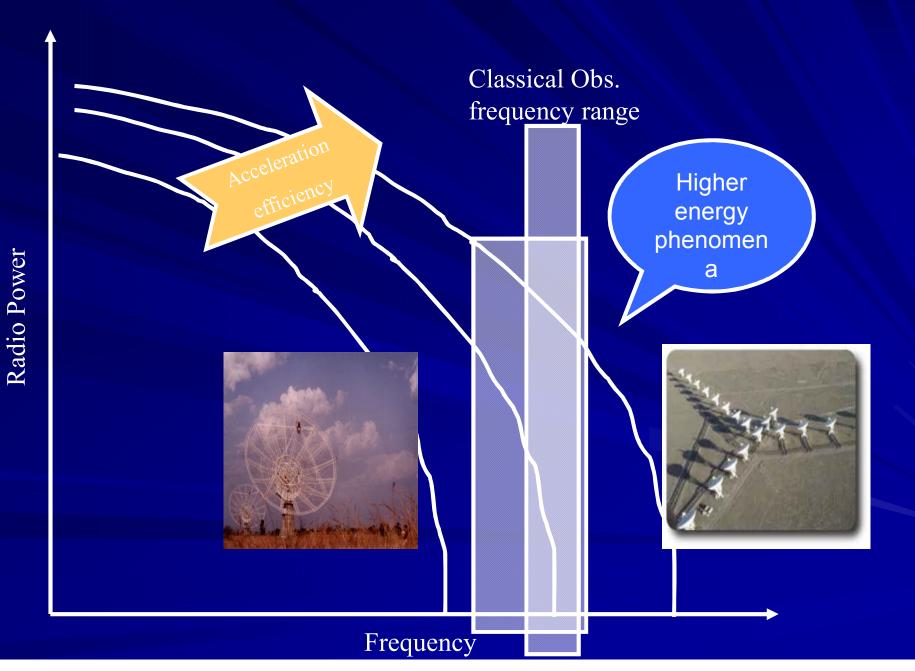
RH are not common

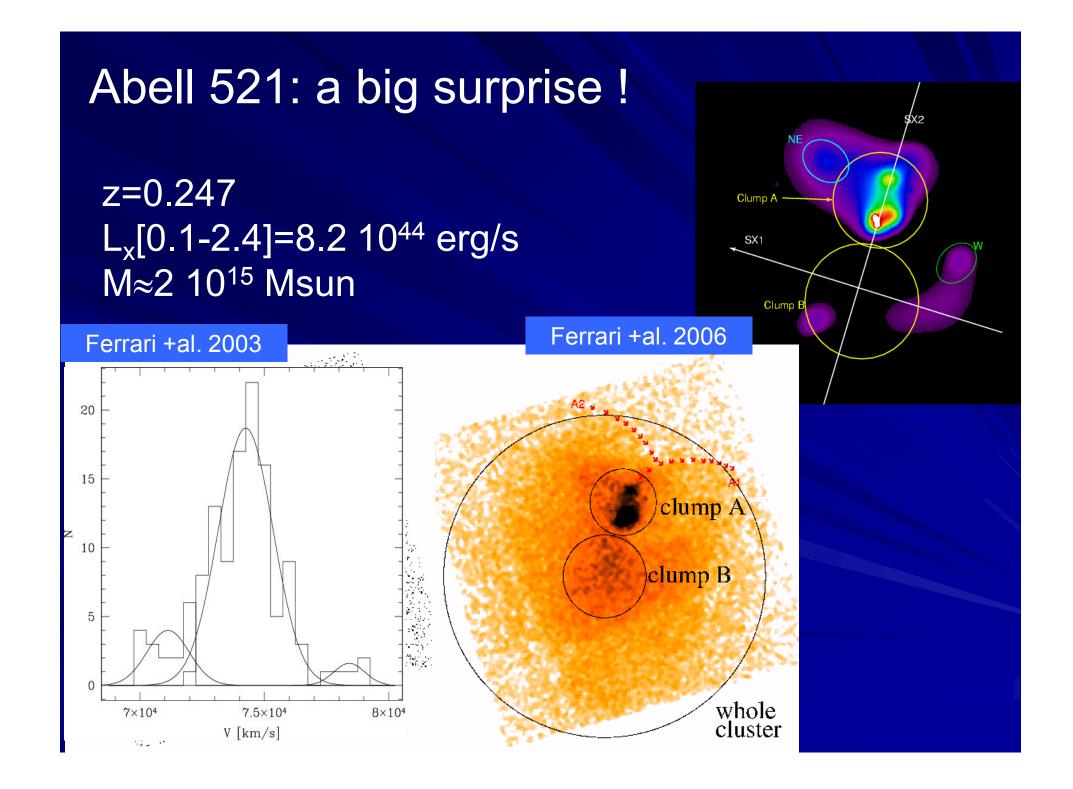
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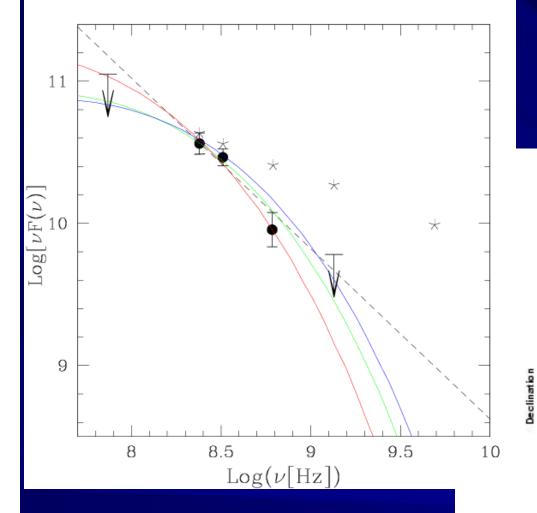






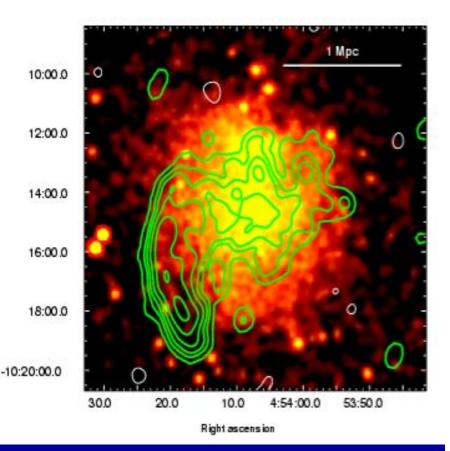


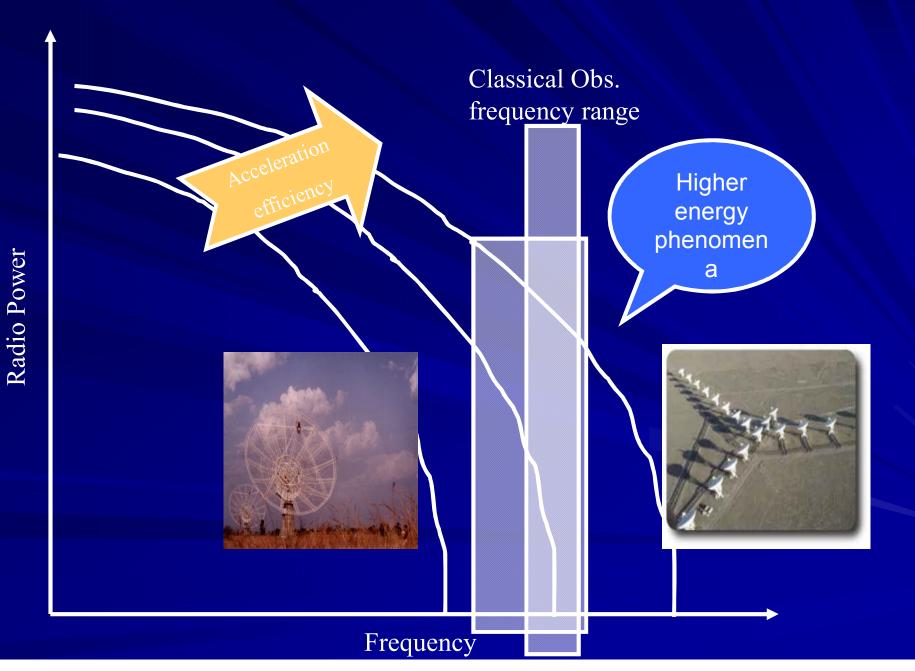
Abell 521: a big surprise !



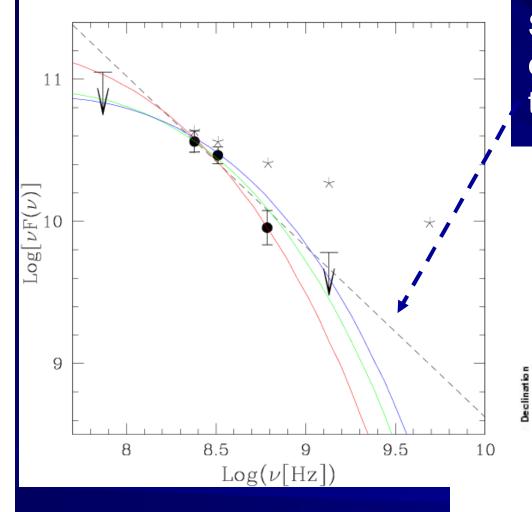
The first Radio Halo with ultra steep spectrum

The first low frequency Radio Halo !

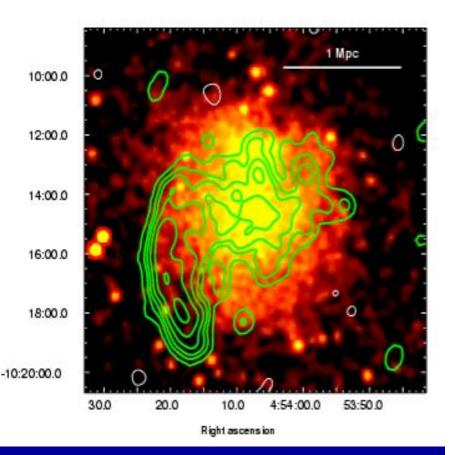




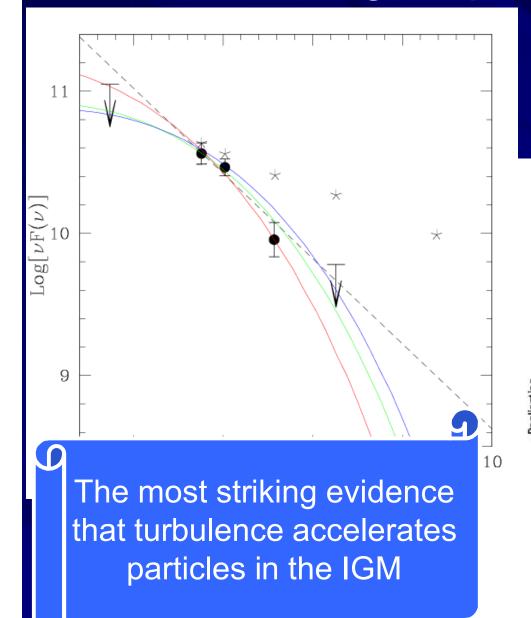
Abell 521: a big surprise !



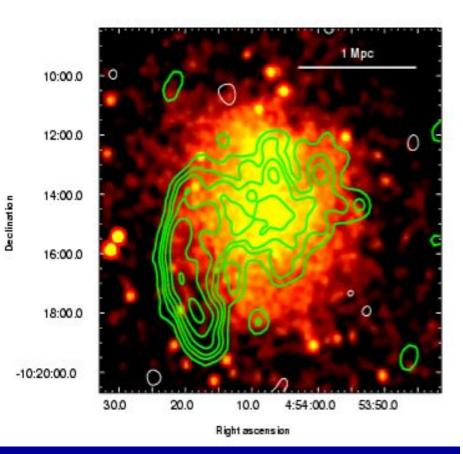
The first Radio Halo with ultra steep spectrum Secondary models require energy in CR protons larger than the cluster thermal energy

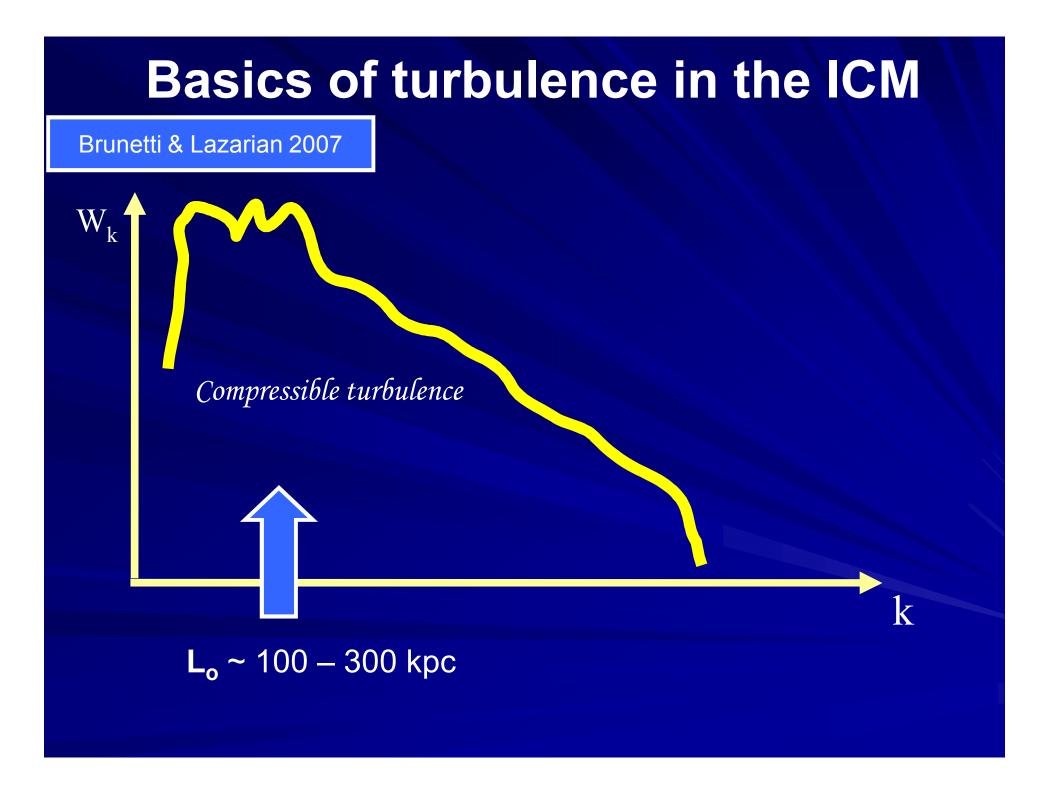


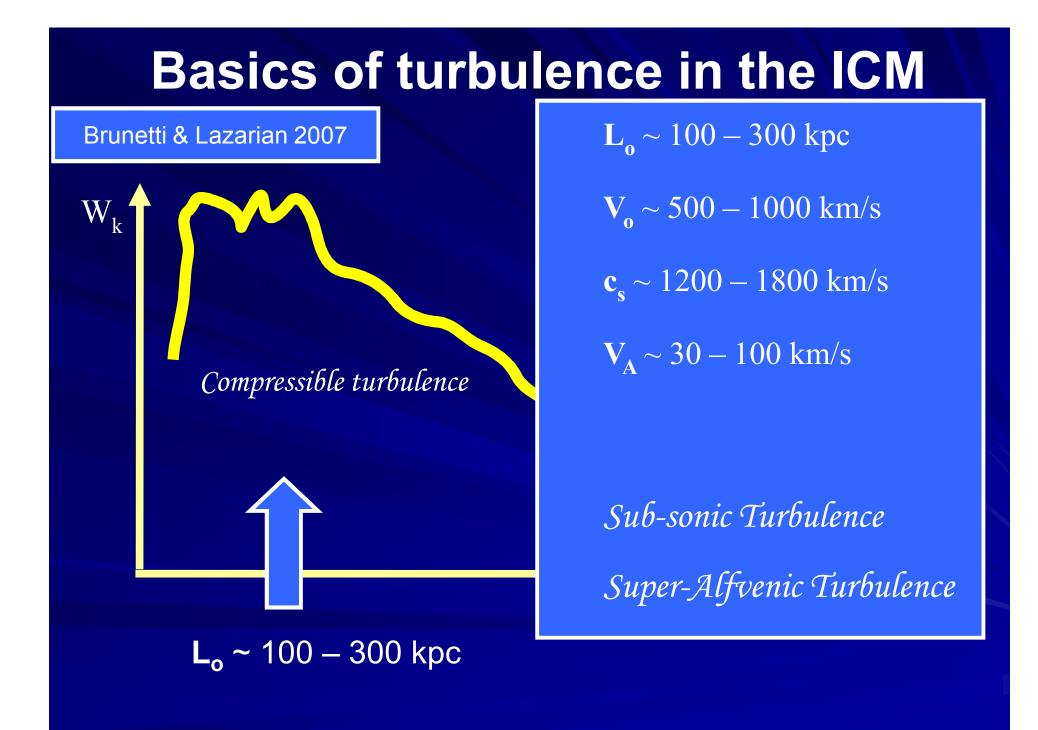
Abell 521: a big surprise !

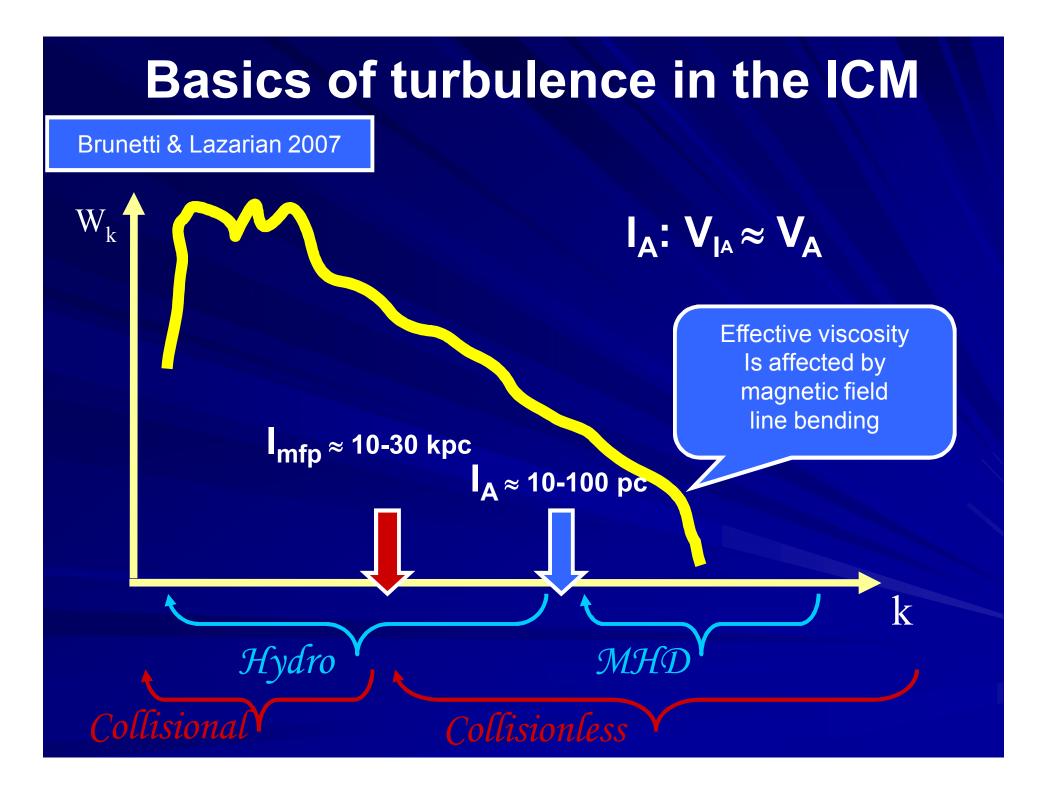


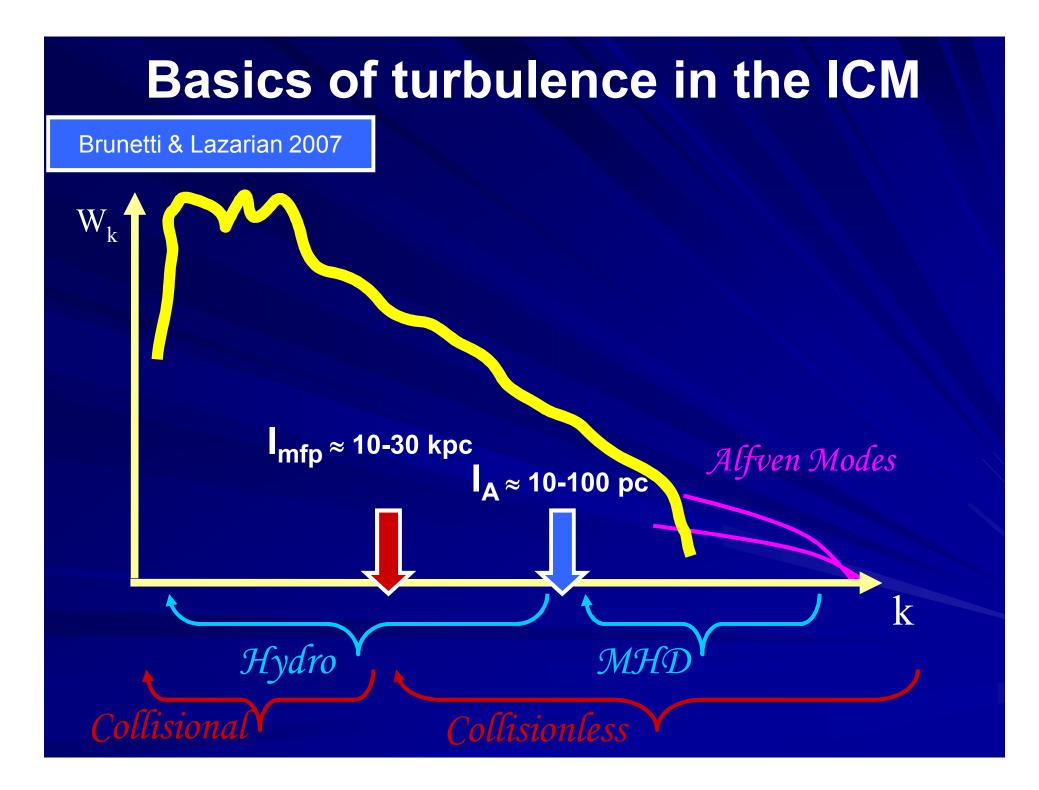
Secondary models require energy in CR protons larger than the cluster thermal energy





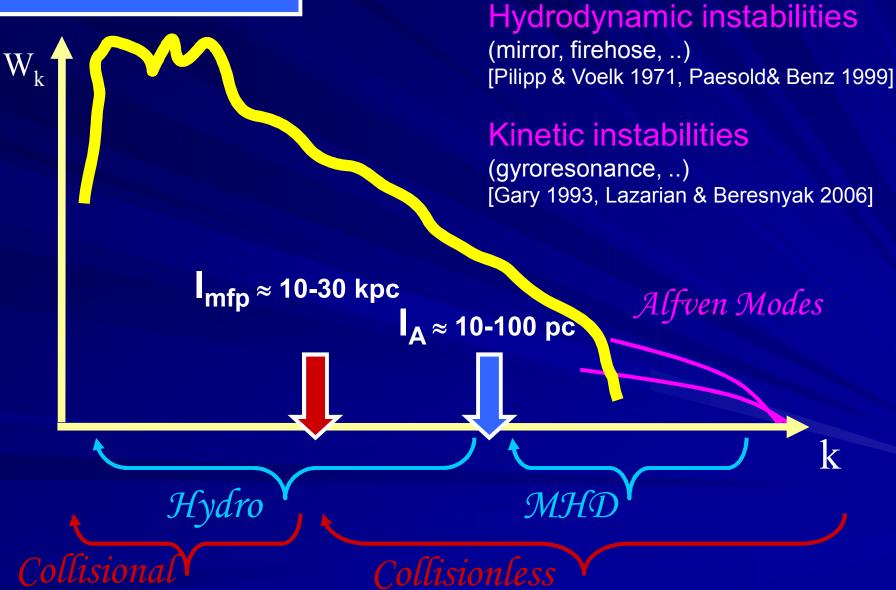


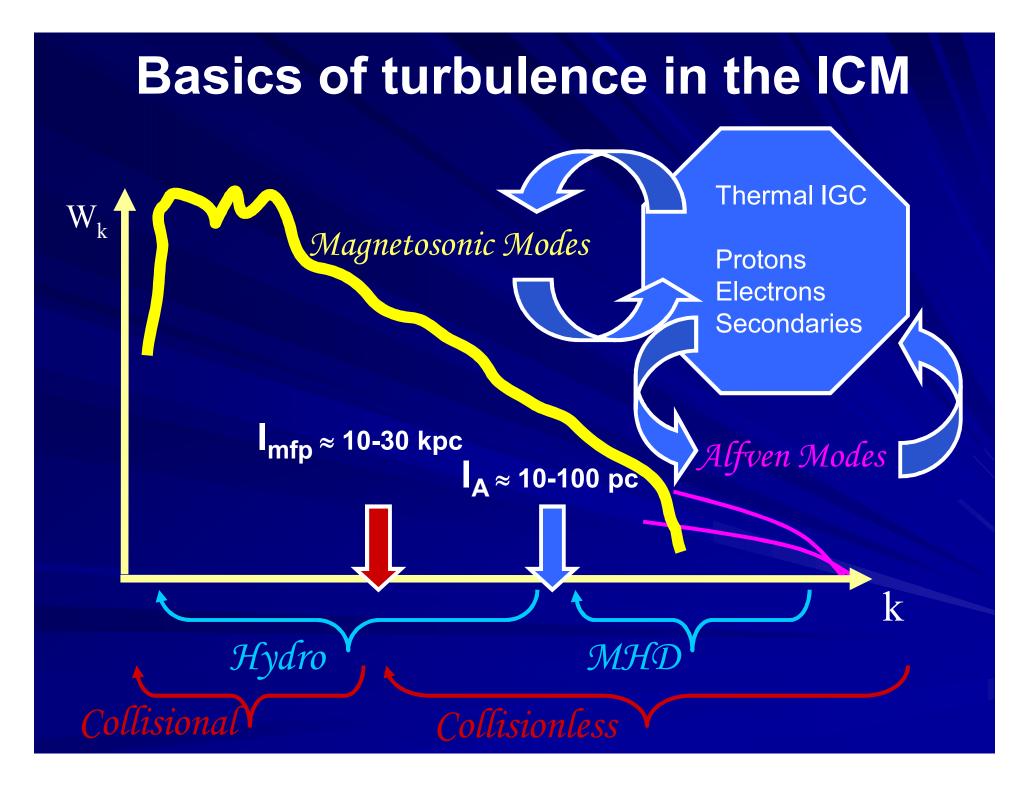






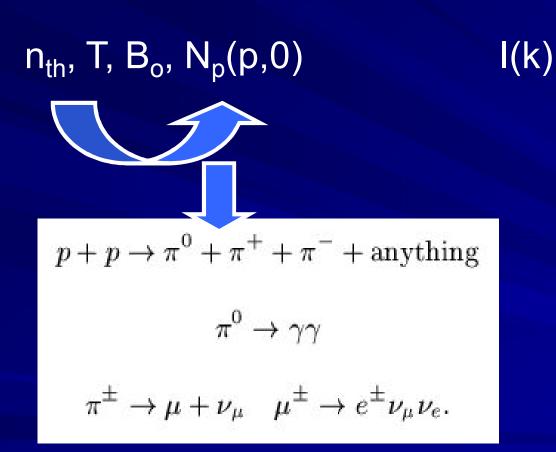
Brunetti & Lazarian 2007

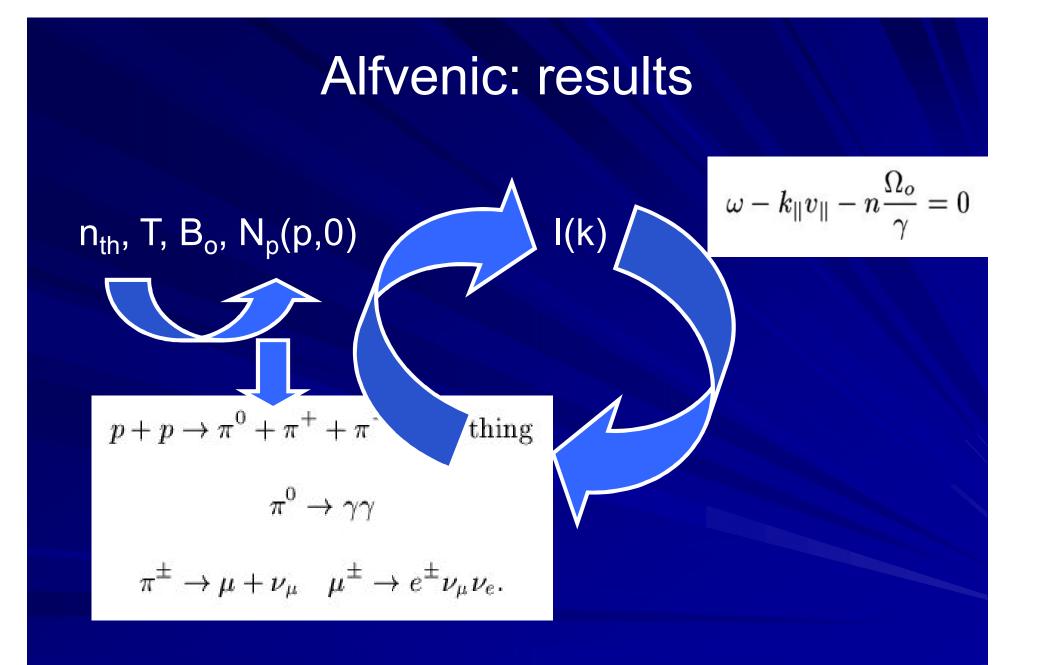




n_{th}, T, B_o, N_p(p,0)

l(k)





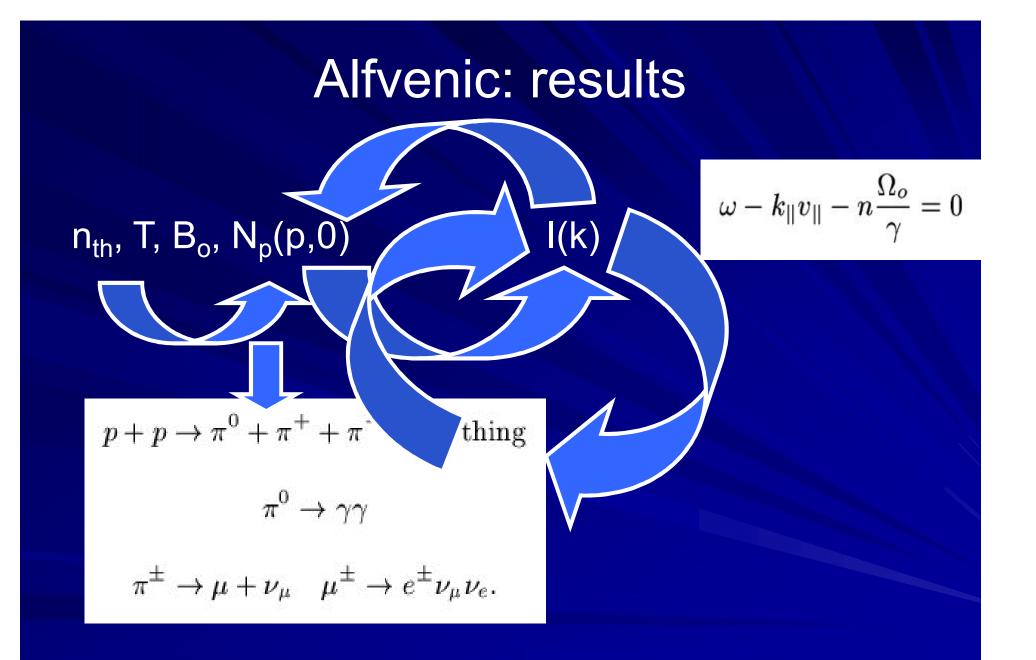
$$n_{th}, T, B_{o}, N_{p}(p, 0) \qquad I(k)$$

$$p + p \rightarrow \pi^{0} + \pi^{+} + \pi^{-} + anything$$

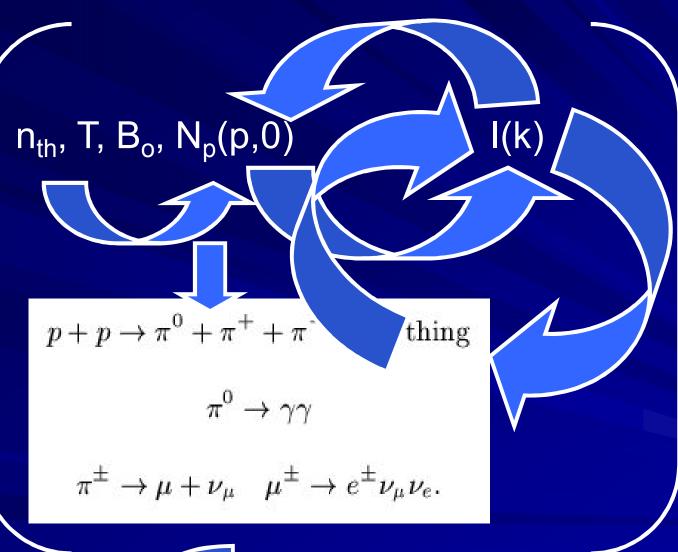
$$\pi^{0} \rightarrow \gamma \gamma$$

$$\pi^{\pm} \rightarrow \mu + \nu_{\mu} \quad \mu^{\pm} \rightarrow e^{\pm} \nu_{\mu} \nu_{e}.$$

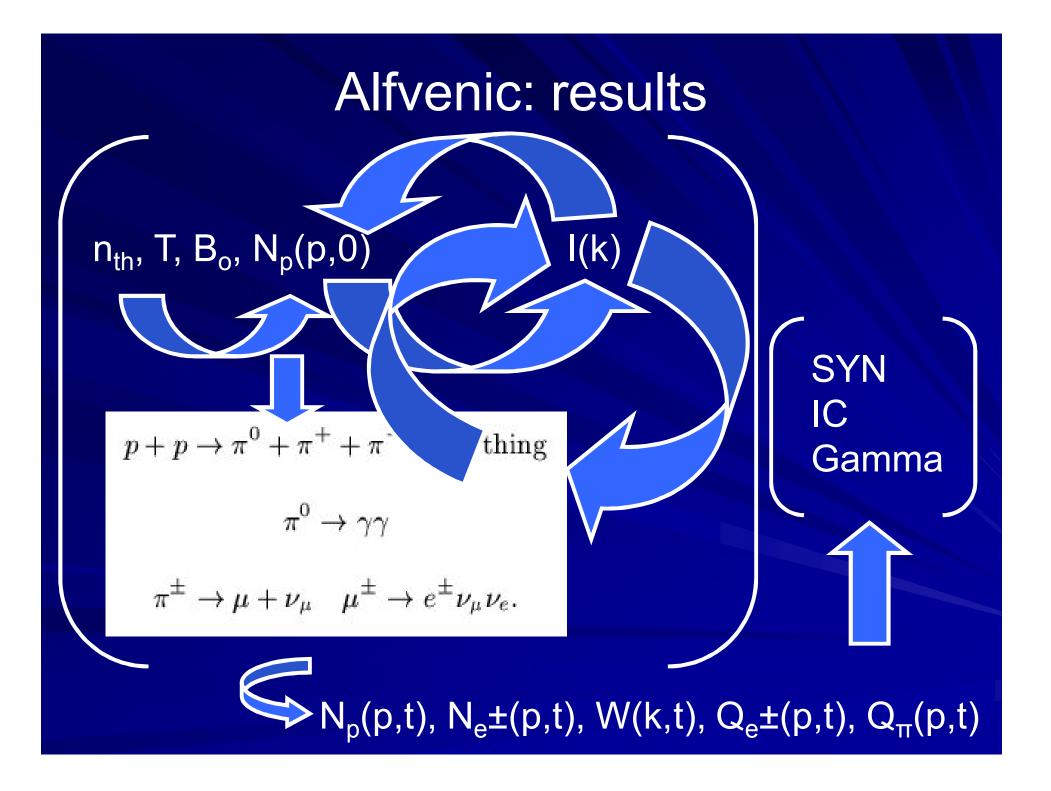
$$\omega - k_{\parallel} v_{\parallel} - n rac{\Omega_o}{\gamma} = 0$$

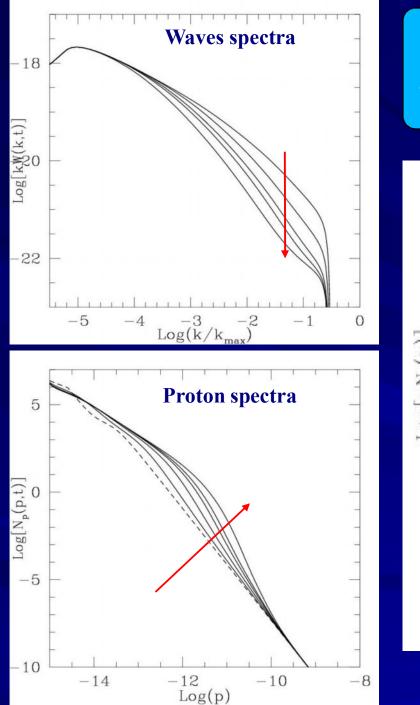


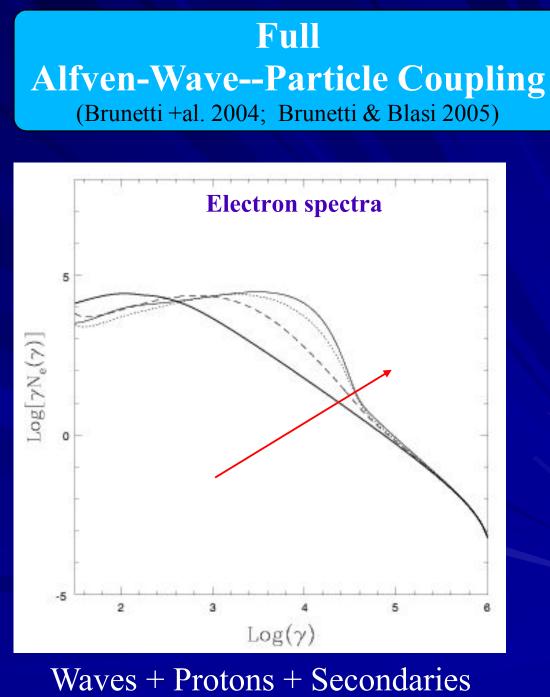


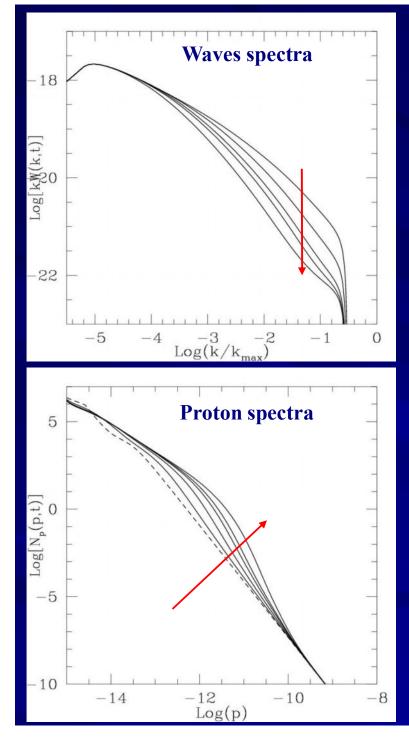


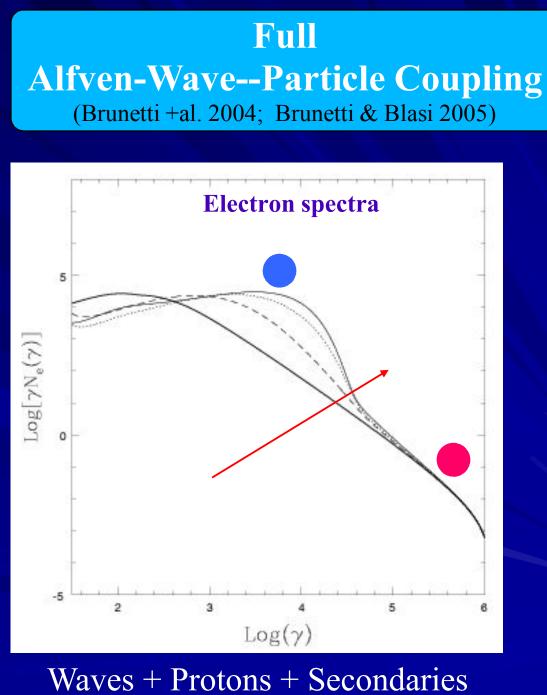
 $\gg N_p(p,t), N_e \pm (p,t), W(k,t), Q_e \pm (p,t), Q_{\pi}(p,t)$





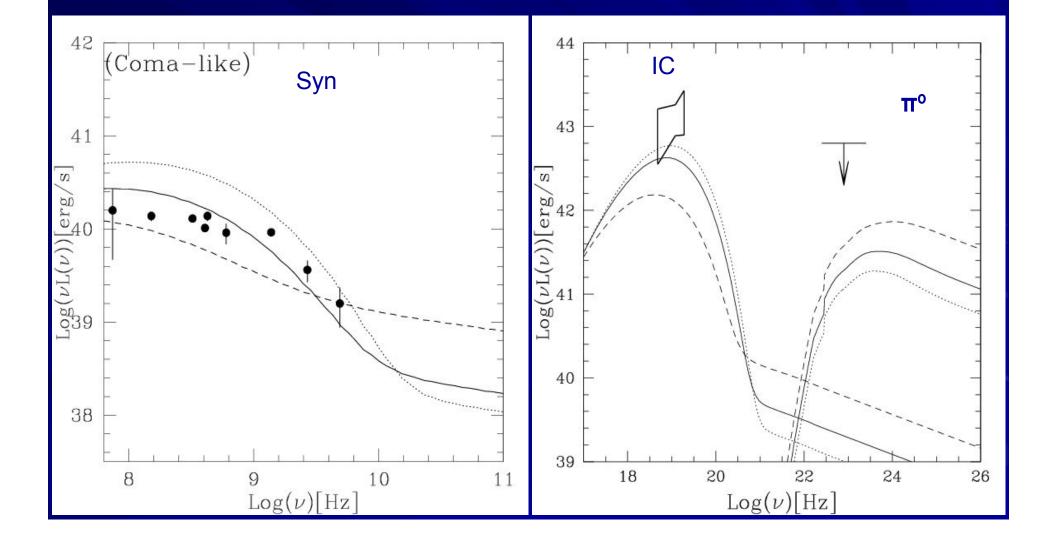




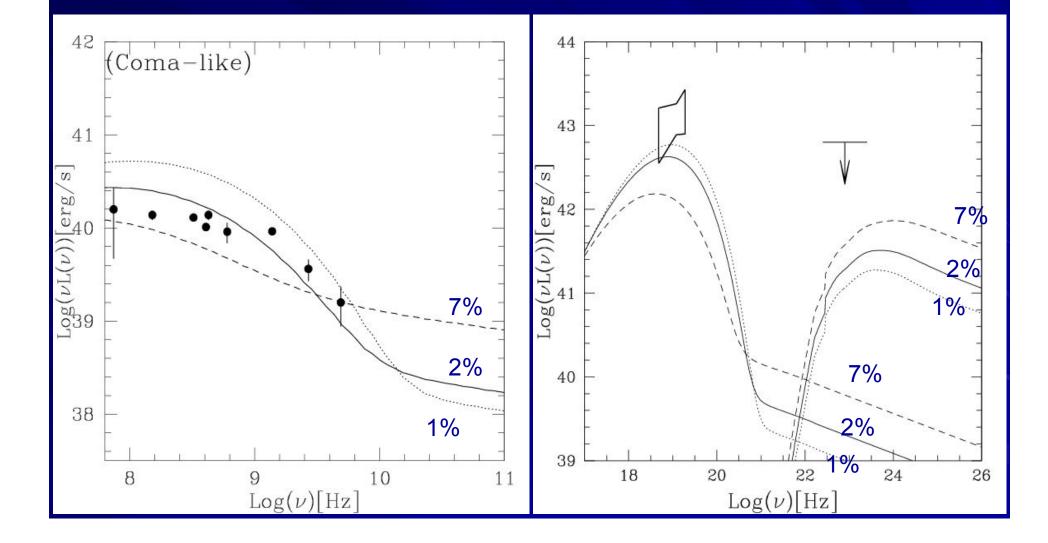


Brunetti, Blasi, Gabici, Cassano, in prep

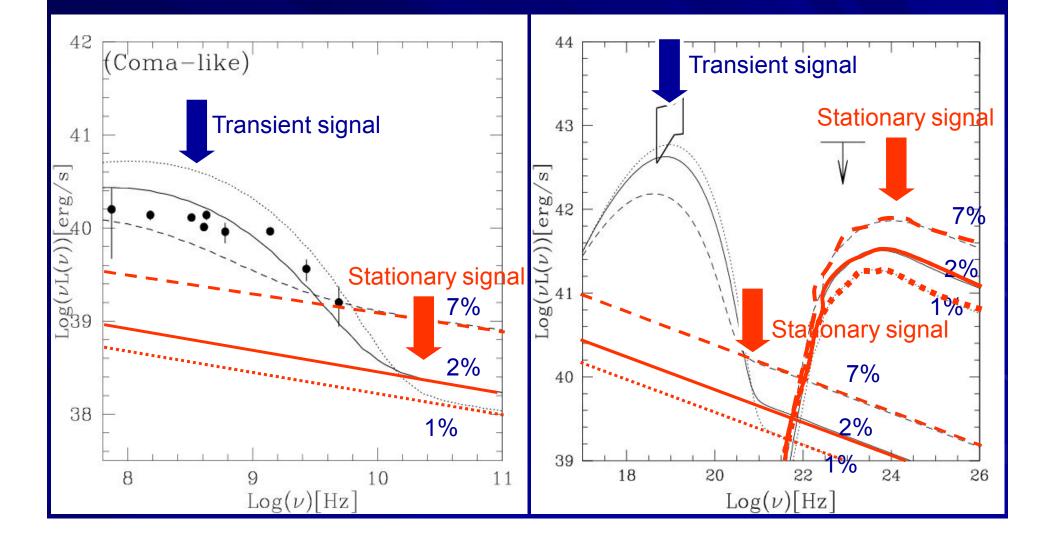
Toy Model: β -profile, $B_o \approx A n_{th}^{2/3}$, $B_o(0)=1.5\mu G$, $W_{CR} \approx f W_{th}$, $P_A \approx Q n_{th}^{5/6}$

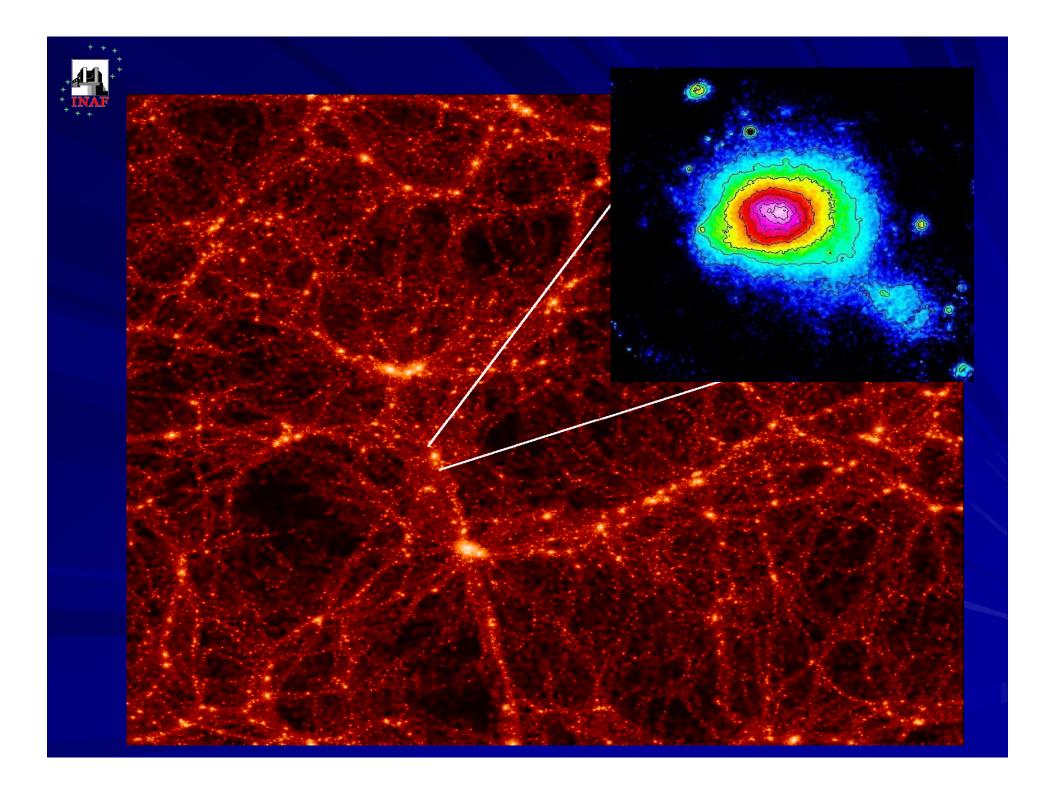


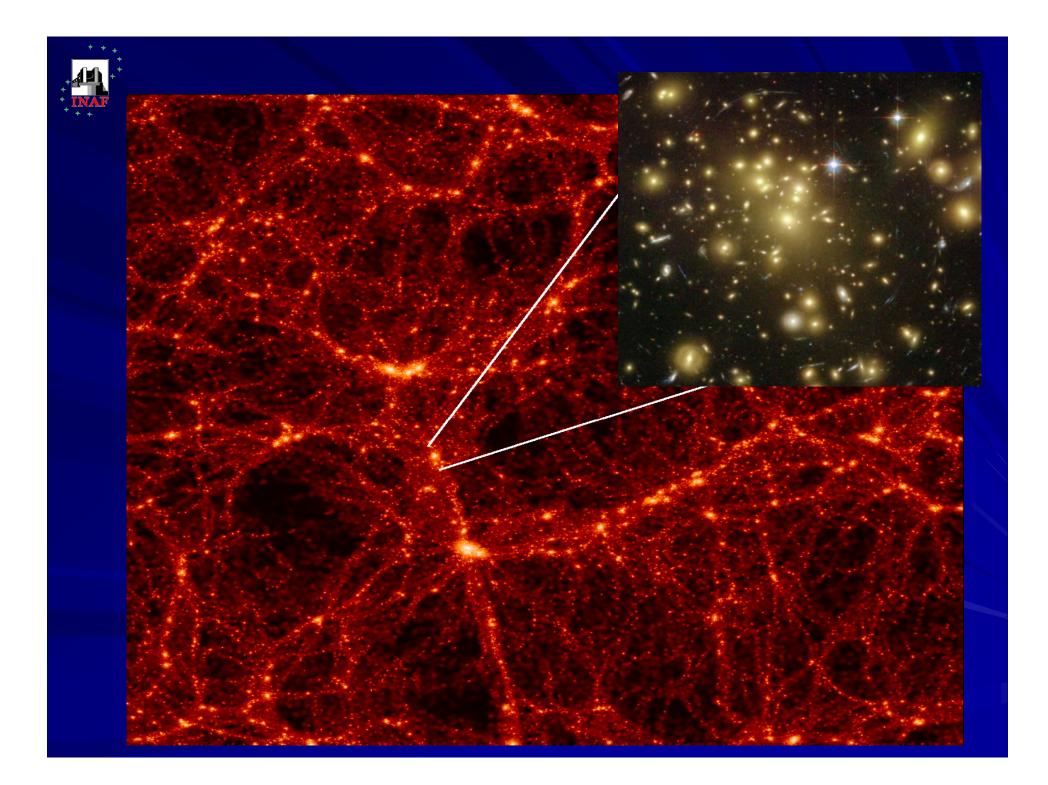
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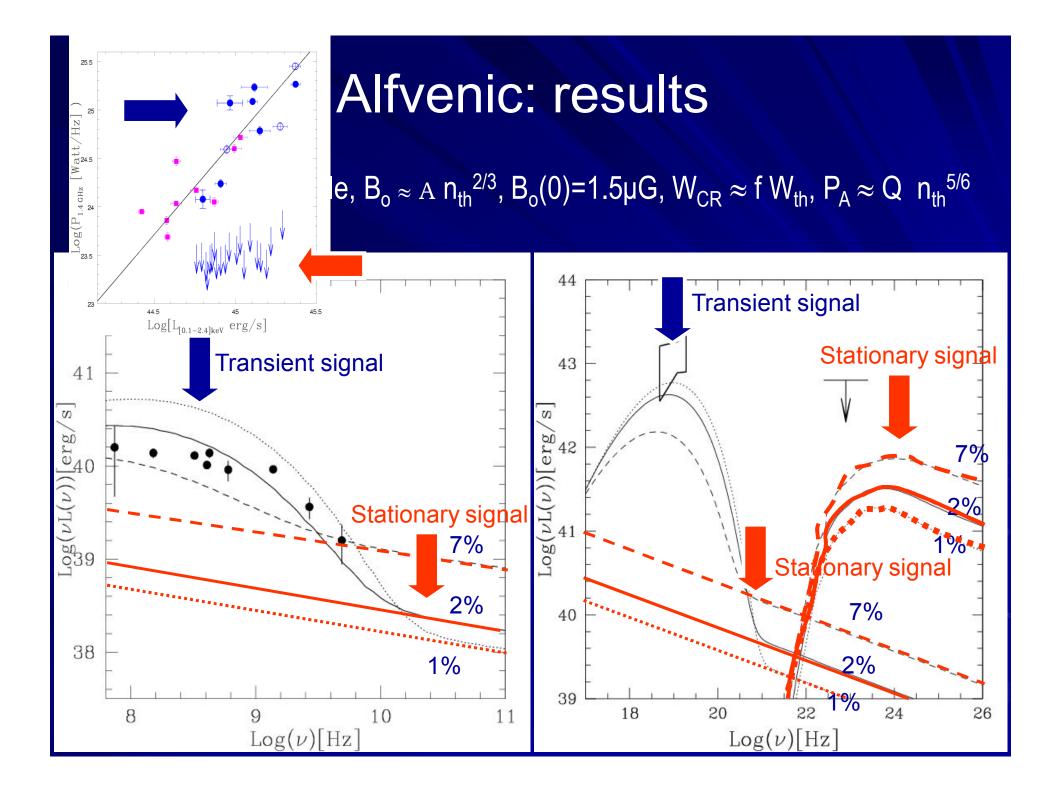


Toy Model: β -profile, $B_o \approx A n_{th}^{2/3}$, $B_o(0)=1.5\mu G$, $W_{CR} \approx f W_{th}$, $P_A \approx Q n_{th}^{5/6}$



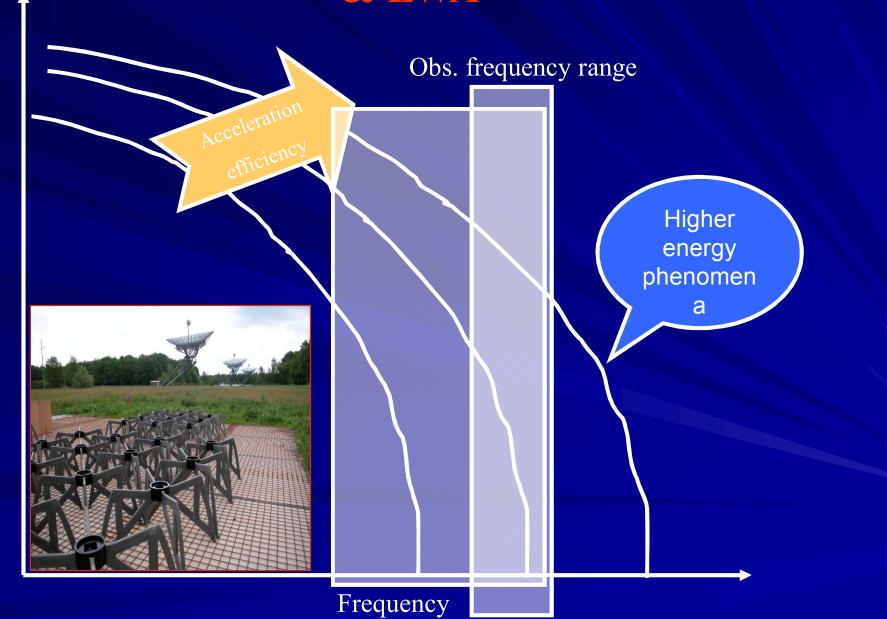








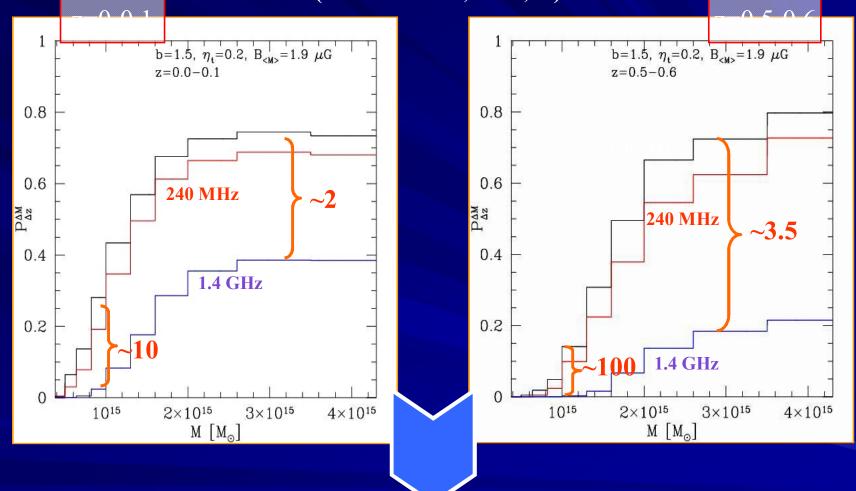
LOW freq observations: LOFAR & LWA



Radio Power

Fraction of GC with RHs at lower frequencies

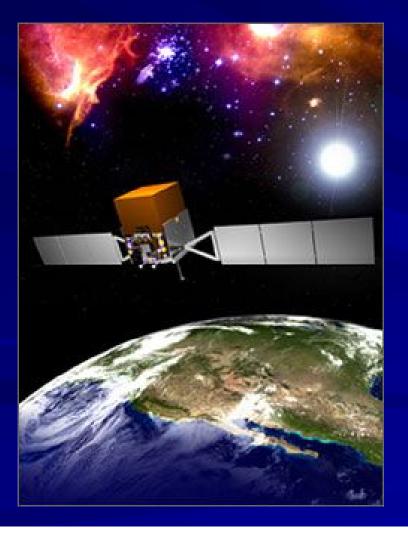
(Cassano et al., 2006,08)

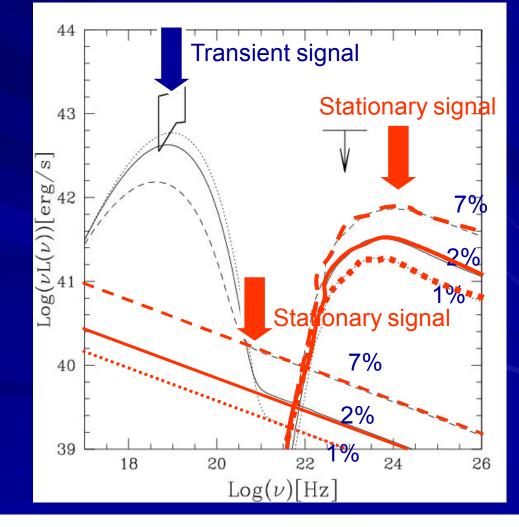


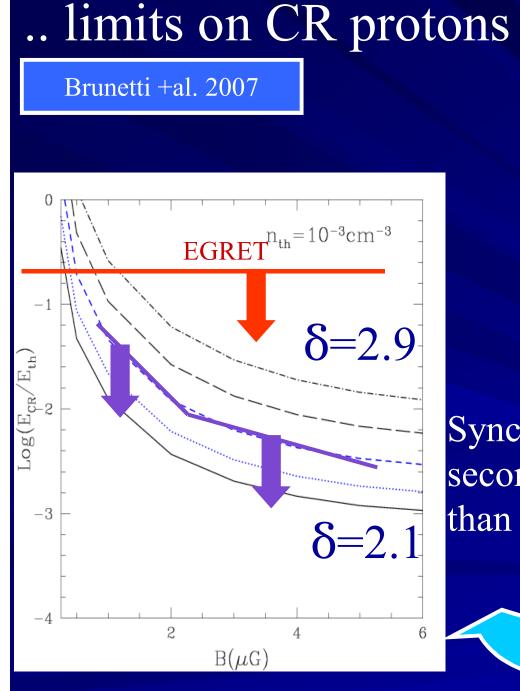
Expectation : Ultra-Steep-Spectrum Radio Halos

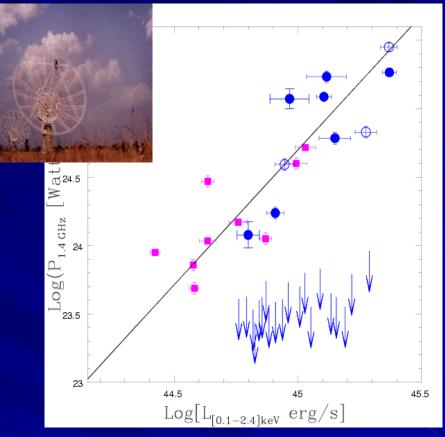


The gamma ray signal yields a direct estimate of CR protons





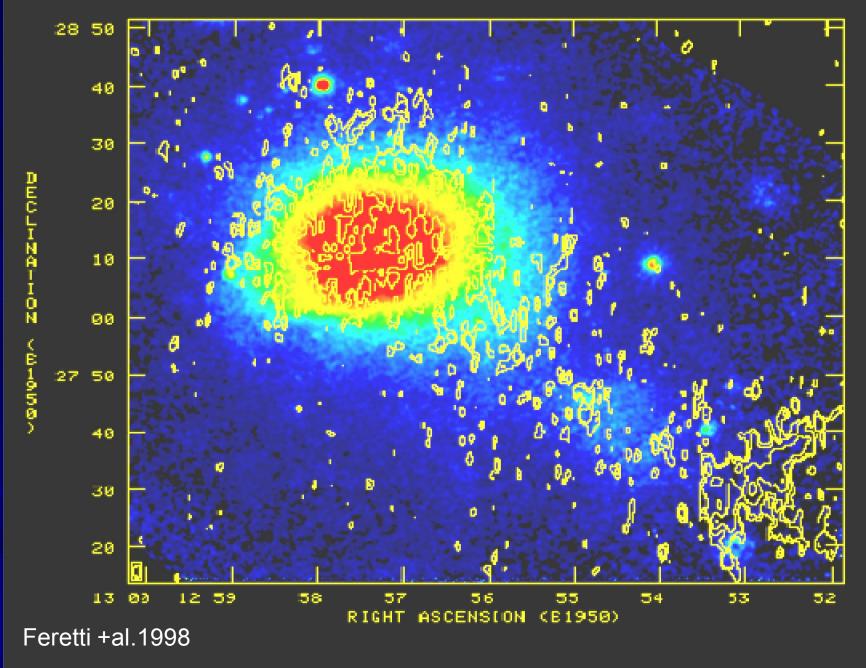




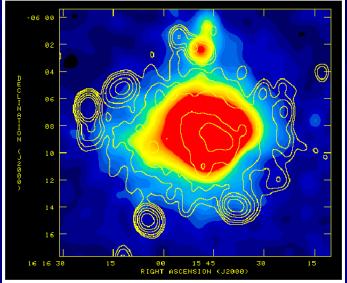
Synchrotron emission from secondaries should be smaller than upper limits .

limit on : B , E_{CRp} , δ

Coma Cluster

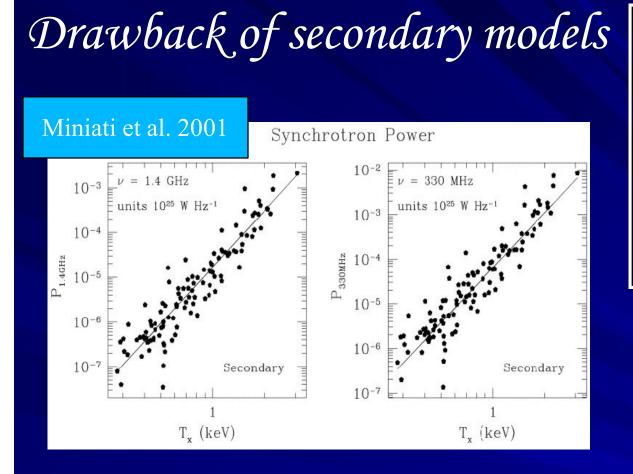


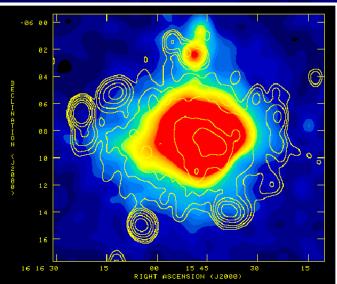
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First possibility: *secondary models*, relativistic electrons continuously injected in the ICM by inelastic proton-proton collisions through productions and decay of charged pions (e.g., *Dennison 1980, Blasi & Colafrancesco 1999, Dolag & Ensslin 2000; Pfrommer & Ensslin 2004*)

Second possibility : *in situ* re-acceleration by MHD turbulence developed in the cluster volume during the merger events (e.g., Brunetti et al. 2001, 2004; Petrosian 2001; Ohno et al. 2002; Fujita et al. 2003; Brunetti & Blasi 2005; Cassano & Brunetti 2005; Brunetti & Lazarian 2007; Petrosian & Bykov 2008)



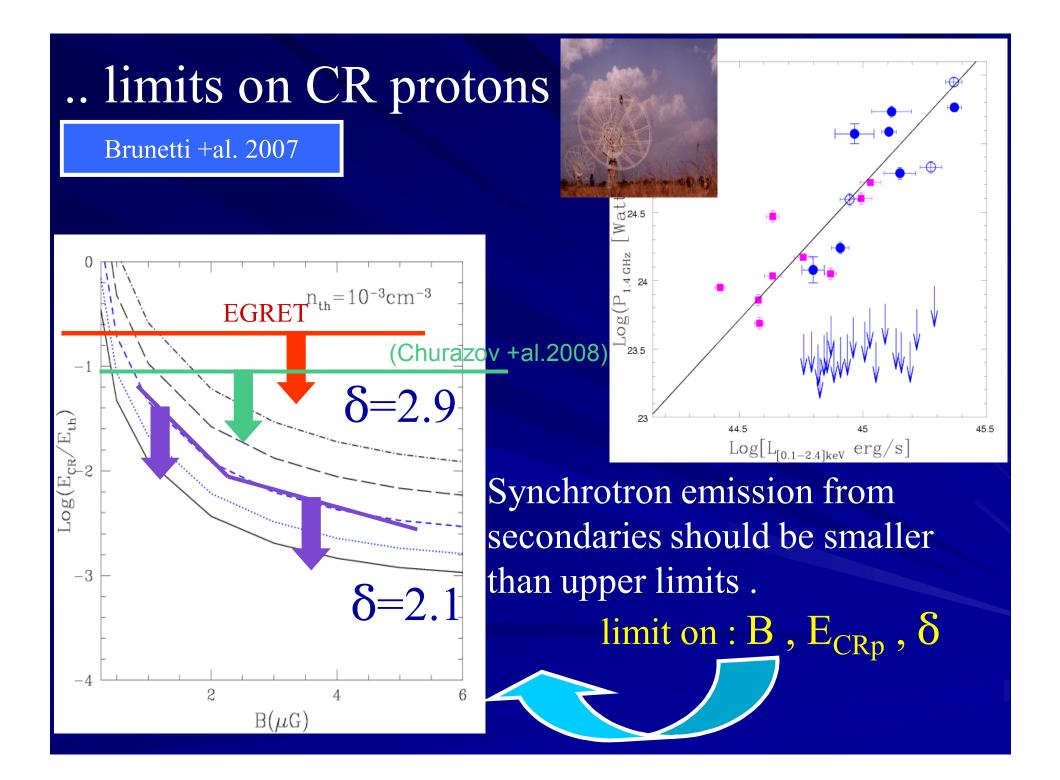


In case clusters have similar magnetic fields :

 \sum

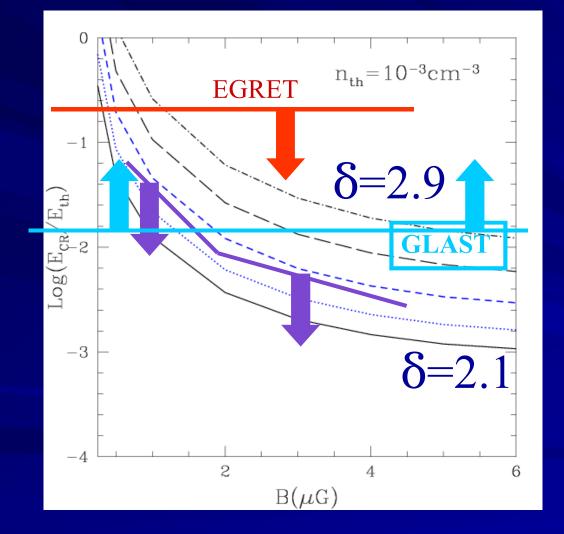
Radio Halos are long living (cluster "life-time") phenomena (difficult to understand connection with mergers)

Clusters with similar thermal properties have similar radio halos



GLAST: CR & B in GC

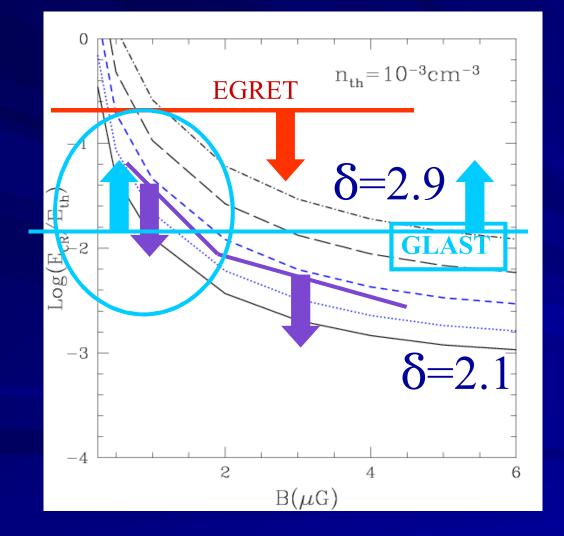
Brunetti +al. 2007



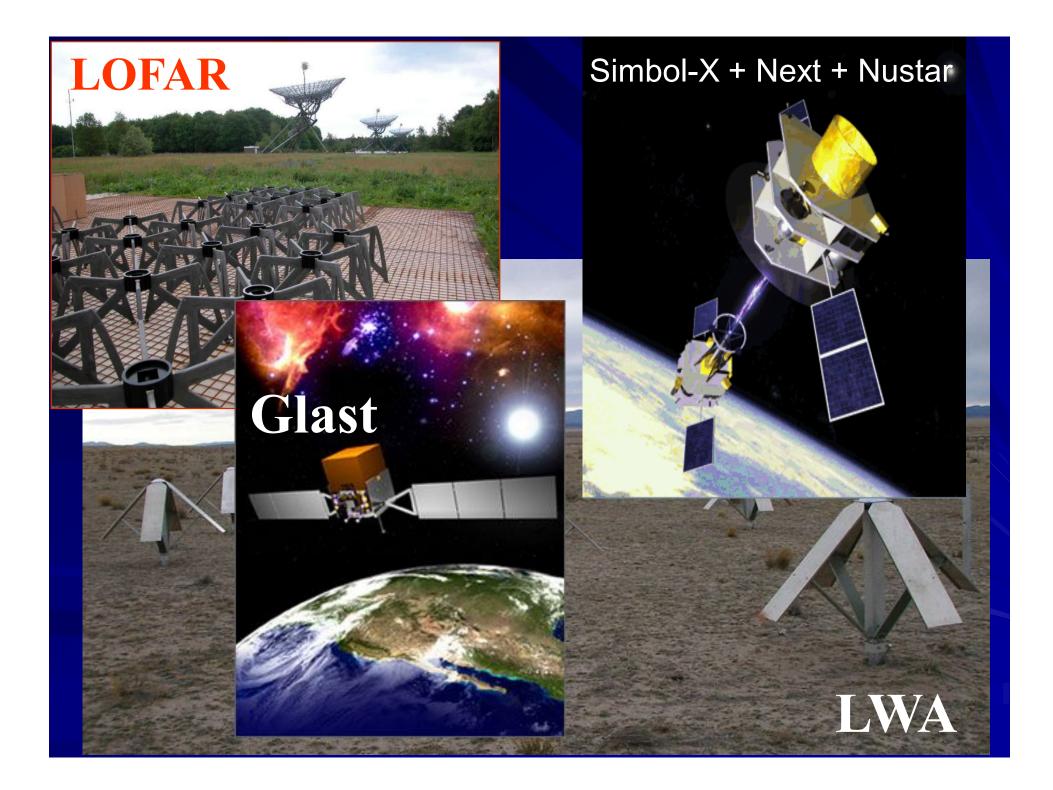
The combination of synchrotron observations (limits) of clusters without RHs with GLAST observations will constrain also B !

GLAST: CR & B in GC

Brunetti +al. 2007

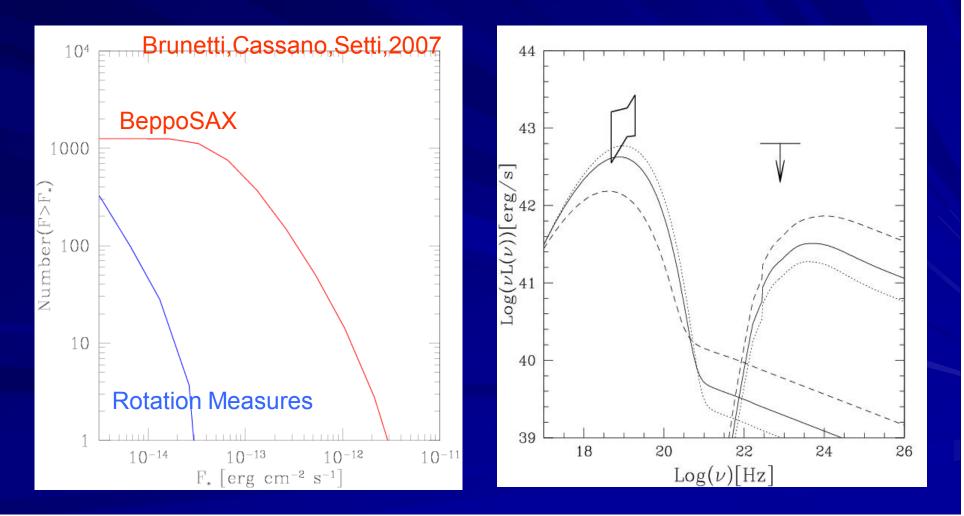


The combination of synchrotron observations (limits) of clusters without RHs with GLAST observations will constrain also B !



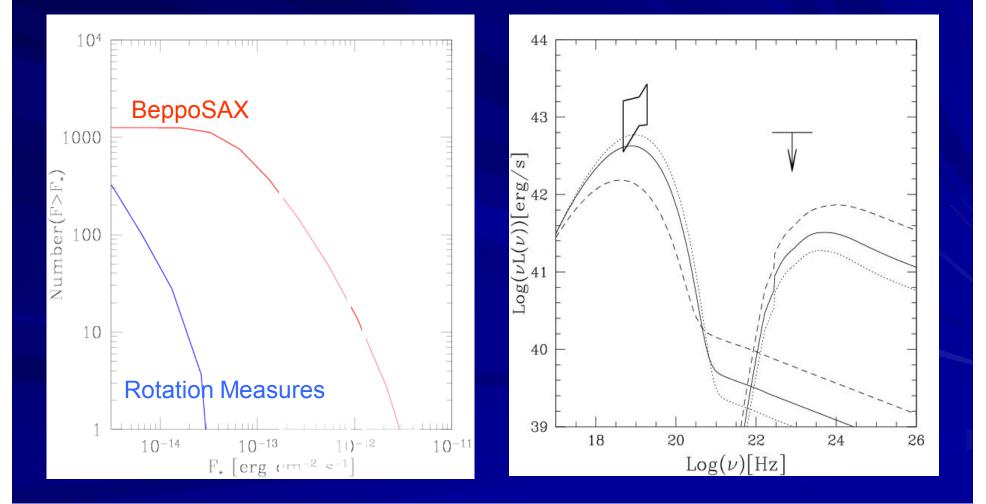
Simbol-X + Next + Nustar

The level of IC emission from the radio emitting regions in Radio Halos depends on the magnetic field.



Simbol-X + Next + Nustar

Future hard X-ray telescopes will give a final(?) answer on the strength of the magnetic field in the IGM



Summary

Turbulent acceleratio

Observations: bi-modality of GC (radio "quiet" & radio "loud")

Observations: break/cut-off in the synchrotron spectrum of RH

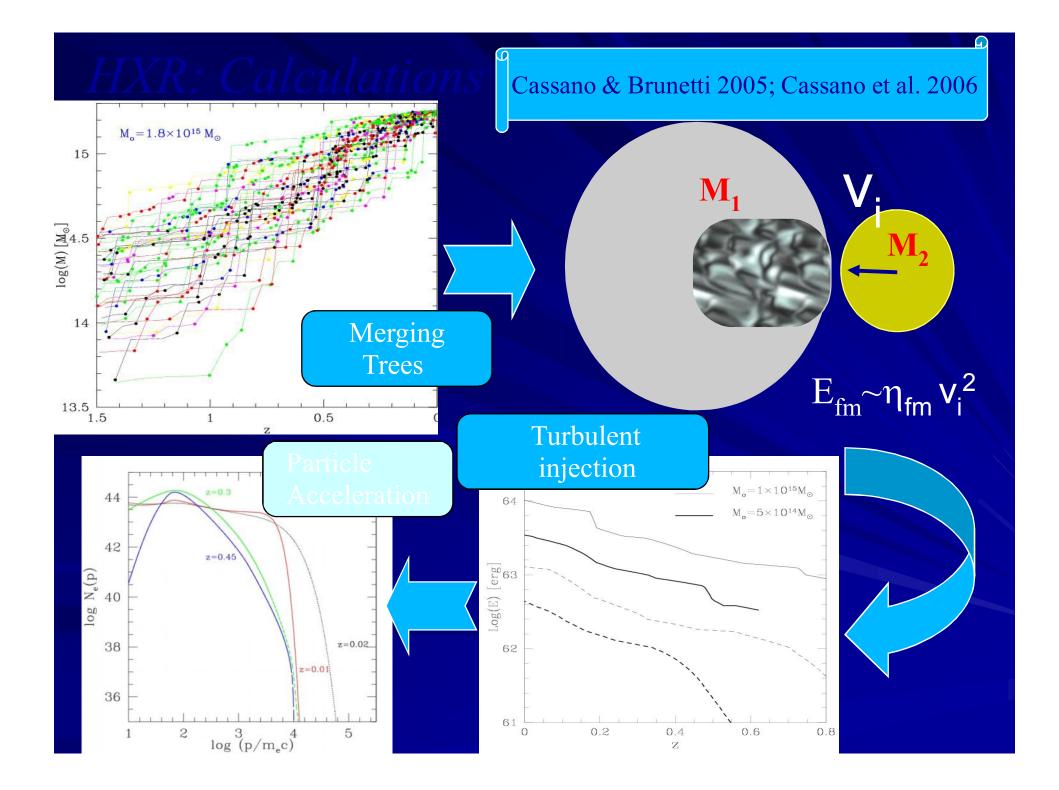
Observations: new class of RH (USSRH)

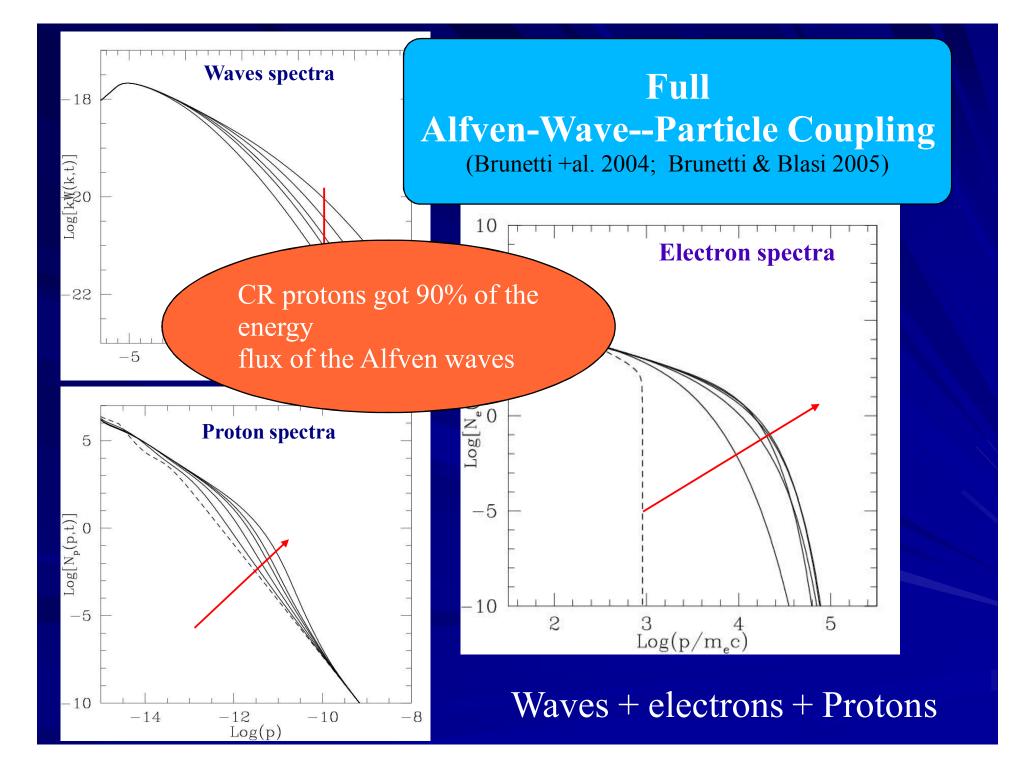
Models: magnetosonic acceleration (... need pitch angle scattering) (... need really turbulent media)

Models: Alfvenic acceleration

(... need injection driven by KPIns)

Future: USSRH and spectral studies of RH (LOFAR, LWA)
Hard-X-rays(Simbol-X,Nustar,Next)
(GLAST, Cerenkov)



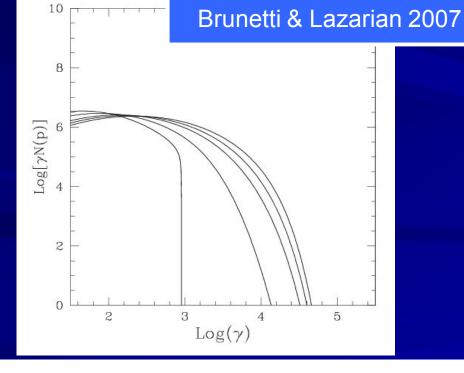


Stochastic Particle Acceleration (formalism)

$$\frac{\partial N_e(p,t)}{\partial t} = \frac{\partial}{\partial p} \left(N_e(p,t) \left[\left(\frac{dp}{dt} \right)_{rad} + \left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right)$$

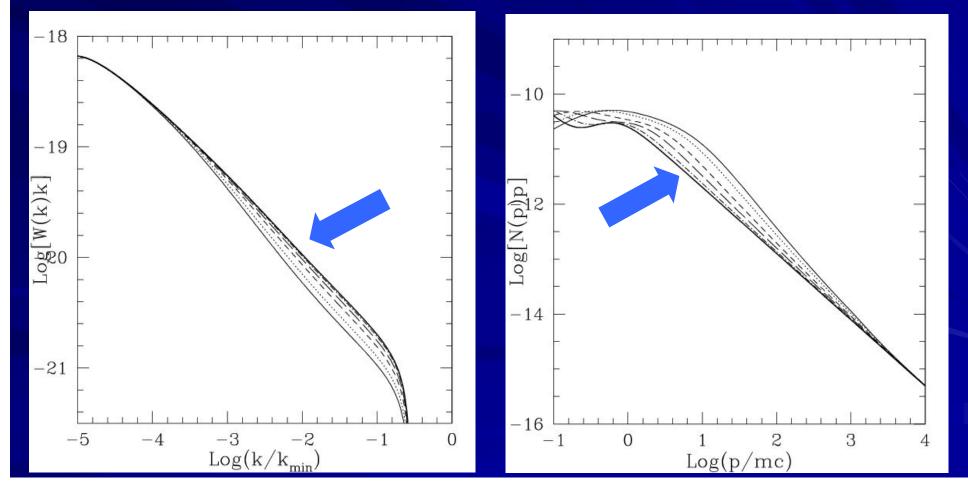
$$\frac{\partial \mathcal{W}(k,t)}{\partial t} = \frac{\partial}{\partial k} \left(k^2 D_{kk} \frac{\partial}{\partial k} \left(\frac{\mathcal{W}(k,t)}{k^2} \right) \right) - \sum_i \Gamma_i(k,t) \mathcal{W}(k,t) + I(k,t)$$

with thermal particles



n_{th}, T, B_o, N_p(p,0), I(k)

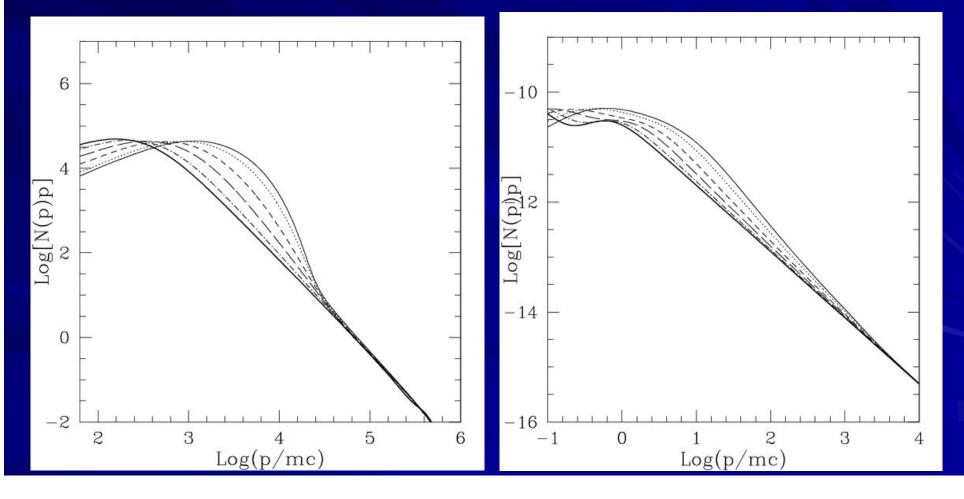
 $N_p(p,t), N_e \pm (p,t), W(k,t), Q_e \pm (p,t), Q_{\pi}(p,t)$



n_{th}, T, B_o, N_p(p,0), I(k)



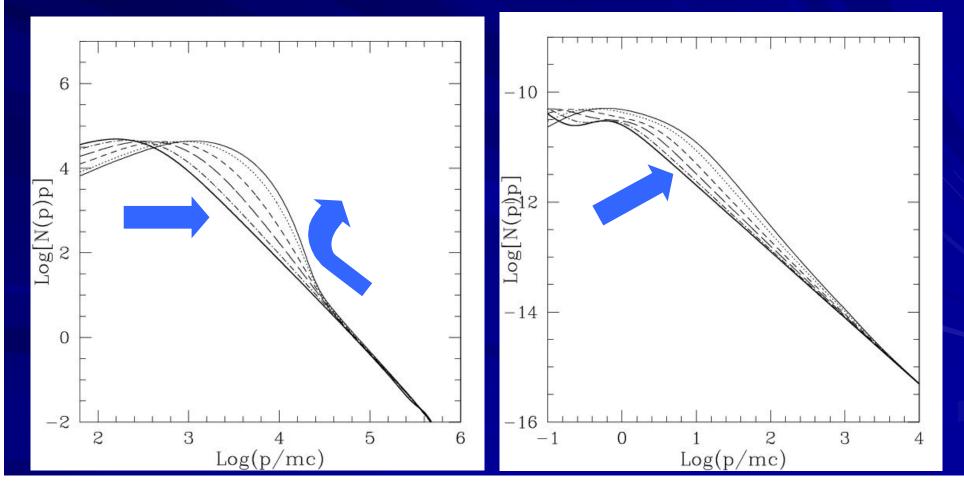
$N_p(p,t)$, $N_e \pm (p,t)$, W(k,t), $Q_e \pm (p,t)$, $Q_{\pi}(p,t)$



n_{th}, T, B_o, N_p(p,0), I(k)

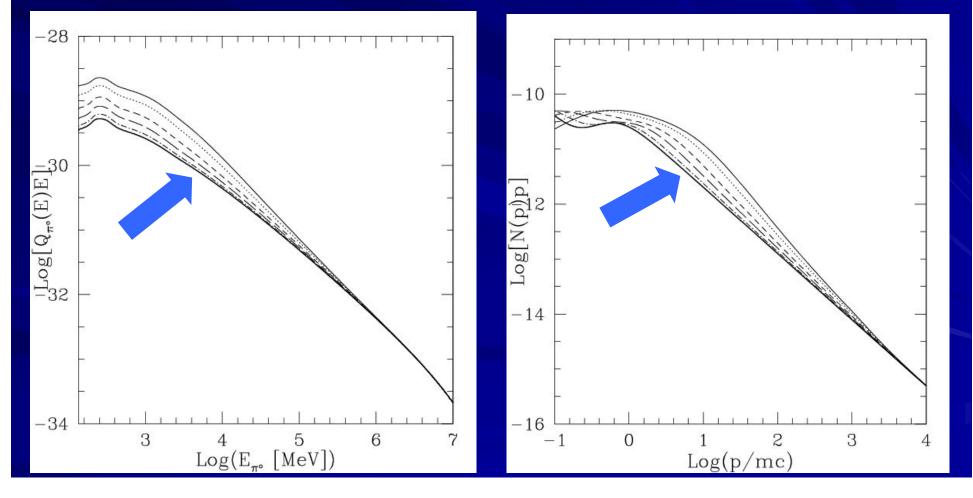


$N_p(p,t), N_e \pm (p,t), W(k,t), Q_e \pm (p,t), Q_{\pi}(p,t)$



n_{th}, T, B_o, N_p(p,0), I(k)

► $N_p(p,t), N_e \pm (p,t), W(k,t), Q_e \pm (p,t), Q_{\pi}(p,t)$



Dissipation of turbulence in the ICM (quasi-linear theory)

Brunetti & Lazarian 2007

• TTD Resonant acceleration

$$\omega - k_{\parallel} v_{\parallel} - n \frac{\Omega_o}{\gamma} = 0$$
 $n = 0$

(provided there is an additional source of pitch angle scattering)

Damping-rate

Energy changing-rate

$$\Gamma = -i \left(\frac{E_i^* K_{ij}^a E_j}{16\pi W} \right)_{\omega_i = 0} \omega_r \qquad \qquad \int d^3 p E_\alpha \left(\frac{\partial f_\alpha(p)}{\partial t} \right) = \int d\mathbf{k} \Gamma^\alpha(k, \theta) \mathcal{W}(\mathbf{k})$$

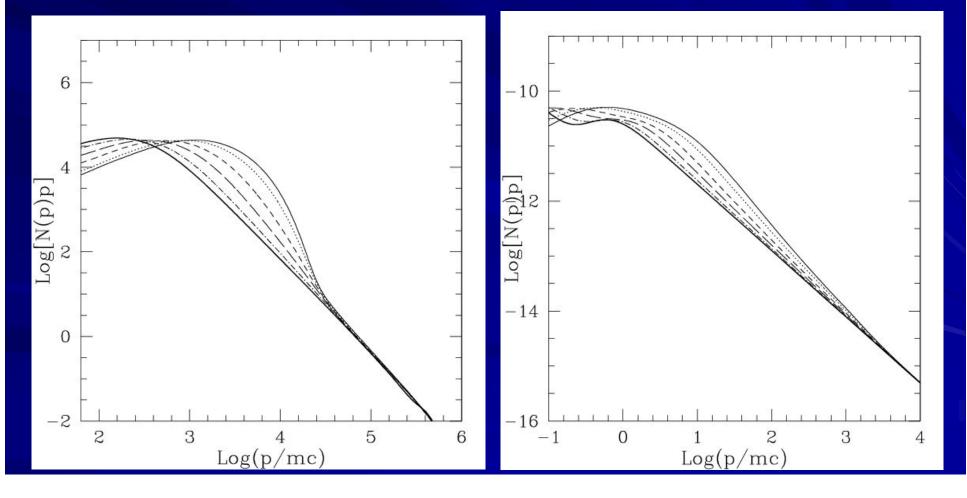
Energy damping-rate

Diffusion coeff. In momentum space

$$D_{pp}(p) = \frac{\pi^2}{2c} p^2 \frac{1}{B_o^2} \int_0^{\pi/2} d\theta V_{ph}^2 \frac{\sin^3(\theta)}{|\cos(\theta)|} \mathcal{H}\left(1 - \frac{V_{ph}/c}{\cos\theta}\right) \left(1 - \left(\frac{V_{ph}/c}{\cos\theta}\right)^2\right)^2 \int dk \mathcal{W}_B(k)k$$

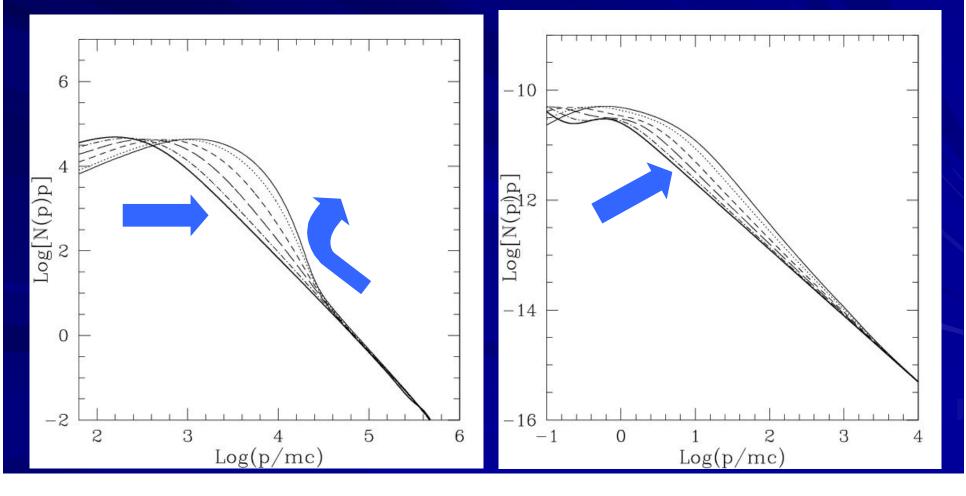
n_{th}, T, B_o, N_p(p,0), N_e(p,0), I(k)

► $N_p(p,t)$, $N_e \pm (p,t)$, W(k,t), $Q_e \pm (p,t)$, $Q_{\pi}(p,t)$



n_{th}, T, B_o, N_p(p,0), N_e(p,0), I(k)

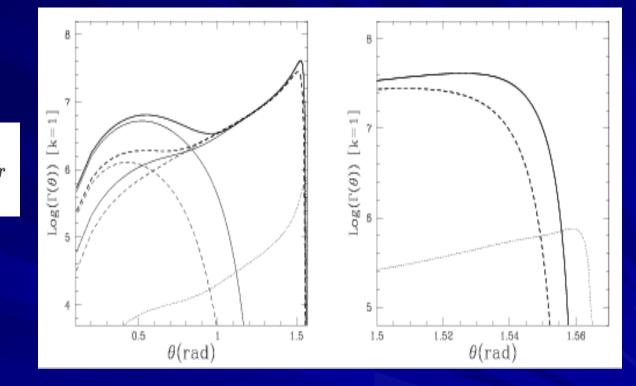






Dissipation of turbulence in the ICM (quasi-linear theory)

$$\Gamma = -i \left(\frac{E_i^* K_{ij}^a E_j}{16\pi W} \right)_{\omega_i = 0} \omega_i$$



Line-bending efficiency >> damping efficiency $\tau_{bb}(\mathbf{k})^{-1} \sim V_{1A} / l_A \qquad \tau_d^{-1} = \Gamma(\mathbf{k})$ Isotropic Effective Damping $l_{diss} \approx 100 \text{ pc}$

Dissipation of turbulence in the ICM (Hydro-turbulence)

Dissipation cut-off developes at scale :

 $\tau_{\rm kk}({\sf I}) \geq \tau_{\rm vis}({\sf I}) = {\sf I}^2/\nu_{\rm K}$

$$\nu_K \sim \frac{1}{3} v_i \times \min\left\{l_A, \, l_{mfp}\right\}$$

(e.g., Lazarian 2006)

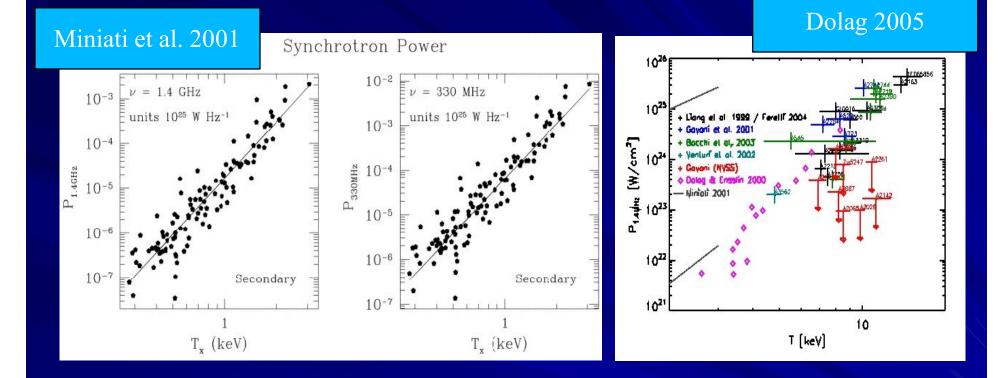
$$Re = \frac{V_o L_o}{\nu_K} > 10^3$$

$$l_{diss} \sim L_o (Re)^{-\frac{3}{4}}$$

Collisionless dissipation !

mfp

Statistical properties of RH

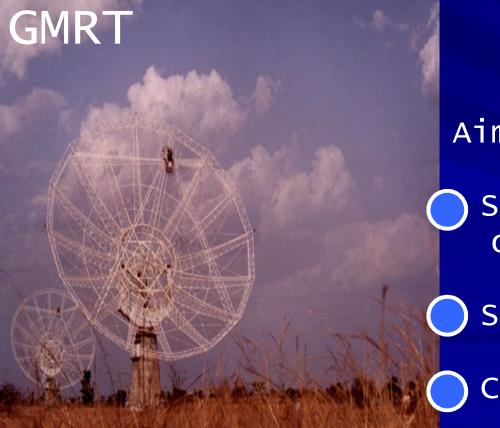


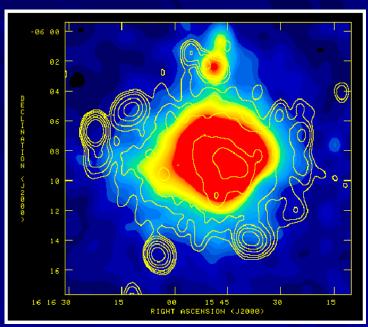
In case electrons are secondaries :

RH are not transient phenomena

GC with similar thermal properties have similar radio properties

Diffuse Radio Emission & origin of emitting electrons





Aims:

Statistical properties of RH

Spectral studies of RH

Connection with mergers

Stochastic Particle Acceleration (formalism)

Brunetti +al. 2004, Brunetti & Blasi 2005, Brunetti +al. In prep

Electrons/Positrons

$$\frac{\partial N_e(p,t)}{\partial t} = \frac{\partial}{\partial p} \left(N_e(p,t) \left[\left(\frac{dp}{dt} \right)_{rad} + \left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_{rad} + \left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_{rad} + \left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_{rad} + \left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_{rad} + \left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + Q_e(p,t) \left[\left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right] + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_e(p,t)}{\partial p} \right)$$

Protons

Q_e : secondaries from (CR)p-p collisions

$$\frac{\partial N_p(p,t)}{\partial t} = \frac{\partial}{\partial p} \left(N_p(p,t) \left[\left(\frac{dp}{dt} \right)_i - \frac{2}{p} D_{pp} \right] \right) + \frac{\partial}{\partial p} \left(D_{pp} \frac{\partial N_p(p,t)}{\partial p} \right) + Q_p(p,t)$$

Waves

$$\frac{\partial \mathcal{W}(k,t)}{\partial t} = \frac{\partial}{\partial k} \left(k^2 D_{kk} \frac{\partial}{\partial k} \left(\frac{\mathcal{W}(k,t)}{k^2} \right) \right) - \sum_i \Gamma_i(k,t) \mathcal{W}(k,t) + I(k,t)$$

most could be with CR protons

Stochastic particle RE-acceleration

• Stochastic turbulent compression (e.g. Ptuskin 1988)

(provided $I >> I_{mfp} \& V_I^2 << c_s^2$)

$$D_{pp} \simeq \frac{2}{9} D p^2 \frac{V_o^2}{L_o^{2/3}} \int_{1/L_o}^{1/l_{cut}} \frac{dy \, y^{1/3}}{c_s^2 + D^2 y^2}$$

With

$$D \sim \frac{c}{3} \max \left\{ l_{cut} , \min\{l_A , l_{mfp}\} \right\}$$

Also ... efficient source of particle pitch-angle isotropization

Dissipation of turbulence in the ICM (quasi-linear theory)

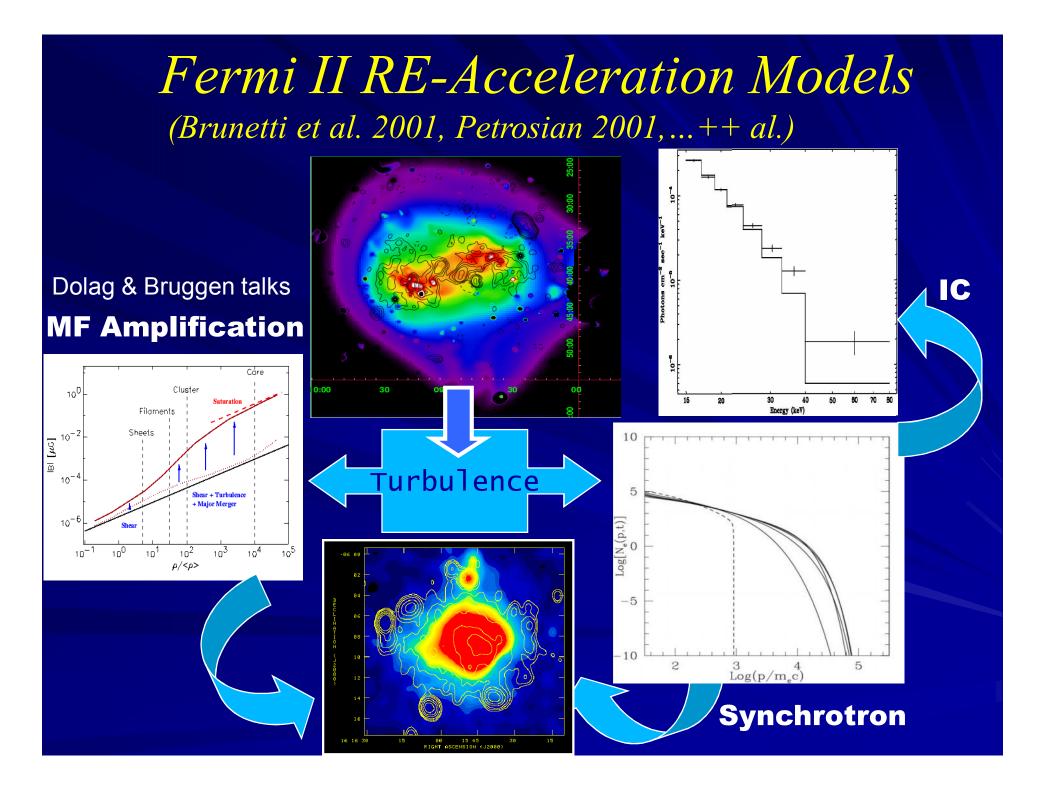
$$\omega - k_\parallel v_\parallel - n rac{\Omega_o}{\gamma} = 0$$

Small scale n = +/-1, +/-2, +/-3, ..Large scale n = 0

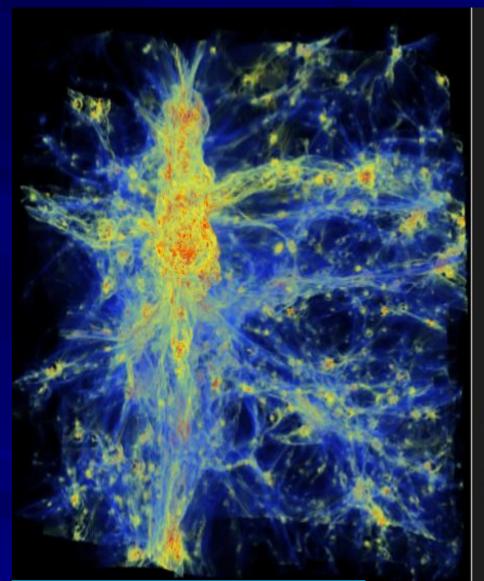
Transit-Time-Damping n = 0 (e.g. Schlickeiser & Miller 1998)

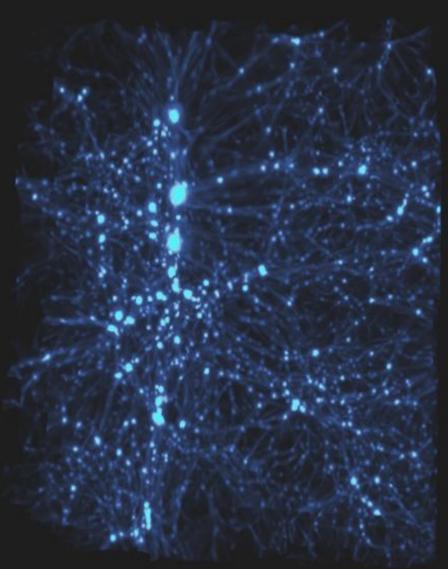
$$\omega = k_{||}v_{||} \qquad \& \qquad V_{\rm ph}^2 = \frac{c_s^2}{2}\frac{\beta_{pl} + 1}{\beta_{pl}} \left\{ 1 + \sqrt{1 - 4\left(\frac{k_{||}}{k}\right)^2 \frac{\beta_{pl}}{(1 + \beta_{pl})^2}} \right\}$$

Brunetti & Lazarian 20



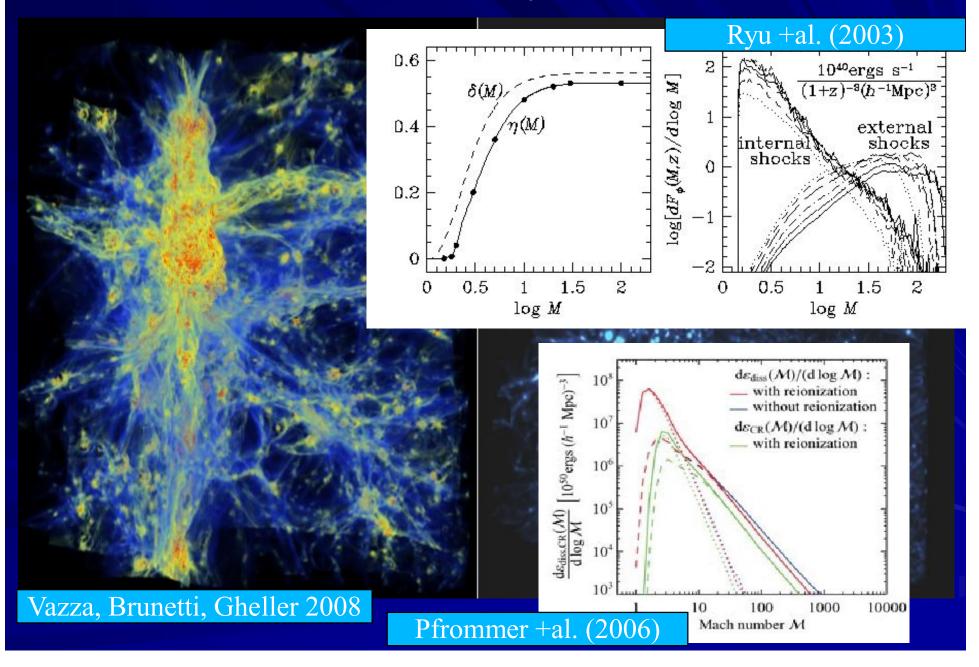
Shocks in Galaxy Clusters



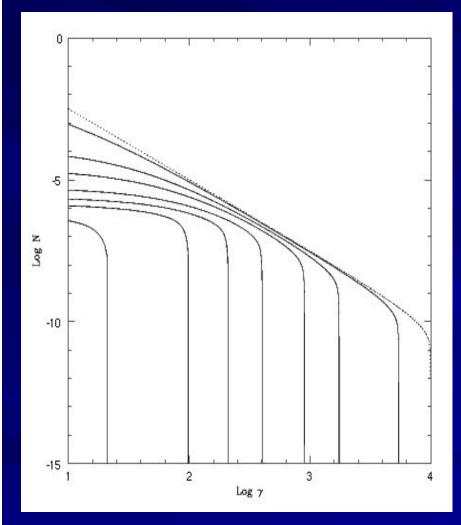


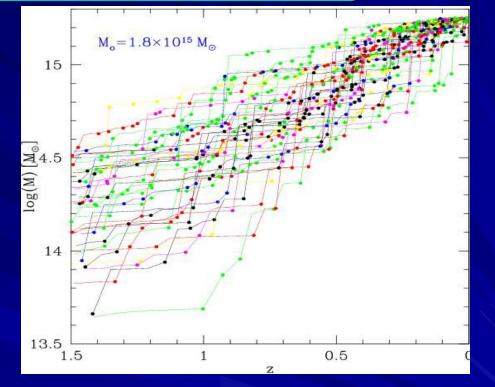
Vazza, Brunetti, Gheller 2008

Shocks in Galaxy Clusters

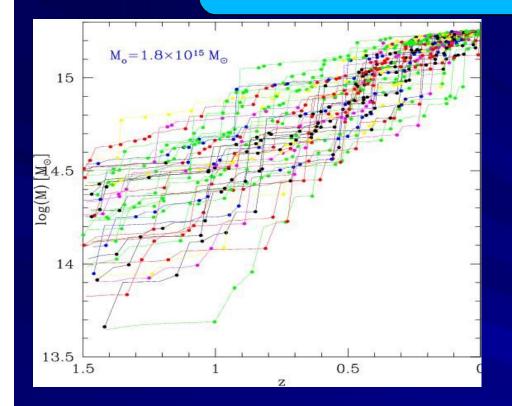


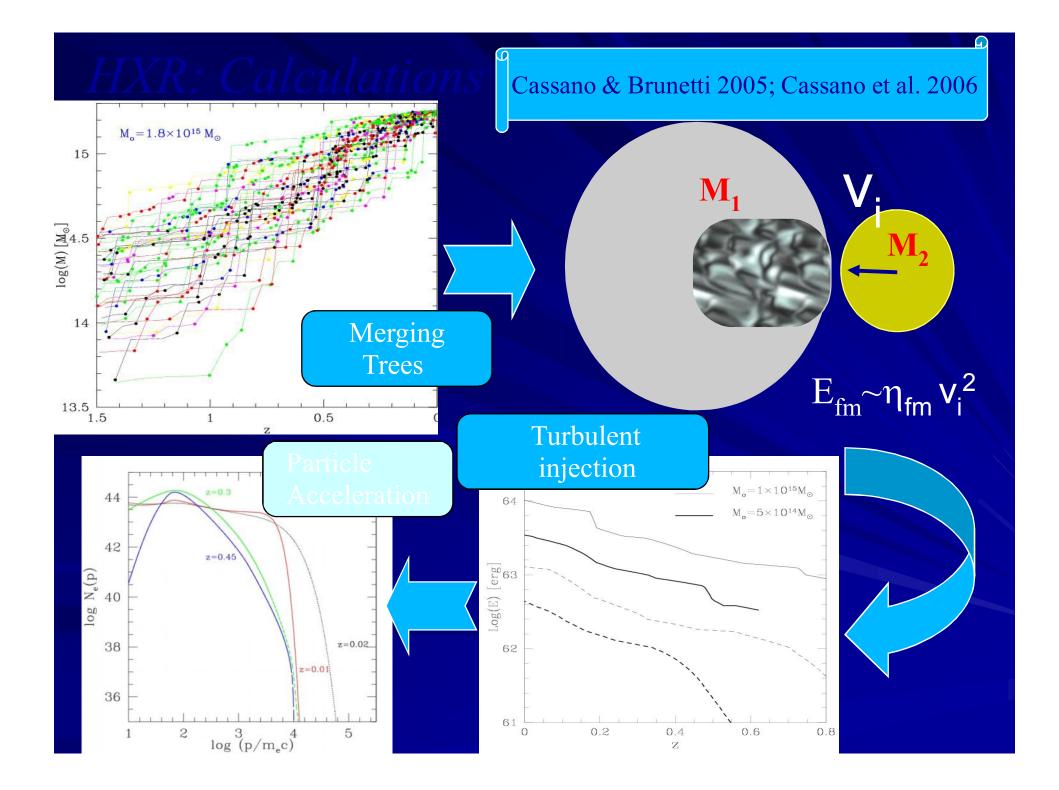




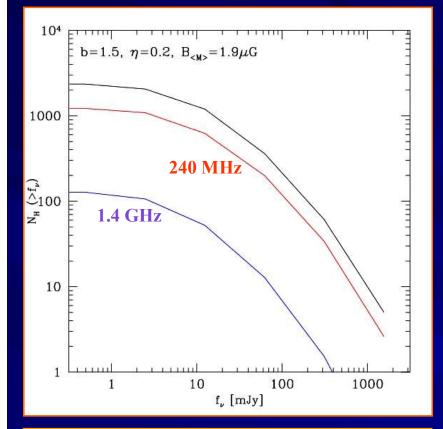


CR Emission



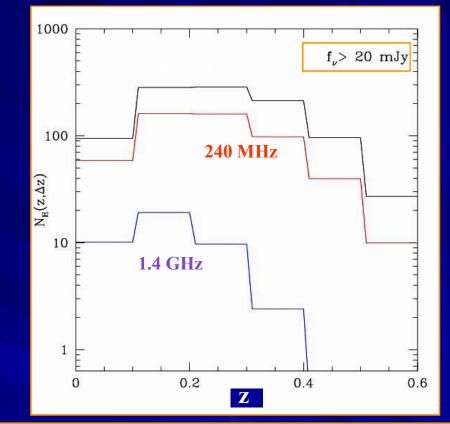


Number Counts of RHs at lower radio frequencies



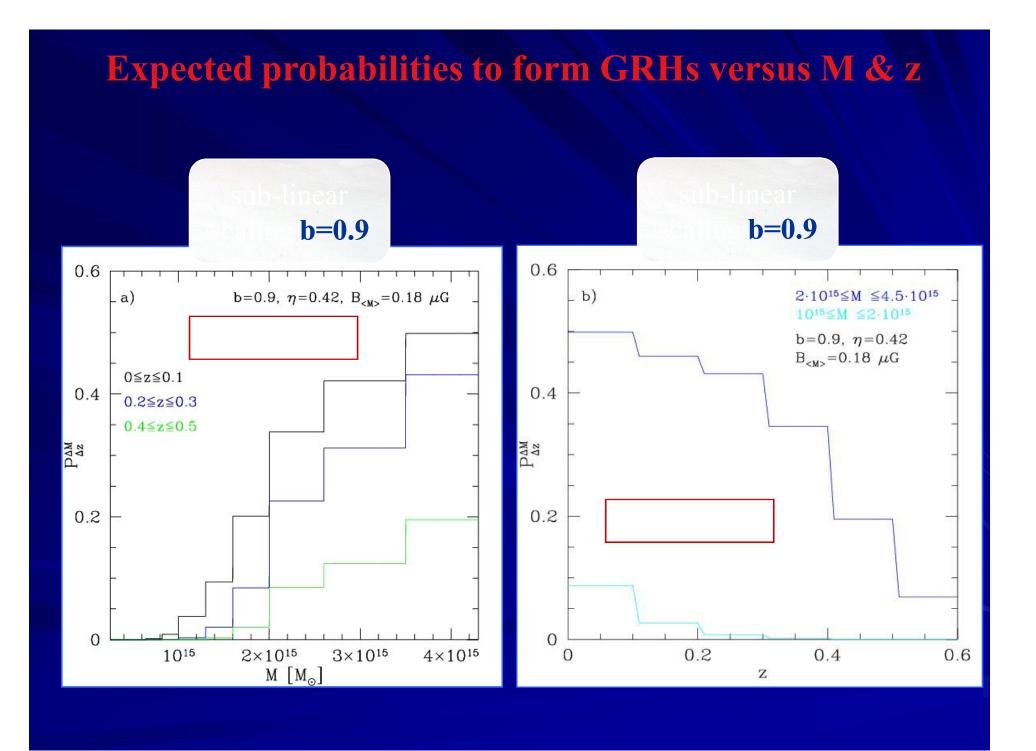
The number of expected RHs increases at lower frequencies by about a factor 10!

LOFAR should be able to catch the bulk of RHs!



The bulk of RHs emitting at GHz frequencies is expected at relatively low redshift: 0.1-0.3.

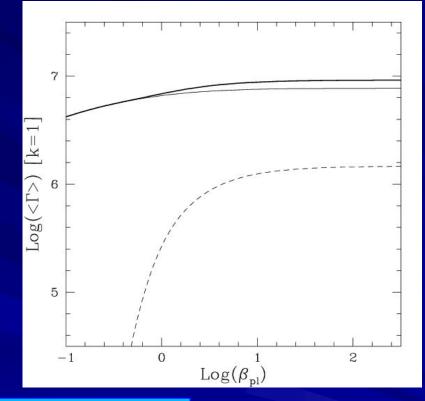
At lower radio frequencies a number of RHs is expected to be discovered at relatively higher redshifts z>0.4.



Dissipation of turbulence in the ICM (quasi-linear theory)

(e.g., Melrose 1968; Barnes 1969; Barnes & Scargle 1973)

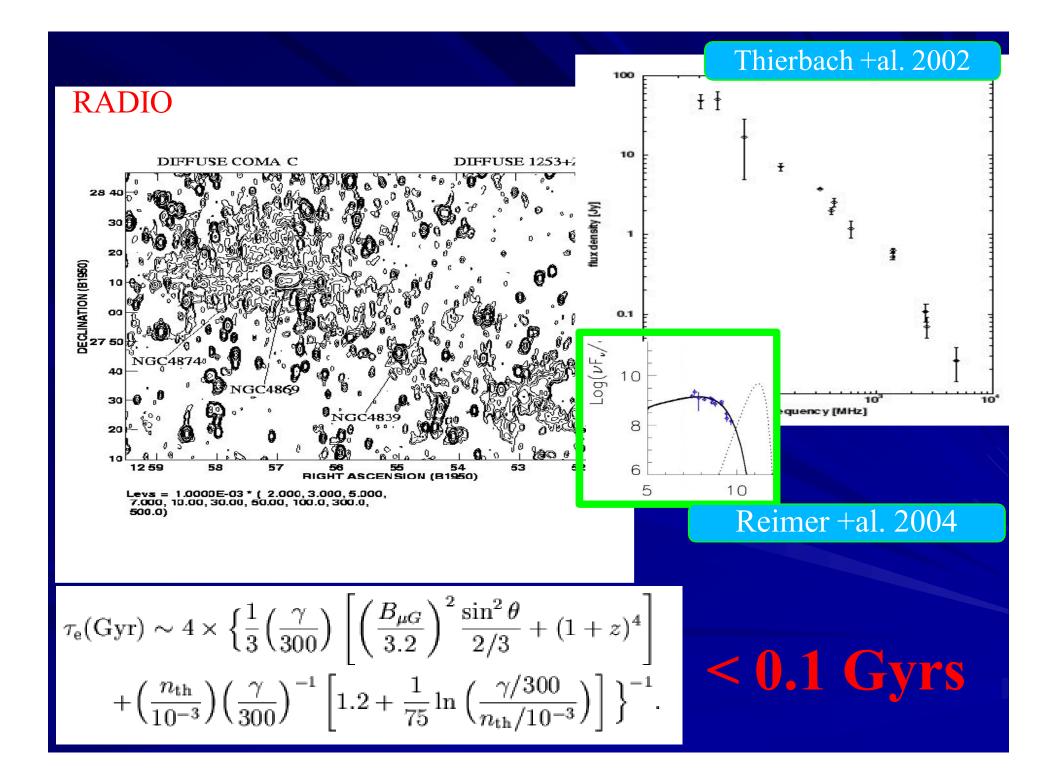
$$\Gamma = -i \left(\frac{E_i^* K_{ij}^a E_j}{16\pi W} \right)_{\omega_i = 0} \omega_r$$

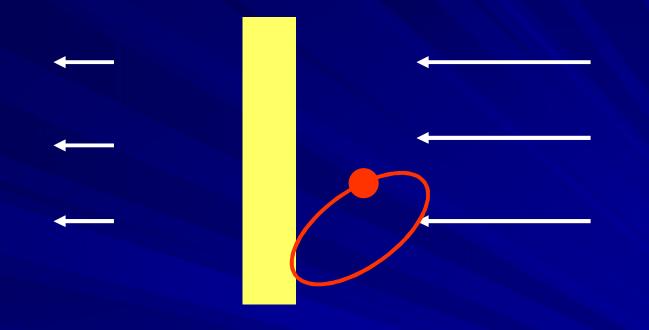


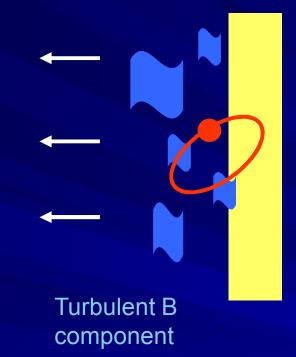
Damping is dominated by thermal electrons since they are faster than the phase velocity of the modes . CR are not important provided they are not a dominant component in the ICM *Damping is simple !*

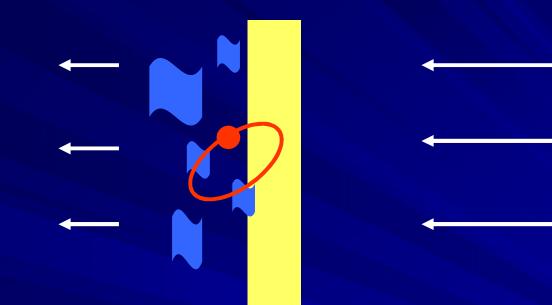
 (\mathbf{k})

Brunetti & Lazarian 20

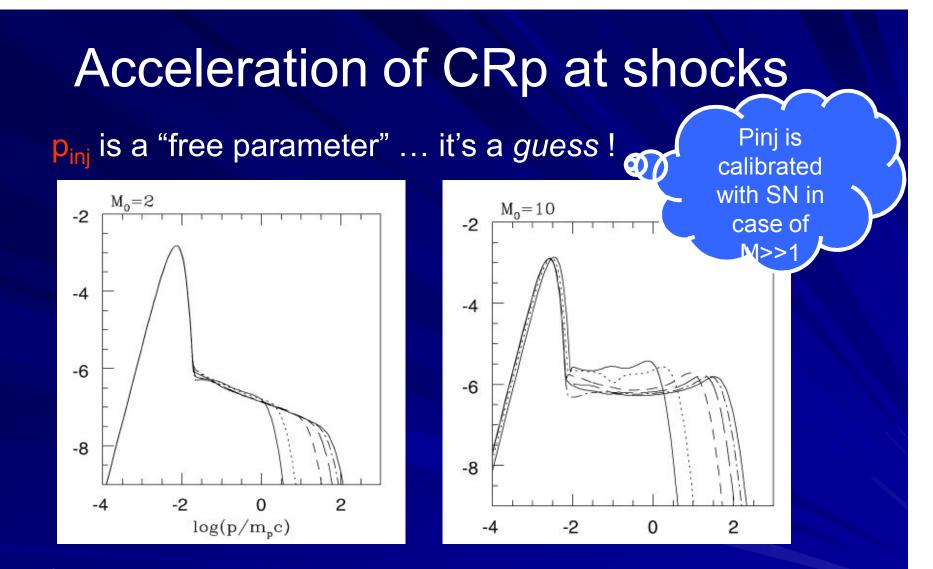








The level of amplification of the turbulent field in the post shock region is necessary to calculate the efficiency of particle acceleration & the minimum momentum of particles that can be accelerated at the shock. This physics is not yet under control !

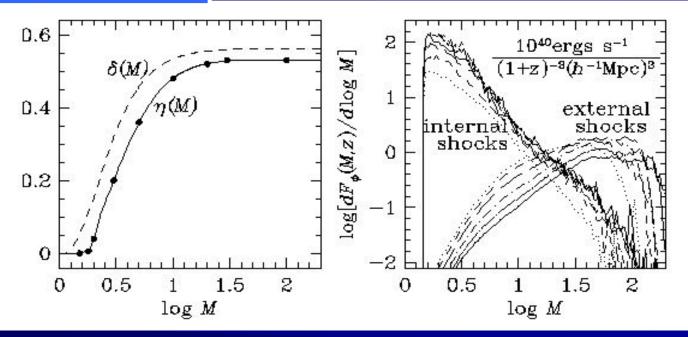


(Kang & Jones 2002; see also Malkov 1997, Kang & Jones 2005, 07, Amato & Blasi 2006)

However going from $p_{inj} = 3.5 p_{th}$ to 3.8 reduces the acceleration efficiency by a factor of 10 !

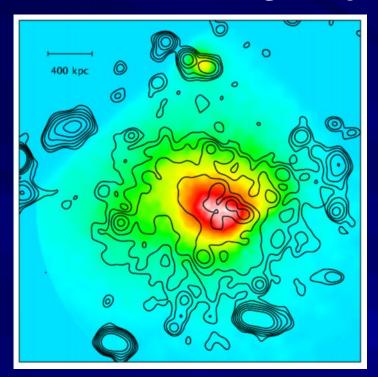
The efficiency strongly depends on Mach number !!!

Ryu et al. 2003

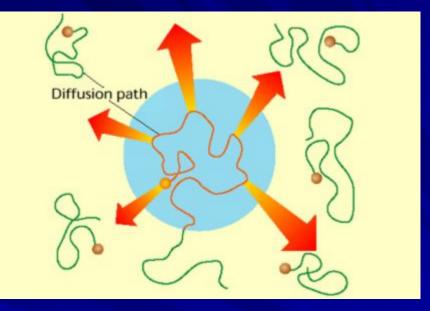


Having M=2 or M=2.5 implies a difference of 10 in terms of acceleration efficiency

Origin of Radio Halos



 $R \approx 2VTD$

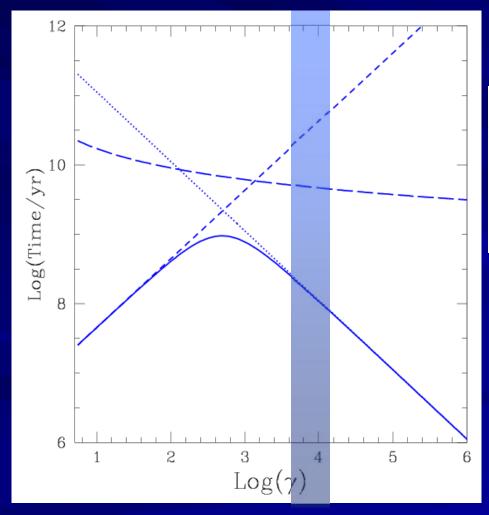


$$\frac{\partial n_p(E_p, r, t)}{\partial t} - D(E) \nabla^2 n_p(E_p, r, t)$$

$$D(E_p) = \frac{1}{3} r_L c \frac{B^2}{\int_{1/r_L}^{\infty} dk P(k)}$$

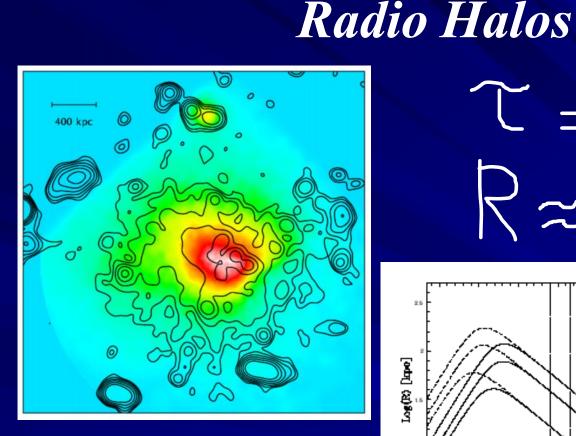
Physics of CR Leptons

$(dE/dt) \sim E / Time$



$$\begin{aligned} \tau_{\rm e}({\rm Gyr}) &\sim 4 \times \left\{ \frac{1}{3} \left(\frac{\gamma}{300} \right) \left[\left(\frac{B_{\mu G}}{3.2} \right)^2 \frac{\sin^2 \theta}{2/3} + (1+z)^4 \right] \\ &+ \left(\frac{n_{\rm th}}{10^{-3}} \right) \left(\frac{\gamma}{300} \right)^{-1} \left[1.2 + \frac{1}{75} \ln \left(\frac{\gamma/300}{n_{\rm th}/10^{-3}} \right) \right] \right\}^{-1}. \end{aligned}$$

The life-time of electrons depends on quantities that are well measured



The diffusion time of the emitting electrons necessary to cover Mpc distances is 100 times larger than their radiative life-time

 $\mathcal{T} = \mathcal{T}_{\mathsf{F}}$ $Z \approx 2VTD$

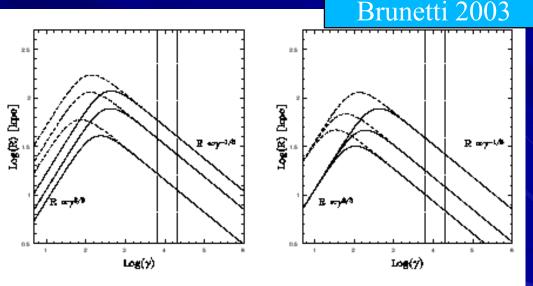
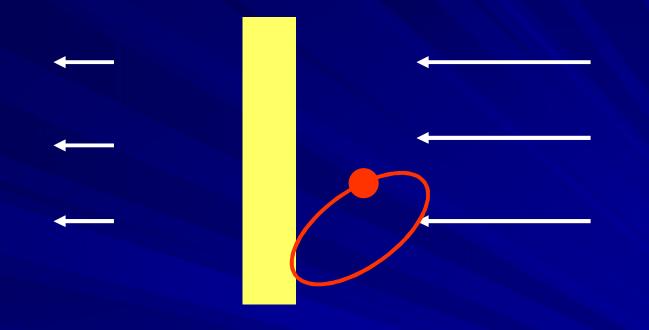
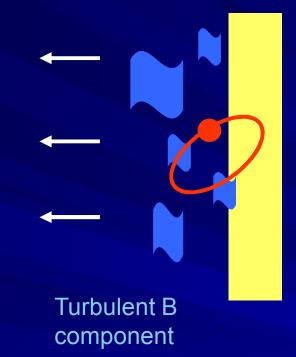
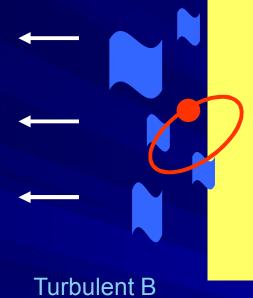


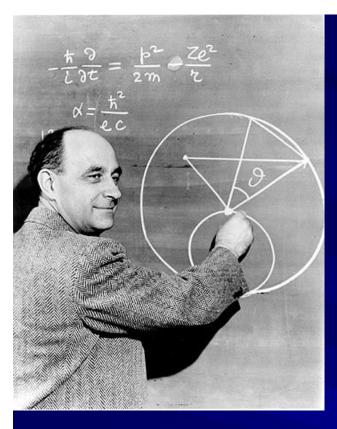
Figure 2. **Panel a)** : Diffusion lengths are reported as a function of γ of the electrons. Calculations are performed at z = 0, for $n_{\rm th} = 10^{-3}$ cm⁻³ (solid lines) and 10^{-4} cm⁻³ (dashed lines) assuming (from the bottom) B=5, 1, and 0.1 μG . **Panel b)** : Diffusion lengths as in panel a), but assuming $B = 1\mu G$, and z = 0, 0.5, and 1.0 (from the top). The energy range of the radio emitting electrons is reported in both panels (vertical dotted lines).



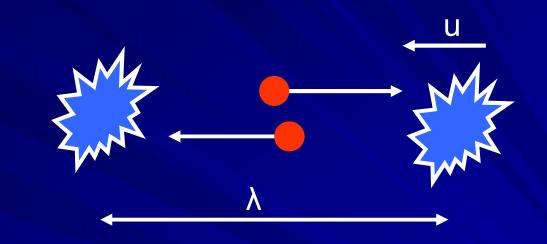




component



First order Fermi Mechanisms (Fermi 1949)



C A

<u>+ C</u>

P=2P~~

Frequency of collisions:

Energy gain per collisions:

Energy gain per second:

Physics of CR Leptons

Sarazin 1999 ; Brunetti 2003 for reviews

 $(dE/dt) / m_e c^2 = b = rate of energy losses in units of m_e c^2$

$$b_{\rm IC}(\gamma) = \frac{4}{3} \frac{\sigma_{\rm T}}{m_e c} \gamma^2 U_{\rm CMB} = 1.37 \times 10^{-20} \gamma^2 (1+z)^4 \, {\rm s}^{-1} \, ,$$

Photon Collisions

$$b_{\rm syn}(\gamma) = \frac{4}{3} \frac{\sigma_{\rm T}}{m_e c} \gamma^2 U_B = 1.30 \times 10^{-21} \gamma^2 \left(\frac{B}{1 \ \mu \rm G}\right)^2 \, \rm s^{-1} \; ,$$

Particle Collisions

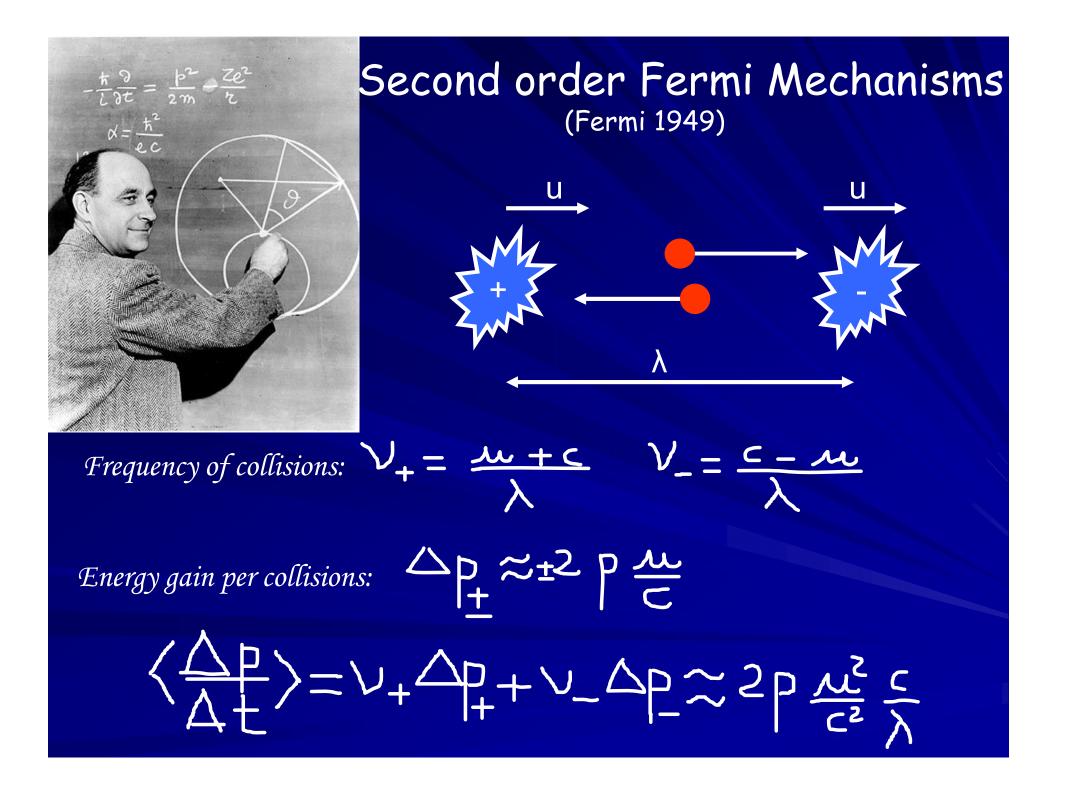
$$b_{\text{Coul}}(\gamma) \approx 1.2 \times 10^{-12} n_e \left[1.0 + \frac{\ln (\gamma/n_e)}{75} \right] \text{s}^{-1}$$
,

 $b_{\rm brem}(\gamma) \approx 1.51 \times 10^{-16} n_e \gamma [\ln (\gamma) + 0.36] \, {\rm s}^{-1}$,

IV - Conclusion

(a) Synchrotron emission from secondaries is an unavoidable process in the IGM

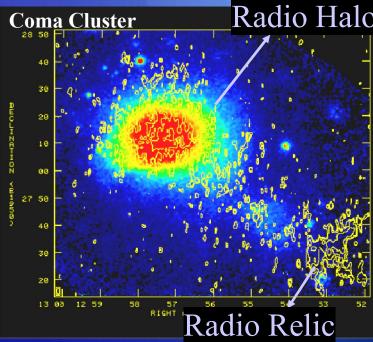
(b) Still this is probably not the (unique) process at the origin of Radio Halos ... the level of CR protons is below that necessary to explain Radio Halos



Non-thermal components

Observational evidences

Diffuse synchrotron radio sources from the ICM (not associated with any individual galaxy): Radio Halos $(L_{1.4GHz} \sim 10^{24} - 10^{26} h_{70}^{-2} Watt/Hz)$



Radio Halo
 • steep spectrum sources (α ~ 1.1-1.5)
 • low surface brightness (μJy arcsec⁻² at 1.4 GHz)
 • at the cluster centre
 • generally regular shape (mimic the X-ray morphology) (-Mpc size)
 Radio Relics (L_{1.4GHz} ~10²³-10²⁵ l²₇₀ Watt/Hz)
 • steep spectrum sources (α ~ 1.1-1.5)
 • at the cluster outskirts

elongated morphology + polarised

Halos and Relics prove the presence of non-thermal componenets, GeV electrons (γ ~10⁴) and μ G magnetic field, mixed with the thermal ICM on Mpc

CR+B : important not only for particle acceleration

Thermal conduction and kin. viscosity in the ICM (e.g., Lazarian 2006)

Heating of the ICM and "cooling flow" problem (e.g., Fujita , Matsumoto, Weda 2004; Guo & Oh 2008)
Diffusion and transport of metals in the ICM (e.g., Voigt & Fabian 2004; Rebusco +al. 2005)
B-Amplification from Cosmological seed fields (e.g., Dolag +al. 1999,02; Subramanian +al. 2006)

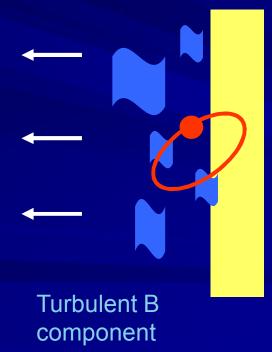
Diffusion and scattering of HE & UHECR in the Universe (e.g., Sigl +al. 2005; Dolag +al. 2005)

$$\bigvee (P) \sim P^{-d} \qquad \delta = 2 \frac{\mathcal{M}^2 + 1}{\mathcal{M}^2 - 1}.$$

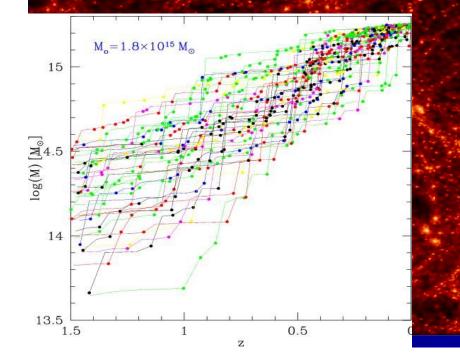
$$J(P) \propto P^{-d}$$

$$\delta = 2\frac{\mathcal{M}^2 + 1}{\mathcal{M}^2 - 1}.$$

Linear Theory (e.g.,Blandford & Eichler 1987)



Mergings are the most energetic events in the Universe ≈ few 10⁶³ergs in a crossing time



1E 0657-56

63-64

Markevitch et al.

z=0.3

0.5 Mpc

